



# The Plasma $E \times B$ Staircase: Turbulence Self-Regulation through Spontaneous Flow Patterning

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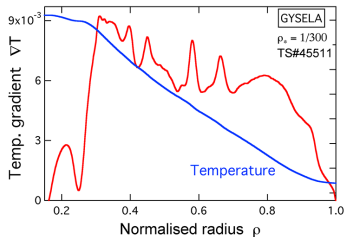
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<sup>5</sup>WCI Center for Fusion Theory, NFRI, Daejeon, Korea

- Ackn.:
- Festival de théorie, Aix-en-Provence 2011 & 2013
  - KITP, Santa Barbara, 2014 & 2015

- T corrugations
- mean  $\mathbf{E} \times \mathbf{B}$  shear
- flow patterning  $\implies$  depart neoclass.
- set of quasi-regularly spaced weak transport barriers

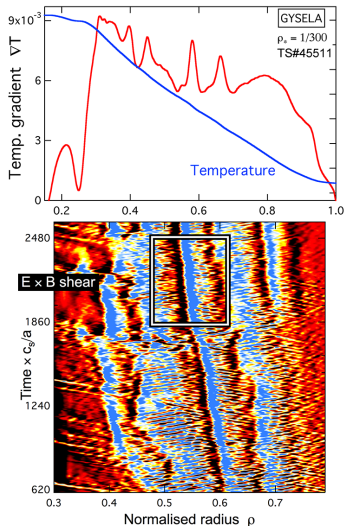
- 
- robust with  $\rho_*$   $\implies$  beneficial for conf.
  - robust with  $\nu_*$
  - observed so far near ITG threshold  
large machines :  $\frac{\text{intern. energy}}{\text{ext. heating}}$   $\nearrow$



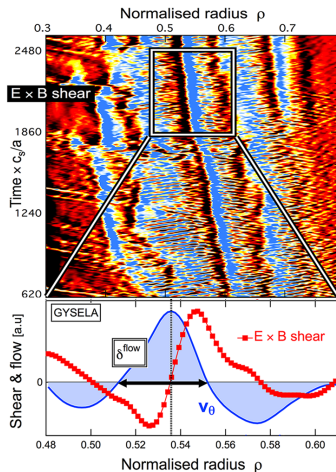
**understand turb. self-organisation?**

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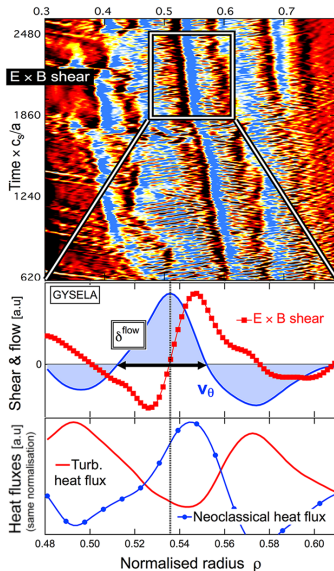
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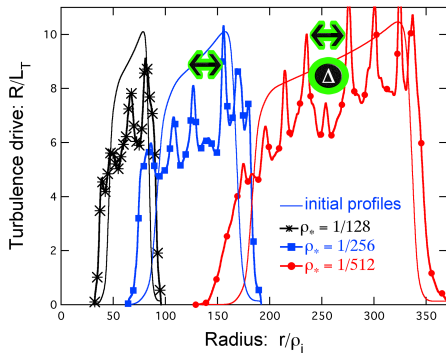
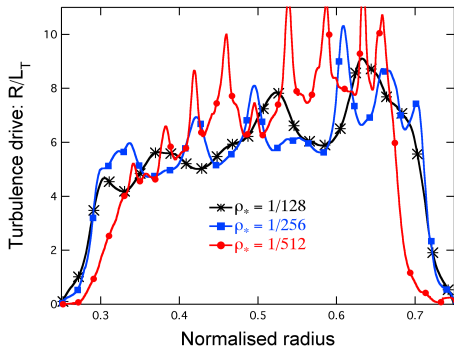


