



Experimental investigation of geodesic acoustic modes on JET using Doppler reflectometry

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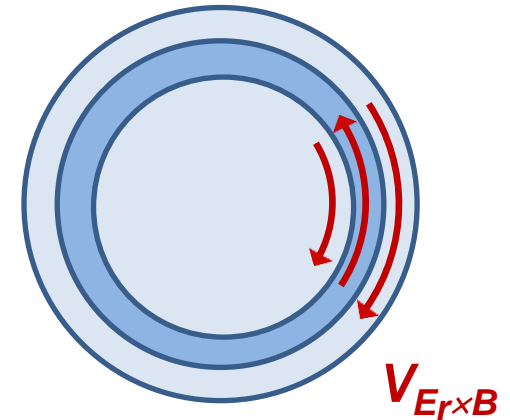
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** See the Appendix of F. Romanelli et al., Proceedings of the 25th IAEA Fusion Energy Conference 2014, Saint Petersburg, Russia*

Geodesic Acoustic Modes



- ❑ GAMs are symmetric ($n = 0, m = 0$) and radially-localized ($k_r \neq 0$) potential structures with $f_{\text{GAM}} \propto c_s/R$
- ❑ $n = 0, m = 1$ density and $n = 0, m = 2$ magnetic structure
- ❑ Oscillating $E_r \times B$ flows driven by turbulence
- ❑ May play role in triggering the L-H transition
- ❑ Strongly damped at low q (exist edge plasma)
- ❑ Extensively characterized in different devices





- ❑ Experiment aim: assessing the possible connection between properties of large scale flows (GAMs) and isotope physics
 - Characterize GAMs and local turbulence in H and D plasmas when approaching L-H transition
- ❑ Characterization performed in hydrogen Ohmic plasmas due to NBI unavailability

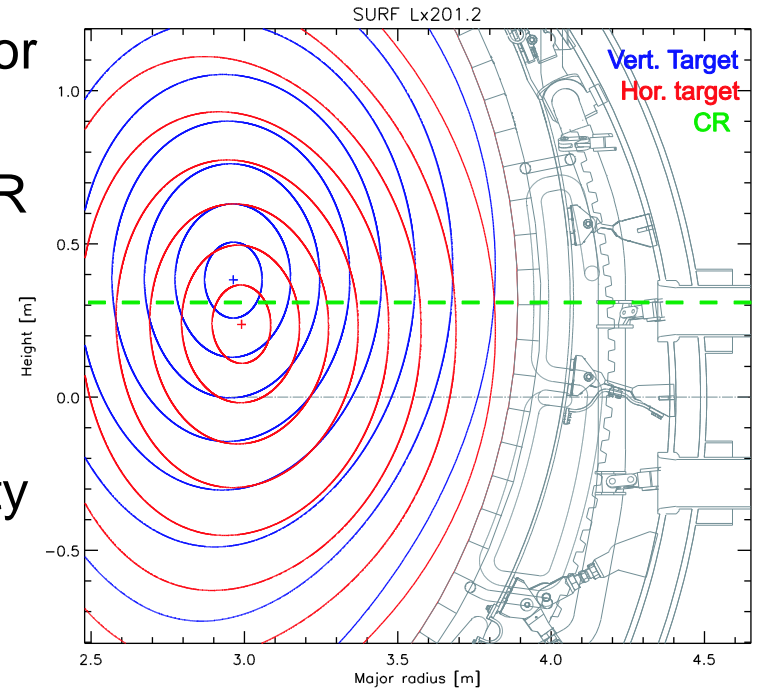
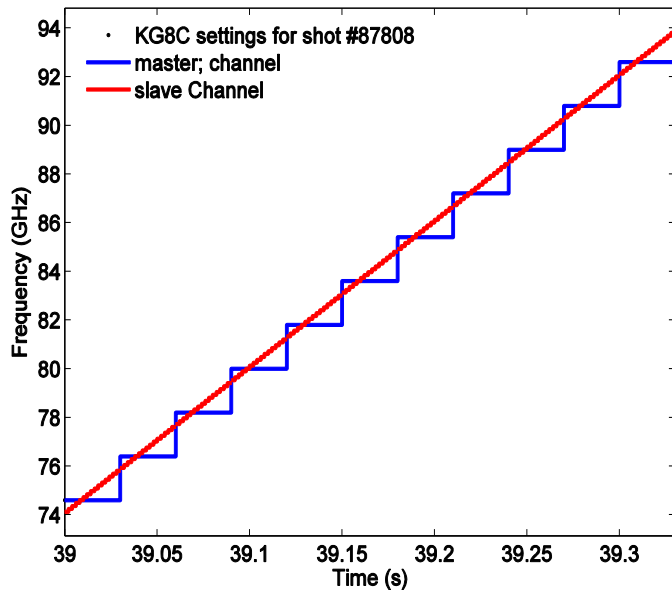
Dataset in H

| Pulse | B_T (T) | I_p (MA) | Density (10^{19} m^{-3}) |
|-------|-----------|------------|--------------------------------------|
| 87801 | 3 | 2.0 | 1.5 – 2.4 – 3.2 |
| 87802 | 3 | 2.5 | 1.6 – 2.8 – 3.6 |
| 87803 | 3 | 2.25 | 1.6 – 2.7 – 3.6 |
| 87804 | 3 | 1.75 | 1.5 – 2.6 |
| 87805 | 3 | 1.5 | 1.5 – 2.5 |
| 87808 | 3 | 2.75 | 2.8 – 3.9 |

JET radial correlation reflectometer (RCR)



- ❑ Radial correlation reflectometer designed for normal incidence
- ❑ For the vertical target configuration the RCR works for Doppler reflectometry
- ❑ Backscattered signal:
 - Amplitude: density fluctuation level
 - Doppler shift: turbulence lab frame velocity



RCR: two X-mode hopping channels

- ❑ **Master:** 11 point frequency sweep (30 ms)
- ❑ **Slave:** 15 point each master step (2 ms)
Full sweep: 330 ms
- ❑ Density range (3 T): $\sim 0.7 - 3 \times 10^{19} \text{ m}^{-3}$