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- \_\_\_\_\_ [Dif-Pradalier PRE 10 & PRL 15]
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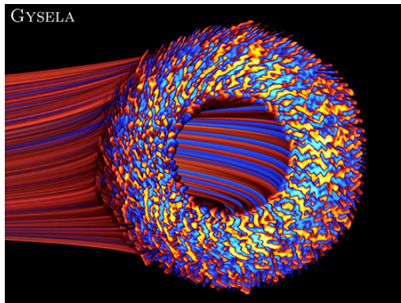


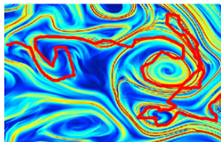
- ▶ # shear layers ↗ as machine size ↗ ⇒ beneficial for confinement
- ↳ spacing saturates past  $\rho_* \sim 1/300$
- ↳ expect 3-5 in a medium-size [ $\rho_* \sim 1/300$ ] tokamak

underlying physics for this mesoscale ?

# Conflicting ways to understanding transport

➔ three generic trends for turb. self-organisation



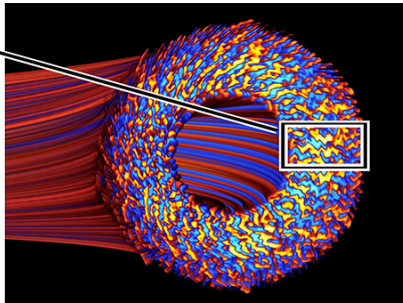


**Random walk,  
diffusion**

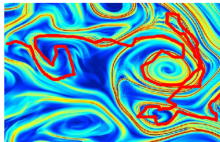
→ ~ 40 years of  
plasma modeling

- density/temp. gradients
- bath of turbulent eddies
- magnetic geometry

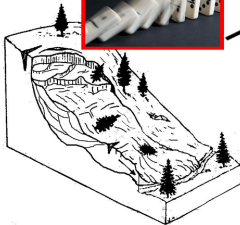
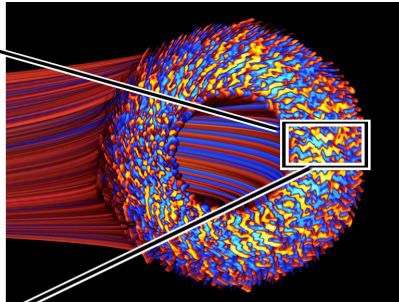
magnetic geom. → "pinned" turbulent eddies →  
random walk transport → **diffusion**  $Q = -n\chi\nabla T$





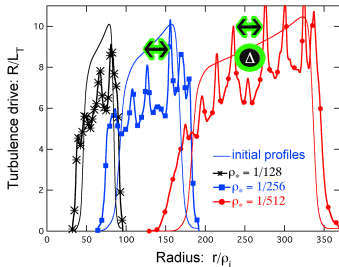


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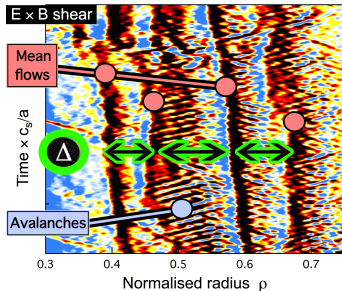
**Avalanches (detrimental)**

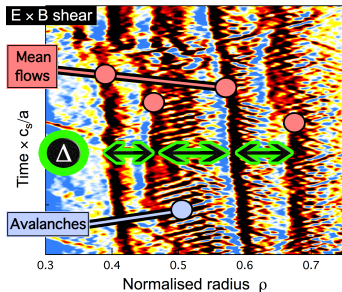
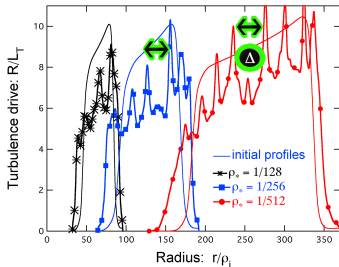
- non-local & non-diffusive
  - self-consistent **profile dynamics**
- hindered in “usual” turb. modeling  
[flux-tube, local or gradient-driven]
- need **flux-driven** approaches



$$Q = -n\chi(r)\nabla T \Rightarrow Q = - \int \mathcal{K}(r, r') \nabla T(r') dr'$$