GAM observation with RCR



Signals digitized: I, Q

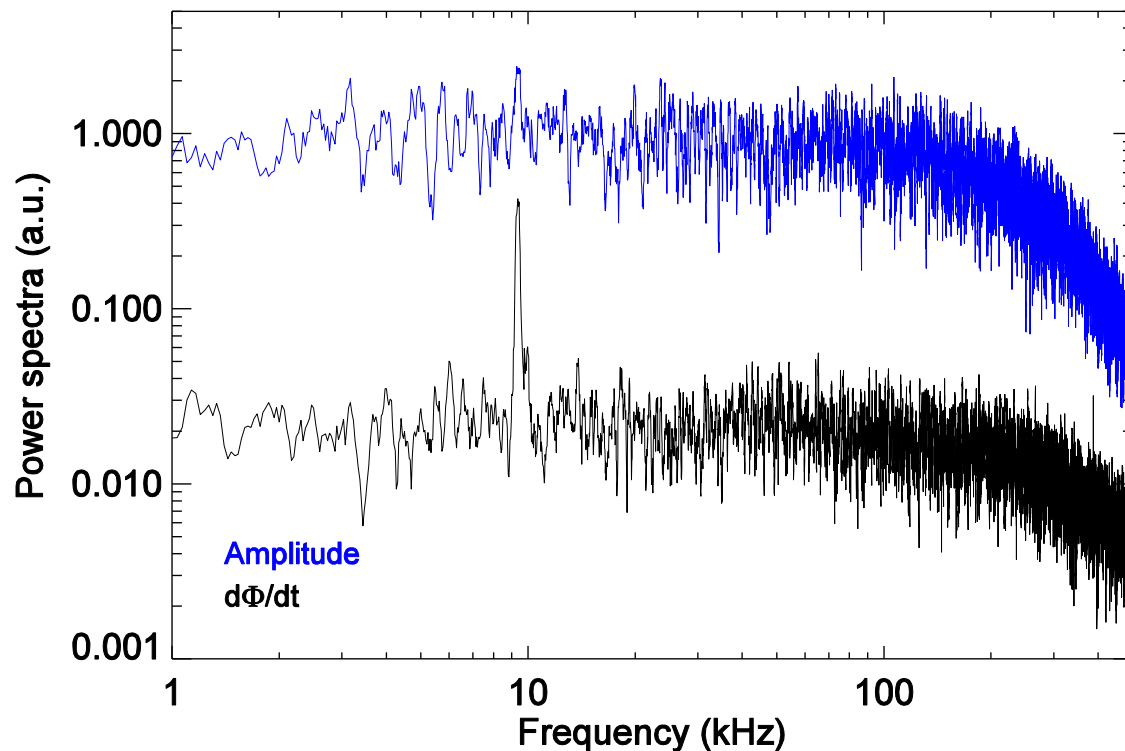
$A = \sqrt{I^2 + Q^2}$

$\phi = \arctan(Q/I)$

$\omega_D = d\phi / dt$

$v_{\perp} \approx \omega_D / k_{\perp}$

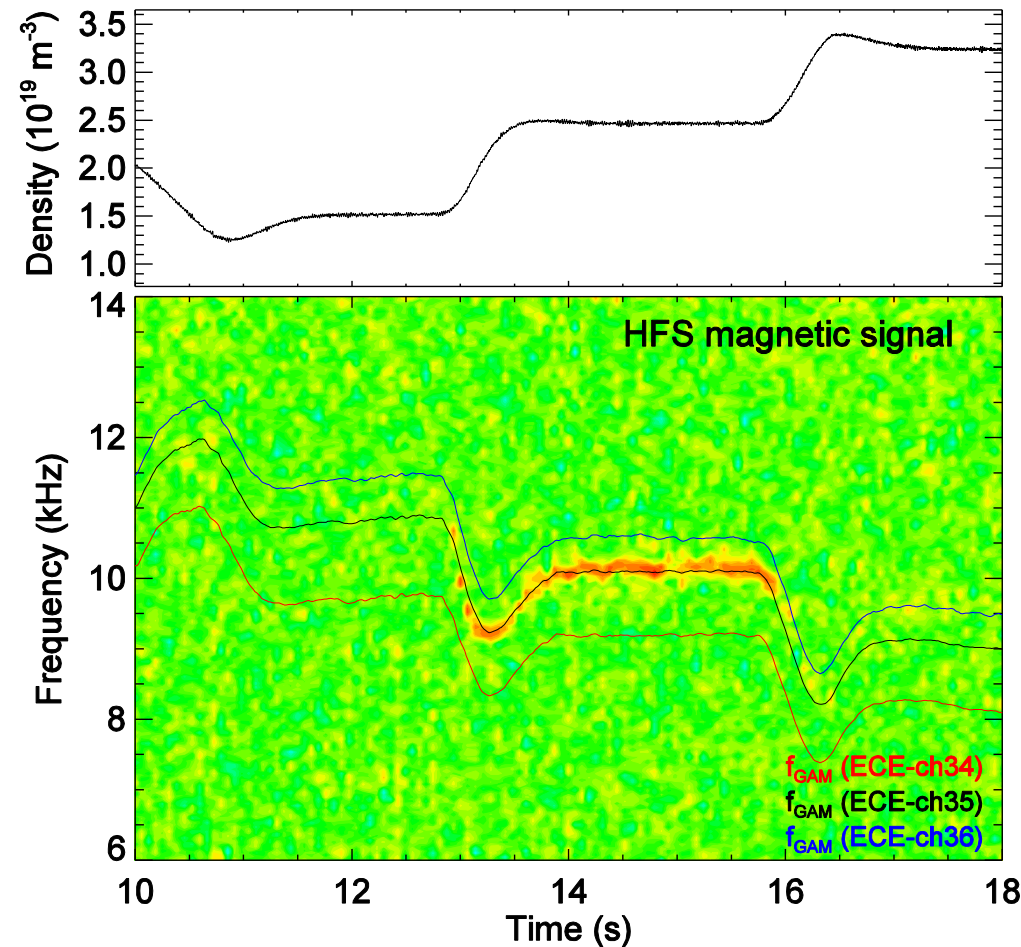
$k_{\perp} \approx 3 \text{ cm}^{-1}$ (from ray tracing calculations)



GAM peak clearly visible in $d\phi/dt$ at ~ 10 kHz but not in the amplitude signal

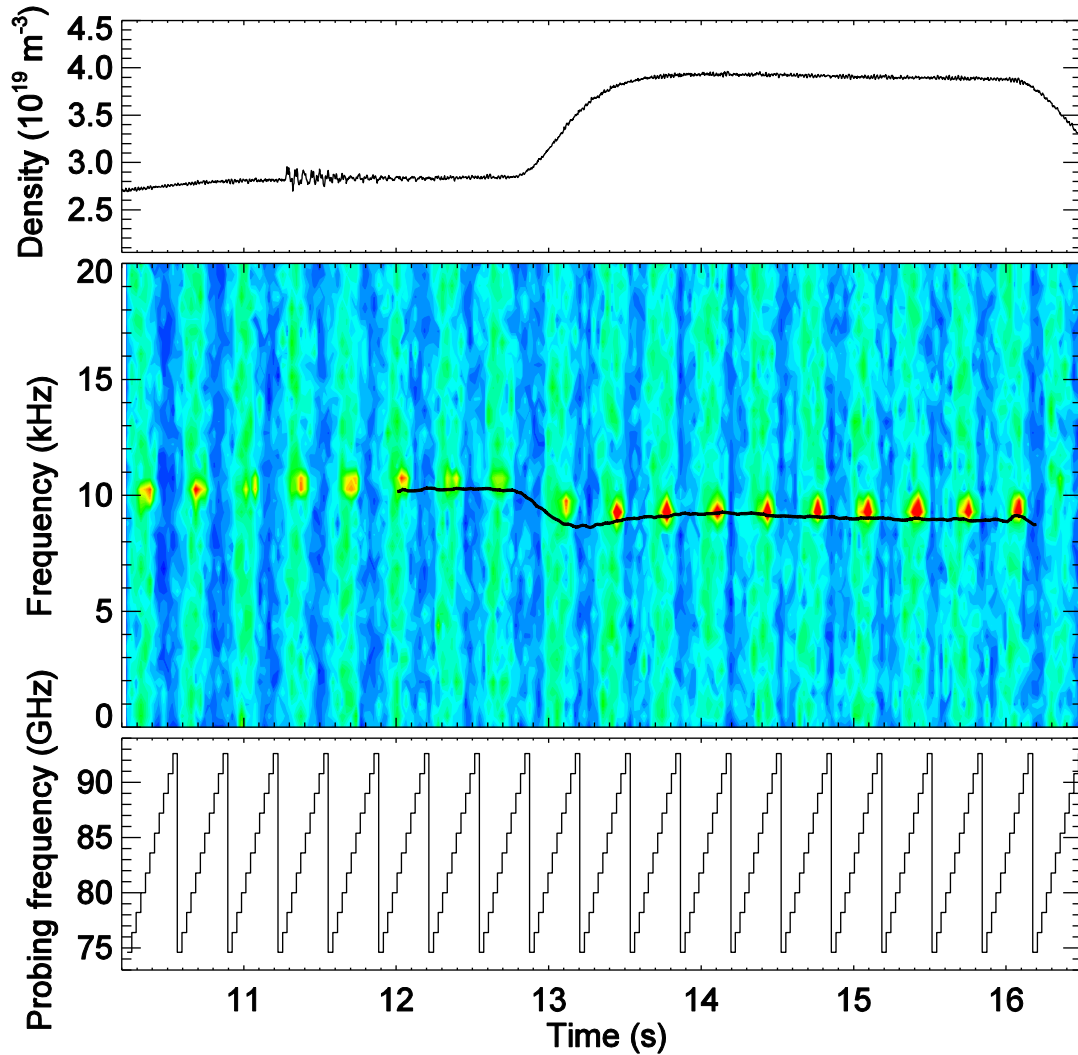
GAM amplitude \sim one order magnitude larger than background