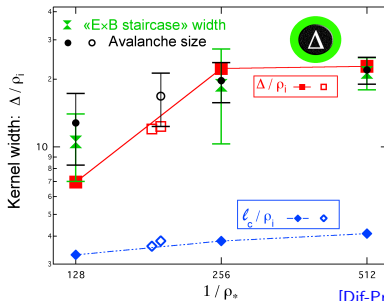
$$\Rightarrow \mathcal{K}(r, r') = \frac{S}{\pi} \frac{\Delta/2}{(\Delta/2)^2 + (r-r')^2}$$





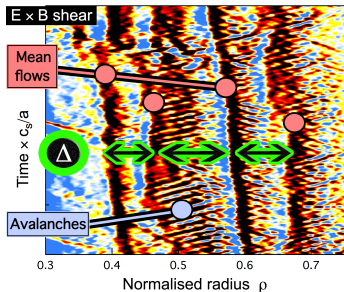
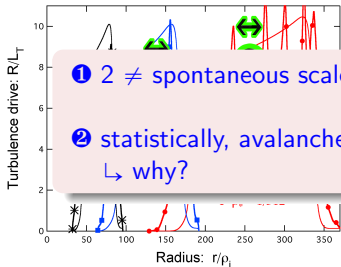
$$Q = -n\chi(r)\nabla T \Rightarrow Q = - \int \mathcal{K}(r, r') \nabla T(r') dr'$$

$$\Rightarrow \mathcal{K}(r, r') = \frac{5}{\pi} \frac{\Delta/2}{(\Delta/2)^2 + (r-r')^2}$$

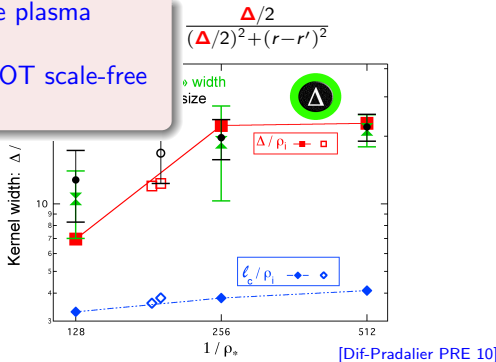


at scales  $\sim \ell_c \equiv$  diffusive-like transport

at scales  $\in [\ell_c; \Delta] \equiv$  scale-invariant, avalanche-mediated, non-local, non-diffusive transp.



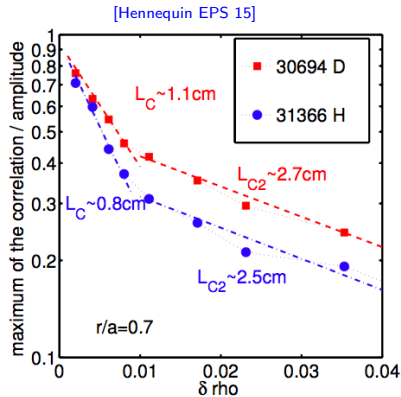
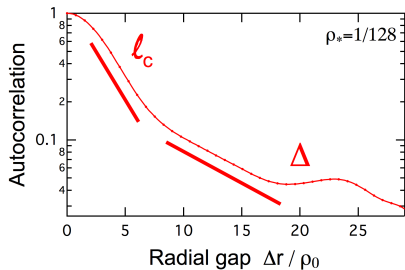
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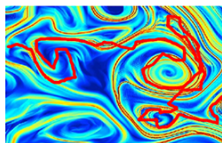
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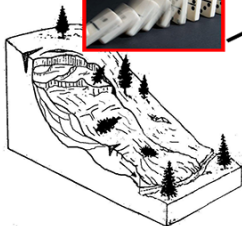
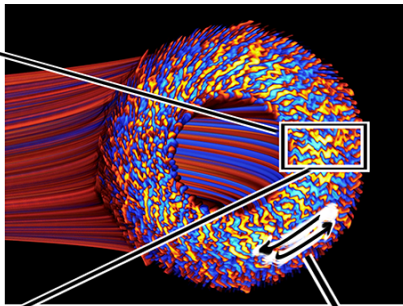
① — Two spontaneous scales for the plasma:  
 $l_c \equiv$  diffusive step &  $\Delta \equiv$  avalanche extension