GAM identification on magnetic signals



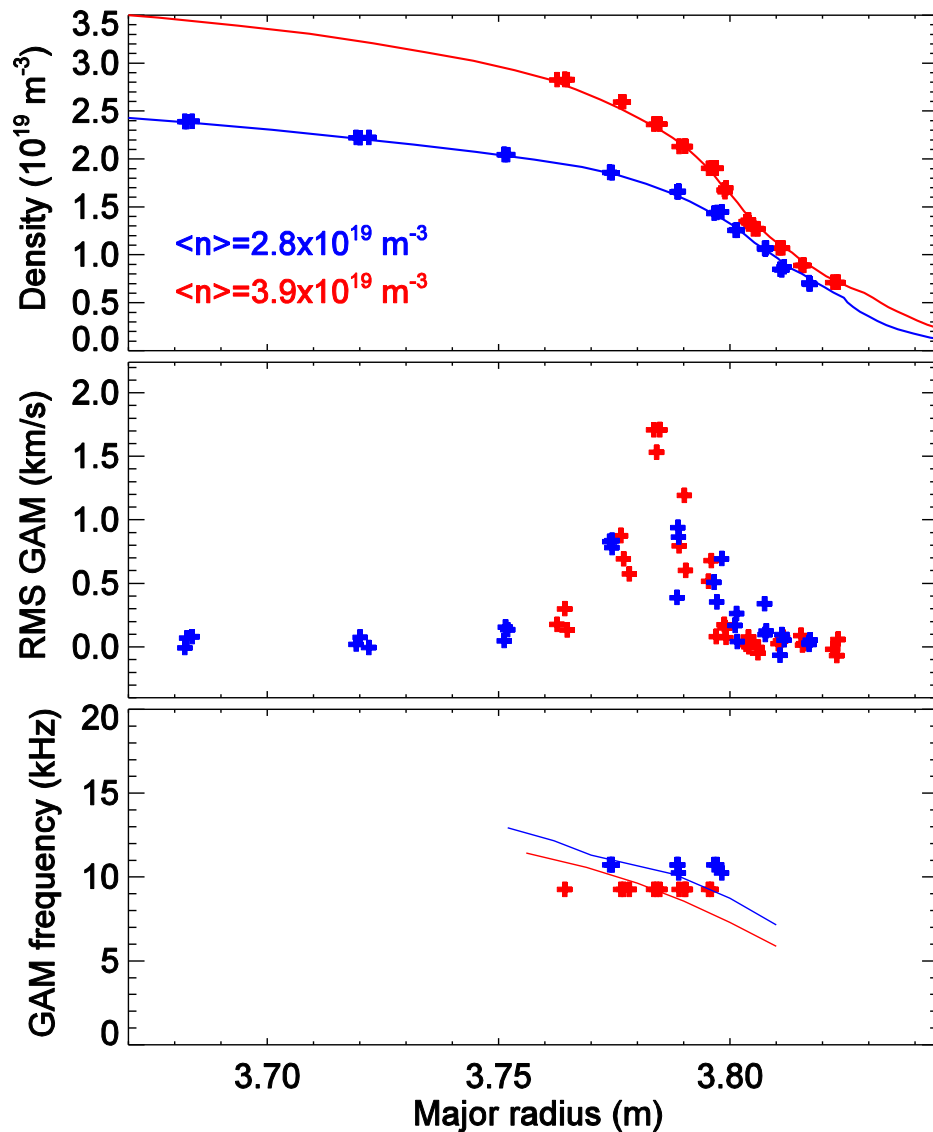
- GAM peak (~ 10 kHz) clearly seen magnetic signals mainly at HFS
- $m = 2$ structure identified
- Frequency follows the local T_e ($f_{\text{GAM}} = c_s / 2\pi R$, $T_i = T_e$, $\gamma_i = 1$)
- Narrow frequency peak ($\Delta f < 0.5$ kHz) suggesting either very localized mode or constant frequency across mode

GAM evolution with RCR



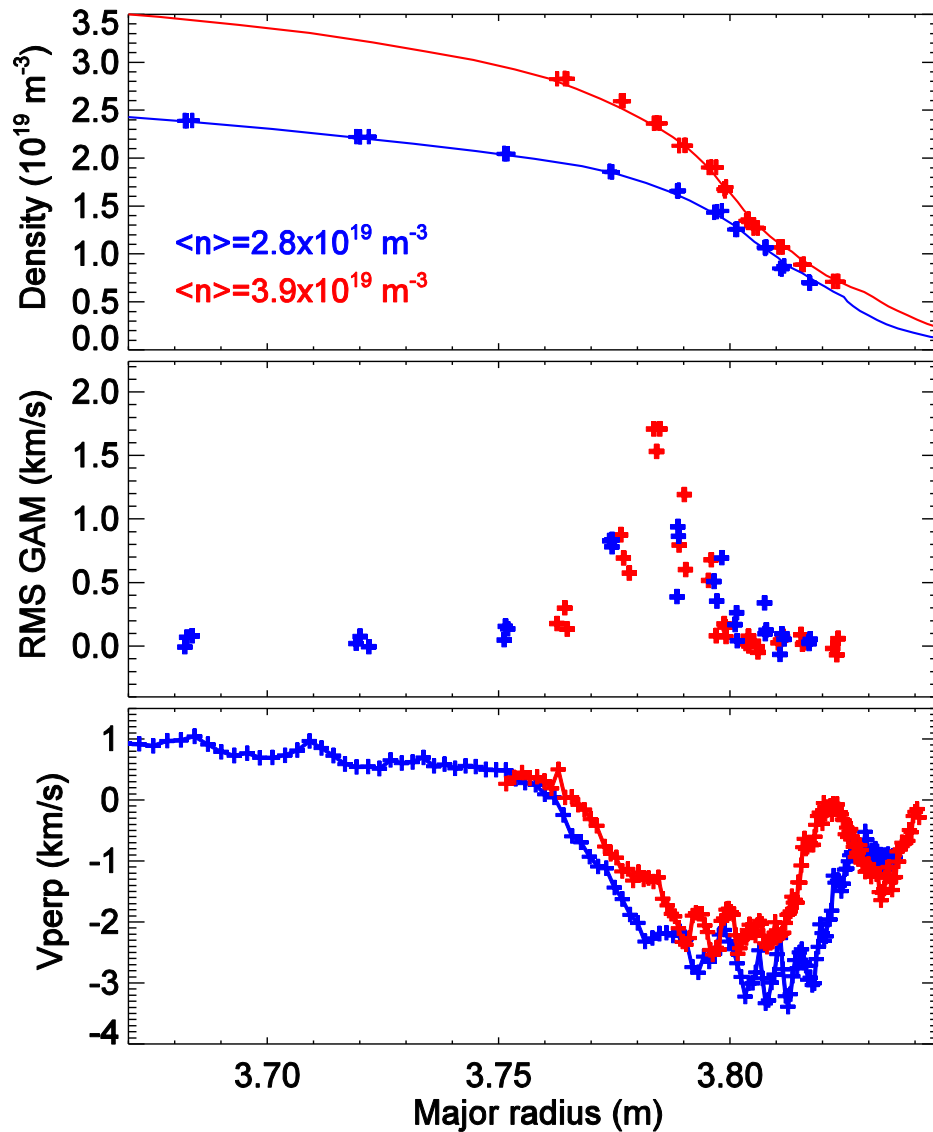
- GAM clearly observed in $d\phi/dt$ signal
- Frequency depends on local temperature
- GAM amplitude:
 - $|\text{FFT}(d\phi/dt)|$
 - rms of $d\phi/dt$ bandpass filtered (f_{GAM})

GAM localization and amplitude



- Profile reflectometer used to determine RCR radial measurement location
- GAM located at the edge density gradient region in a narrow layer of ~ 2 cm
- GAM frequency roughly constant with radius
- Larger f_{GAM} radial variation expected from T_e profile

GAM localization and amplitude

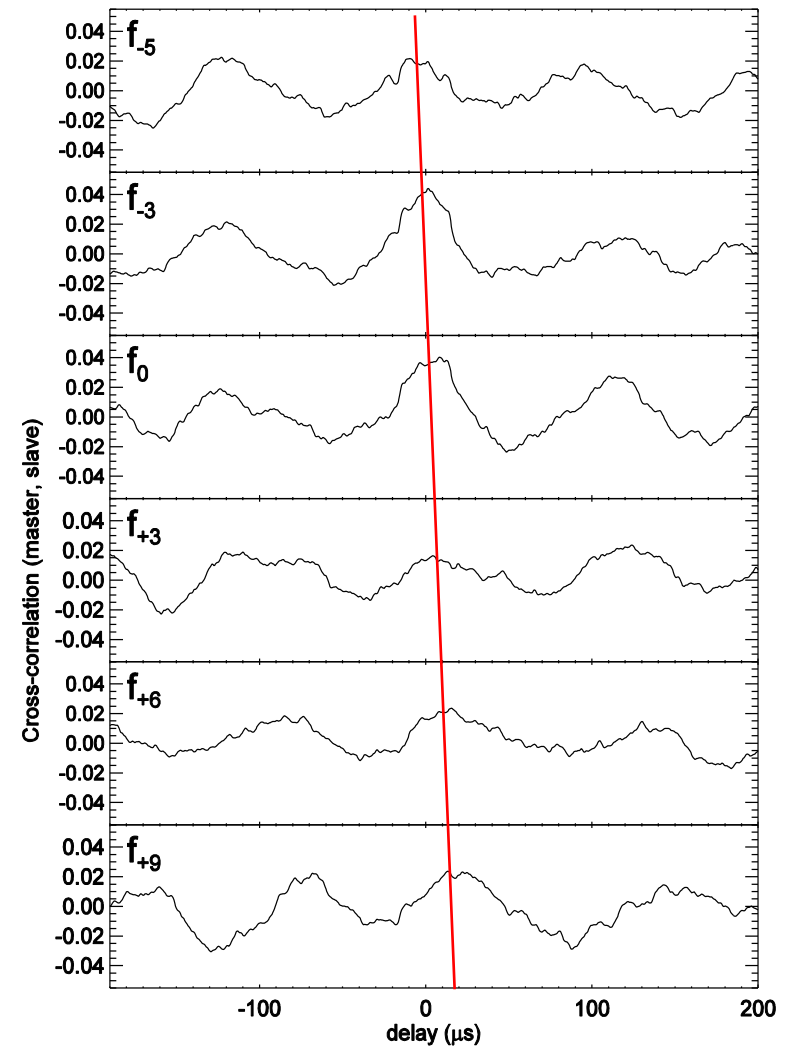
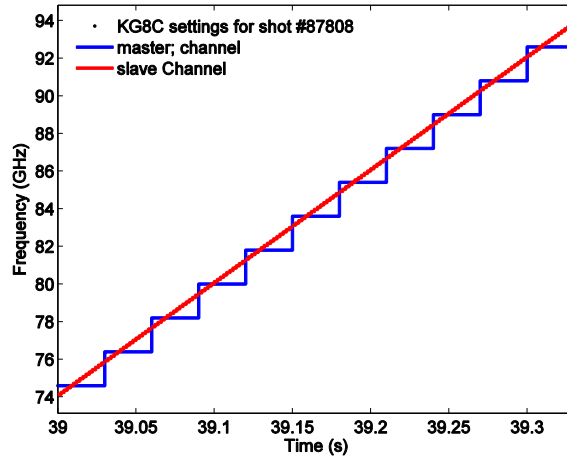


- Mean v_{\perp} also determined from Doppler reflectometer
- GAM located in the E_r well with amplitude up to 50% of mean v_{\perp}
- Similar to AUG/DIII-D observations
- Amplitude and width of the GAM region and E_r well depends on the density

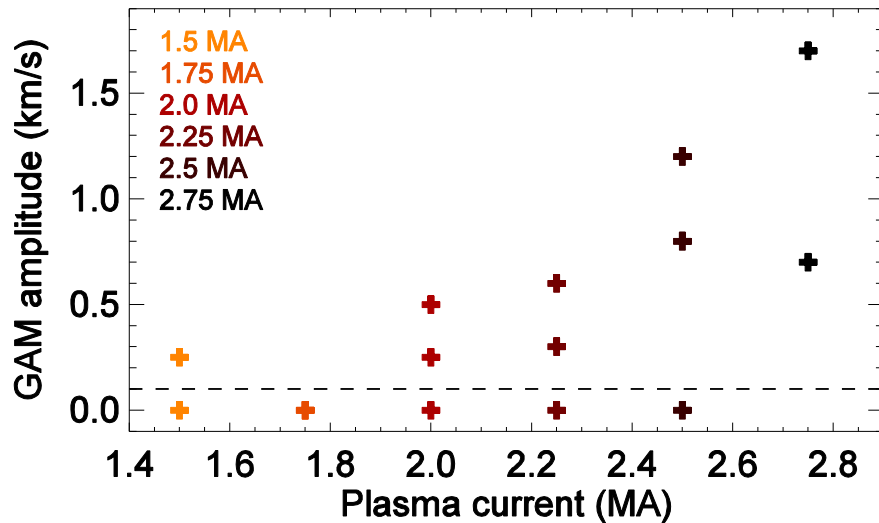
GAM radial structure



- ❑ Correlation between master and slave signals used to determine radial structure
- ❑ Delay $\sim 20 \mu\text{s}$ across slave freq: $\sim T_{\text{GAM}}/5$
- ❑ Slave sweep covers $\sim 0.5 \text{ cm}$ radially $\Rightarrow \lambda_r \approx 2.5 \text{ cm}$
- ❑ Radial wavelength $> \sim$ dimension of GAM existence region