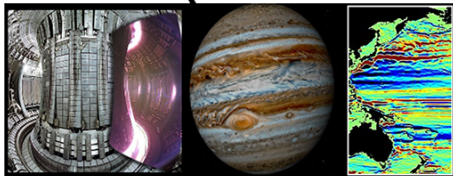
② — Why are avalanches not scale-invariant?  
 → three generic trends for turb. self-organisation



Random walk,  
diffusion



**Avalanches (detrimental)**



**Zonal/Mean flows (beneficial)**

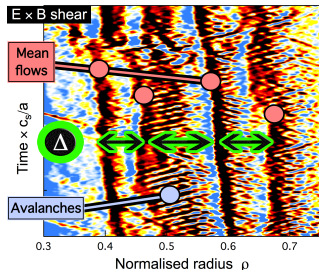
**The conundrum:** two spontaneous patterning trends. . . yet mutually-exclusive

- ▶ avalanches: correlated over-turnings  $\Rightarrow$  extended, intermittent transport
- ▶ zonal flows: limit the turbulent mixing, regulates transport via shearing

## How can they coexist?

- 1 they do not: both exist in different regions of parameter space?
- 2 separation in time: both exist alternately?
- 3 separation in space  $\Rightarrow$  the “ $E \times B$  staircase”

[Dif-Pradalier PRE 10; Kosuga PRL 13; Dif-Pradalier PRL 15]



the  $E \times B$   
staircase

- ▶ ZFs concentrate into thin layers & endure  $\Rightarrow$  mean flows (MFs)
- ▶ MFs organise into a lattice of quasi-regularly spaced transp. barriers
- ▶ avalanches propagate in-between

## A LOT TO UNDERSTAND STILL...

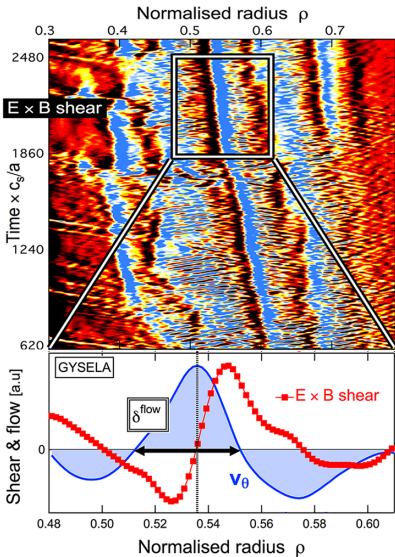
➔ WHAT PREDICTIONS? [a large dataset from GYSELA]

Goal: provide figures for the experimental characterisation of the structure

- ☛ radial resolution  $\delta^{\text{flow}} \sim 10\rho_i \sim 1\text{cm}$  [~ 1.5cm, Fujisawa PRL 04]
- ☛ radial extension: 2 – 3 layers  $\sim 60 - 100\rho_i$
- ☛ long-lived mean flows  $\geq 10\text{ms}$ ... [~ 1.5ms, Fujisawa PRL 04]
- ☛ ... yet fast meandering  $\sim 1\text{ms}$

needed: profiles, fast & high-res. [Hornung, following]