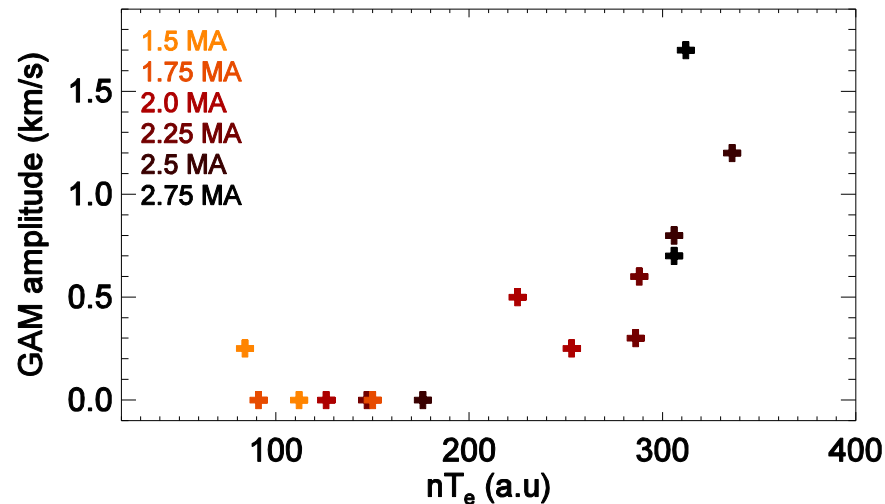
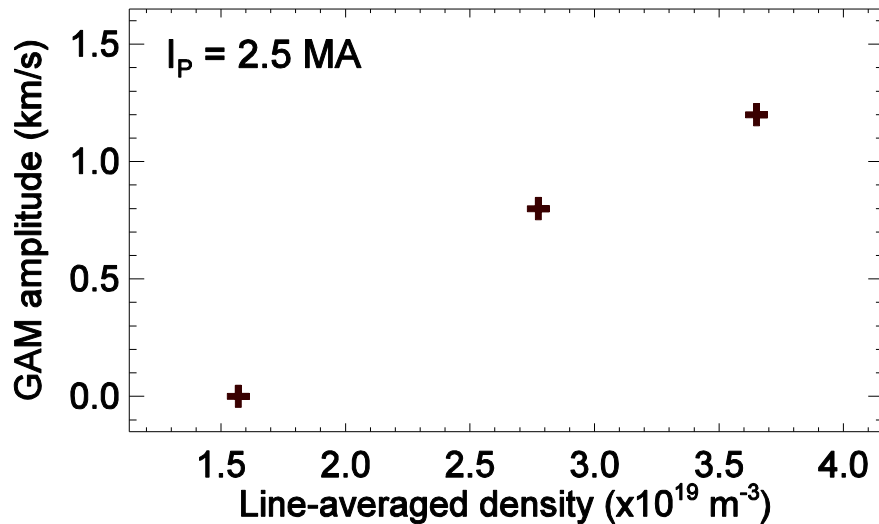
GAM dependence on plasma current and density



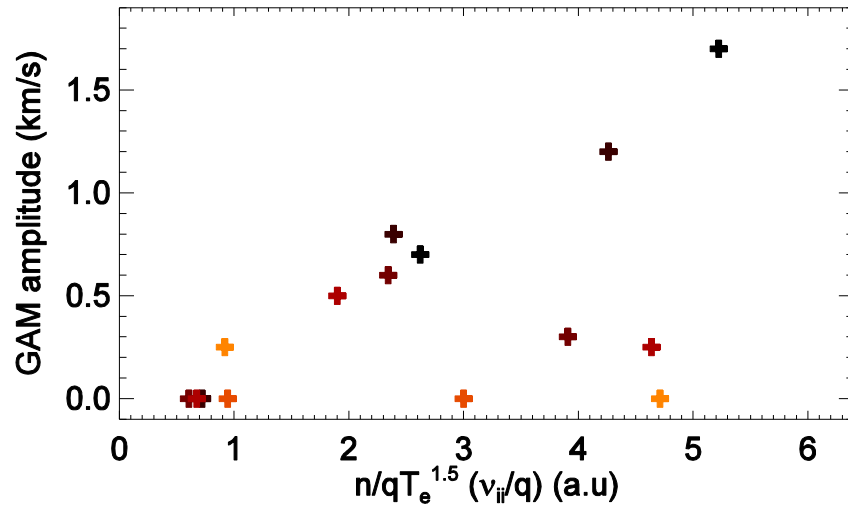
Dataset (H): $1.5 < I_p < 2.75$ MA,
 $1.5 < n < 4 \times 10^{19} \text{ m}^{-3}$

GAM amplitude

- Increases with I_p and n (at high I_p)
- Depends on the edge pressure



GAM damping rate

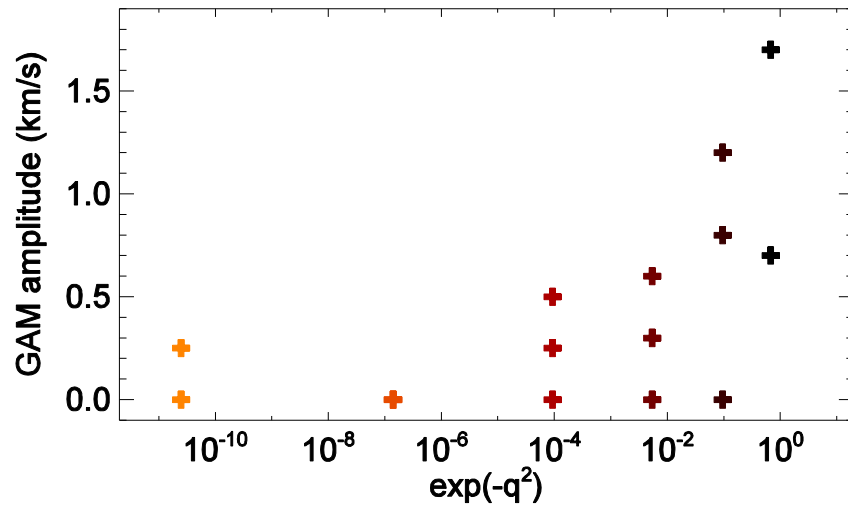


A_{GAM} determined by its drive and damping mechanisms

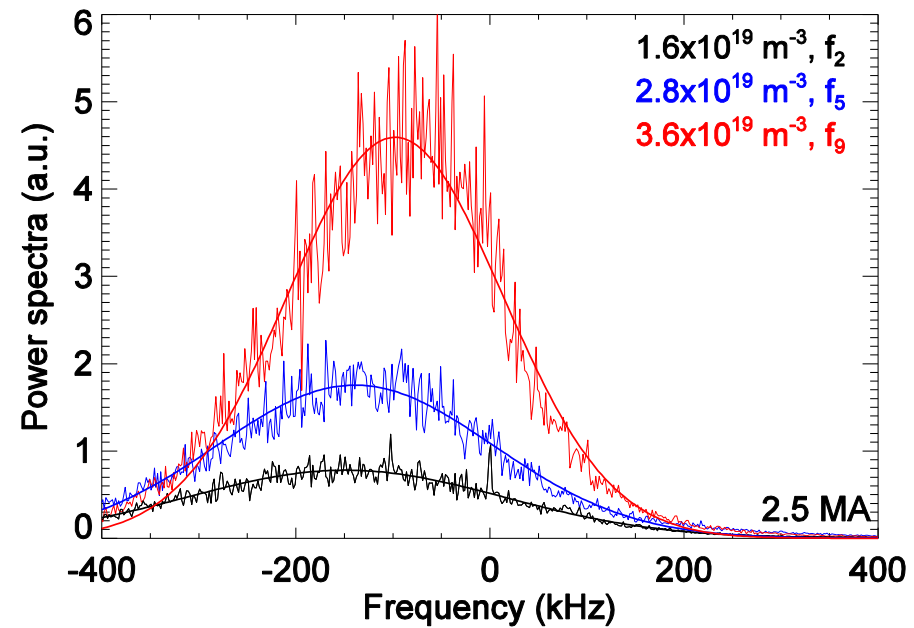
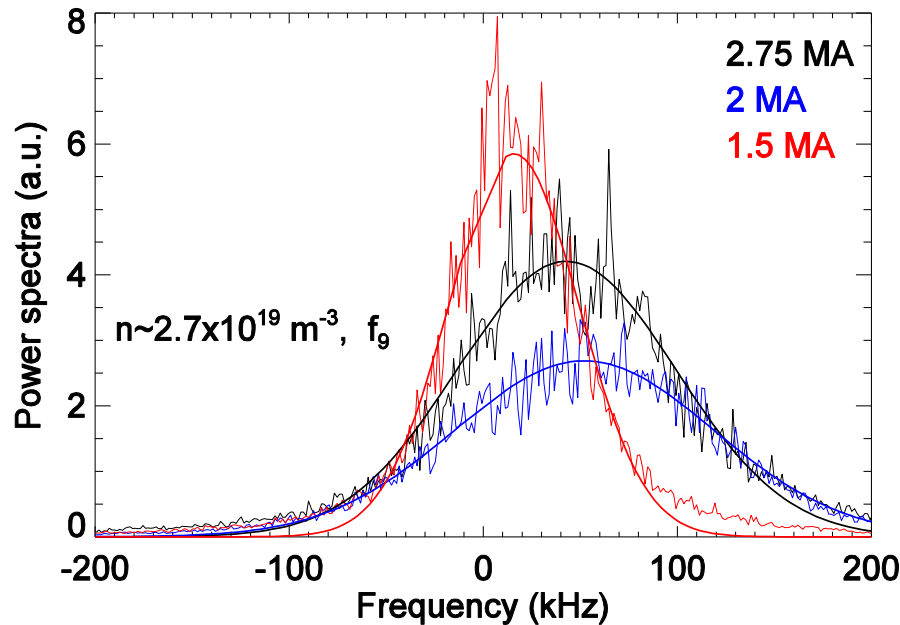
□ Collisionless damping rate: $\gamma \sim \exp(-q^2)$
Collisional damping rate: $\gamma \sim v_{ii}/q$

□ A_{GAM} increases with both $\exp(-q^2)$ and v_{ii}/q in disagreement with theory