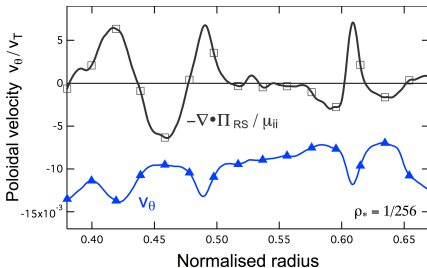


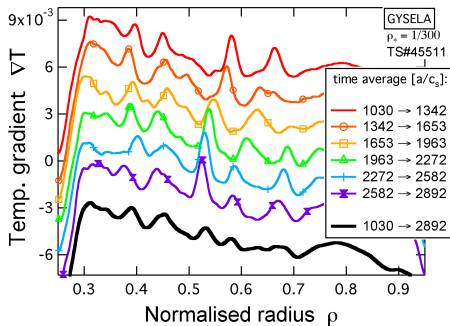
$\delta^{\text{flow}} \sim 10\rho_i \sim 1\text{cm}$  [Dif-Pradalier, PRL15]

see G. Hornung

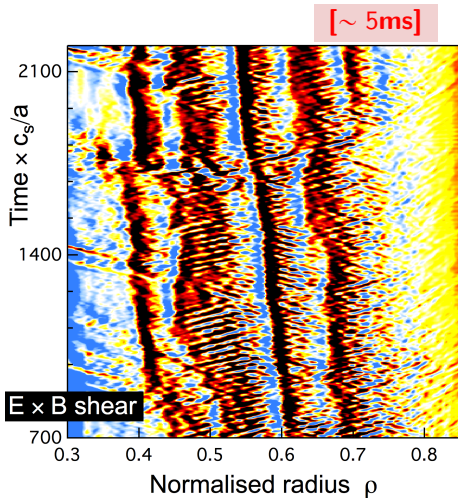
local deviation from neoclassical prediction

[Bell PPCF 98, Crombé PRL 05, Dif-Pradalier PRL 09]





no fast-sweep  $\Rightarrow$  no exp. observation



- ▶ spacing: regular, staircase: dynamical
- ▶ not pinned to a precise location
  - ↳ meanders, disappears, reforms
  - ↳ not linked to low-order  $q$  rationals

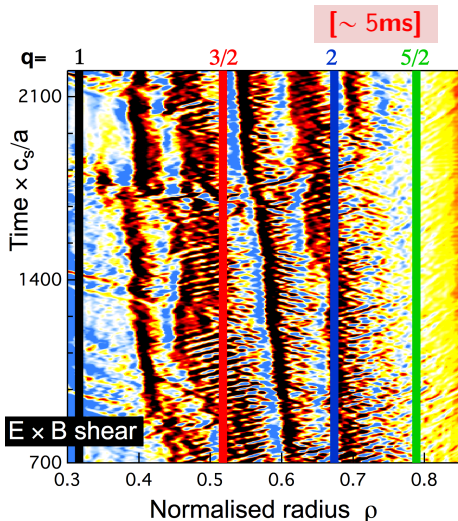
see G. Hornung

in a nutshell: 3 – 5 steps in medium size tokamak

resolve:  $\sim$ ms range &  $\sim$ cm range  
over  $\sim 10$ 's cm

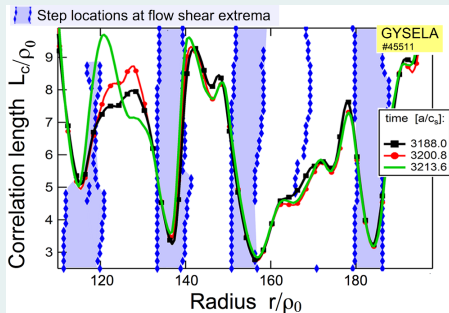
“straightforward” is uneasy

↳ “higher order”  $\equiv$  **correl. meas.**

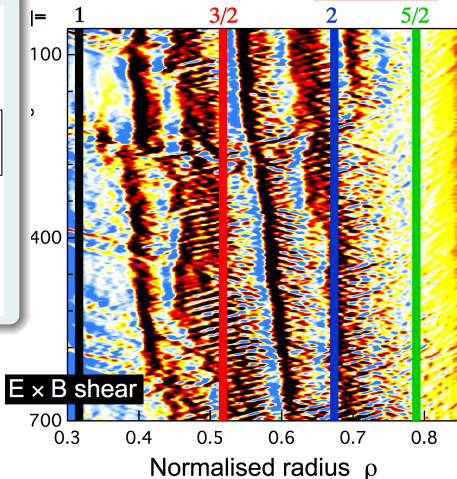


[experimentally...actually good news]

## Radial correlation length of turb. fluct.



[~ 5ms]



in tokamak  
 resolve: ~ms range & ~cm range  
 over ~ 10's cm

"straightforward" is uneasy  
 ↳ "higher order"  $\equiv$  **correl. meas.**

- ▶ **Turb.-borne long-lived & localised polo. flows  $\neq$  neoclass.** [Dif-Pradalier PRL 09]
- ▶ **Identification of the  $E \times B$  staircase** [Dif-Pradalier PRE 10]
  - nonlocal transport & stochastic avalanches
  - mesoscales are key  $\Rightarrow$  2 typical correlation lengths [Hennequin, this conference]
- ▶ **Model for staircase emergence** [Kosuga PRL 13; PoP 14]
  - clustering instability through a time-delay between flux & gradient
  - staircase with mesoscale step
- ▶ **Predictions v.s. experiments** [Dif-Pradalier PRL 15]
  - no link to low-order  $q$  rationals
  - radial turb. correlations  $\Rightarrow$  observation on ToreSupra & predictions
- ▶ **Experimental characterisation of the  $E \times B$  staircase** [Hornung, in preparation]

**NB:** may imply important changes in common (diffusive-based) turb. models

[Nakata NF 13; Villard PPCF 13; Norscini, Varenna 14; Cartier-Michaud, Varenna 14; Imadera, IAEA 15]

# Additional material

(i) Gyrokinetic equation for the gyrocenters  $\bar{f}(\mathbf{x}_g, v_{\parallel,g}, \mu)$

$$\frac{\partial \bar{f}}{\partial t} + (\mathbf{v}_{E \times B} + \mathbf{v}_D) \cdot \nabla \bar{f} + v_{\parallel} \nabla_{\parallel} \bar{f} + \frac{dv_{\parallel}}{dt} \partial_{v_{\parallel}} \bar{f} = \mathcal{C}(f) + \mathcal{S}(f)$$

$$\mathbf{v}_{E \times B} = \frac{1}{B^*} \mathbf{b} \times \nabla \bar{\phi}$$

$$m \frac{dv_{\parallel}}{dt} = - (\mu \nabla B + e \nabla \bar{\phi}) \frac{\mathbf{B}^*}{B^*}$$

$$\mathbf{v}_D = \frac{mv_{\parallel}^2 + \mu B}{eB^*} \mathbf{b} \times \frac{\nabla B}{B}$$

$$\mathbf{B}^* = \mathbf{B} + \frac{mv_{\parallel}}{e} \nabla \times \mathbf{b}$$

(ii) Quasi-neutrality equation for the particles

$$\frac{\delta n_e}{n_{eq}} = \frac{\delta n_i}{n_{eq}}$$

$$\frac{e}{T_e(r)} (\phi - \langle \phi \rangle) = \frac{2\pi B}{mn_0} \iint d\mu dv_{\parallel} J_0 (f - f_{init}) - \frac{1}{n_0(r)} \nabla_{\perp} \cdot \left[ \frac{n_0}{B \omega_c} \nabla_{\perp} \phi \right]$$

## The GYSELA framework in a nutshell

(i) Gyrokinetic

$$\frac{\partial \bar{f}}{\partial t} + (\mathbf{v} \mathbf{E} \times \mathbf{B})$$

$$\mathbf{v} \mathbf{E} \times \mathbf{B} = \frac{1}{B_{\parallel}^*} \mathbf{b}$$

$$\mathbf{v}_D = \frac{mv_{\parallel}^2 + \mu B}{eB_{\parallel}^*}$$

(ii) Quasi-neu

$$\frac{e}{T_e(r)} (\phi - \langle \phi \rangle)$$

- ▶ full-f  $\equiv$  **no scale separation** in space & time

[Grandgirard '06]

- ▶ flux-driven  $\equiv$  sources of heat, momentum & vorticity

- open system, known sources
- **self-consistent** force balance  $E_r$ ,  $\nabla p$ ,  $v_{\theta}$  &  $v_{\varphi}$
- **stiffness** of the mean profiles
- **locality & non-locality**  $\equiv$  on equal footing

[... , Dif-Pradalier '10, Ghendrih '12, Gurcan '13, Kosuga '13]

- ▶ **adiabatic electrons & electrostatic**

self-organising framework  $\Rightarrow$  what transport to be expected?