□ GAM amplitude not determined by its damping rate

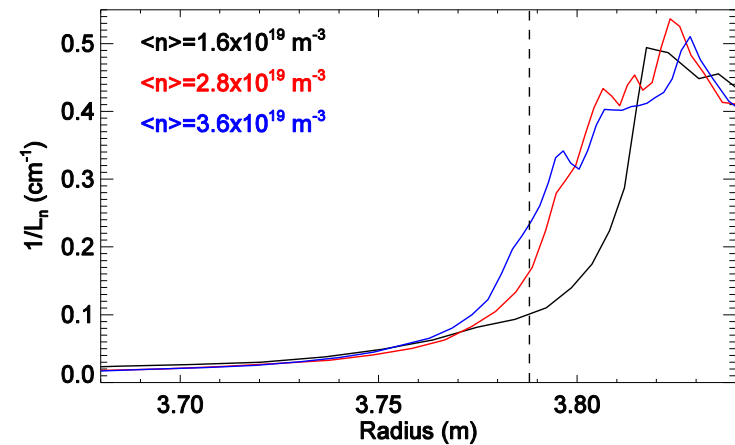
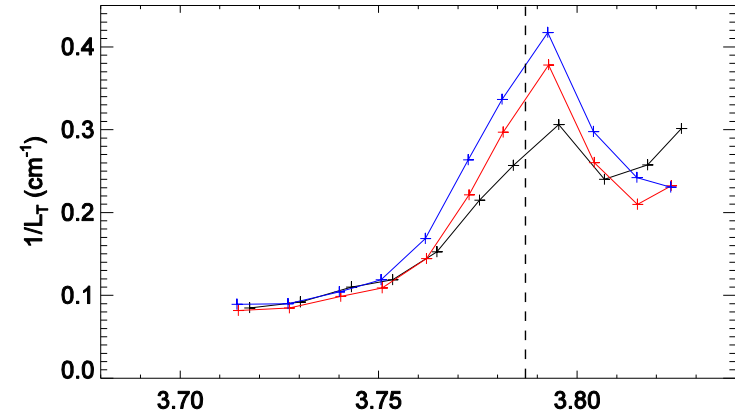
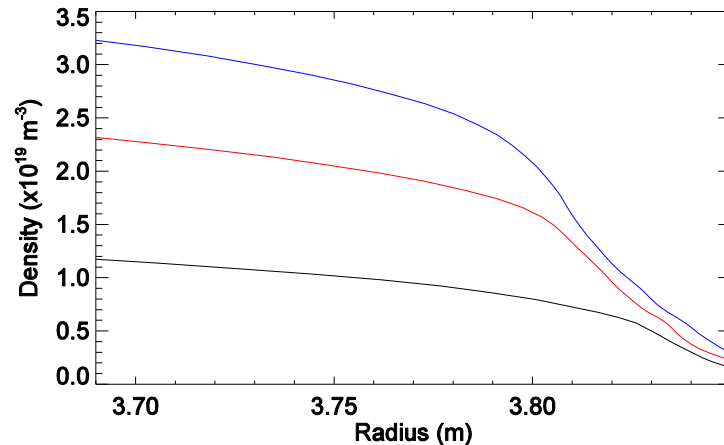
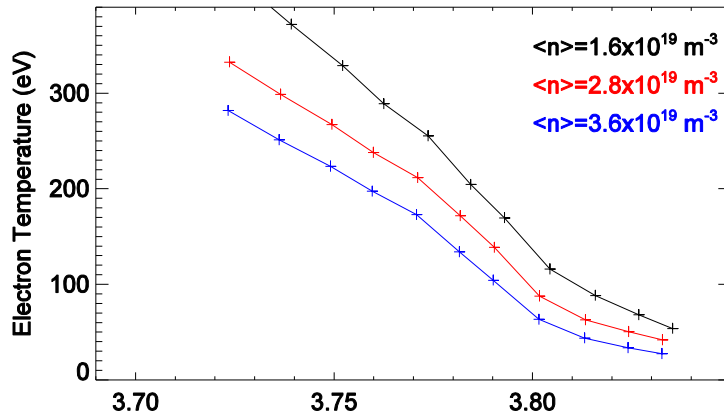


GAM drive: turbulence – density fluctuations



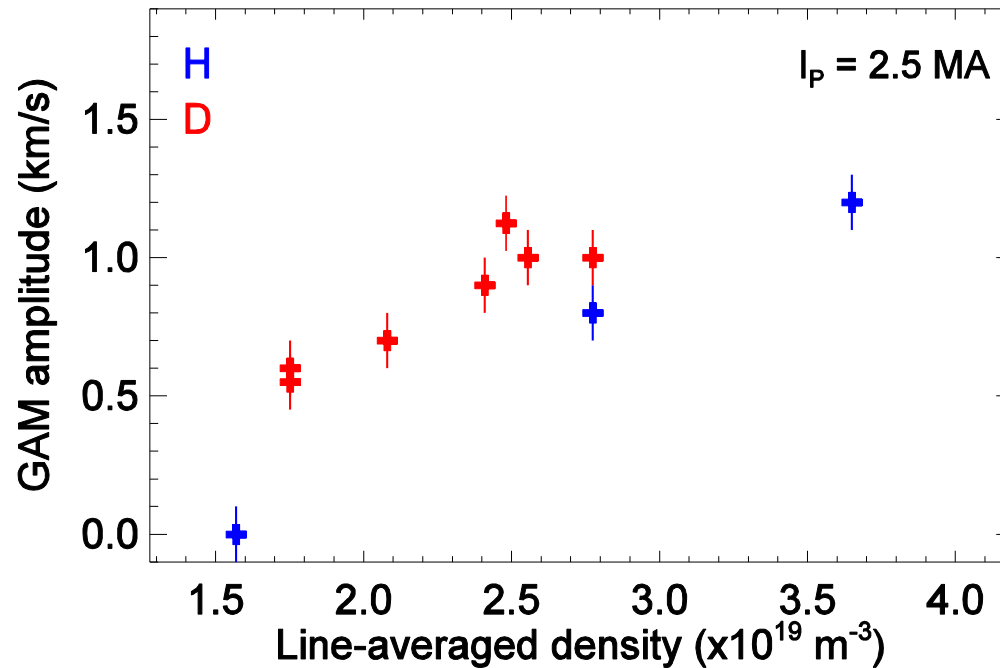
- Density fluctuations \propto amplitude backscattered signal
- Fluctuations increase with $\langle n \rangle$ and I_p , similarly to the GAM amplitude dependence

GAM drive: turbulence – scale length



- **Density dependence:** T_e and n inverse scale length increase with $\langle n \rangle$ in the GAM region, indicating an increase in the turbulence drive (in agreement with the larger GAM amplitude)

Isotopic effect



- ❑ GAM amplitude is apparently larger for D plasmas but complex dependence on plasma parameters prevents a more definitive conclusion
- ❑ Consistent with results on TEXTOR, FT-2, TJ-II and ISTTOK supporting a multi-scale mechanism to explain the isotope effect



- ❑ GAMs investigated JET edge plasma using Doppler reflectometry
- ❑ GAM clearly observed with RCR and magnetics, located edge density gradient region in a narrow layer with frequency radially constant
- ❑ GAM amplitude up to 50% of the mean E_r
- ❑ A_{GAM} in agreement with turbulence drive but not with theoretical damping rate
- ❑ A_{GAM} is apparently larger for D plasmas but complex dependence on plasma parameters prevents a more definitive conclusion

Outlook

- ❑ Explore the importance of GAMs when approaching L-H transition