## ① what we do know

### Globally-organised pattern spontaneous & weak transport barriers

[Dif-Pradalier PRL 15]

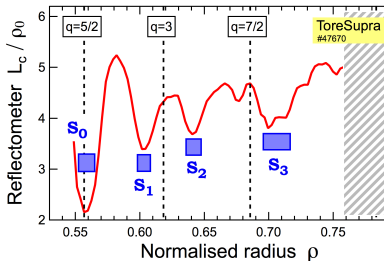
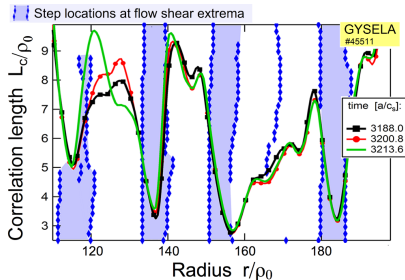
Visible on	$\gamma_{\mathbf{E} \times \mathbf{B}}, \nabla p, L_c, v_\theta, v_\parallel, \langle \tilde{v}_{Er} \tilde{v}_{E\theta} \rangle, \langle \tilde{v}_{Er} \tilde{v}_\parallel \rangle$
$\rho_\star = \rho_i/a$	$1/75 \rightarrow 1/512$
Step spacing	Outer scale of avalanche distribution constant [ $\sim 20 - 30\rho_0$ ] for $\rho_\star \leq 1/300$ [1]
No. of steps	$1 \rightarrow 2[\rho_\star = 1/75]; 3 \rightarrow 5[\rho_\star = 1/300]; 5 \rightarrow 7[\rho_\star = 1/512]$
Flow thickness	$\delta^{\text{flow}} \sim 10\rho_0$
Collisionality $\nu_\star$	$0.001 \rightarrow 1$
$\langle\langle R/L_T \rangle\rangle$	$4 \rightarrow 8$
$\langle\langle R/L_n \rangle\rangle$	$1 \rightarrow 4$
$\eta = L_n/L_T$	$2 \rightarrow 8$
Meandering	Stay at $\sim$ constant drive, follow $\nabla p(t)$
Strength	$ \gamma_{\mathbf{E} \times \mathbf{B}}  \sim$ constant for $\rho_\star \geq 1/300$
Resonant $q$	No correlation with low-order rationals

## ② what is still [somewhat] unknown

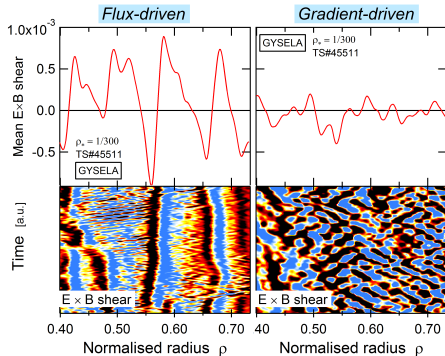
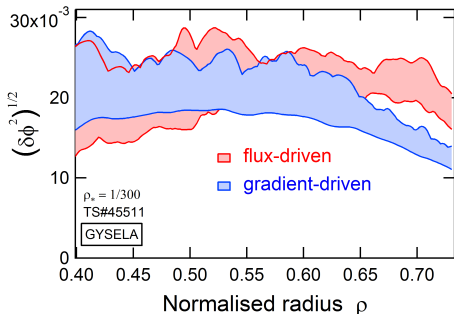
- ▶ mechanism whereby large-scale org. occurs [Kosuga PoP 14]
- ▶ robustness & domain of existence [Hornung, in prep.]
- ▶ impact on macro-transport properties & modeling

## ④ why (we should) care?

- because quite **robust in flux-driven** modeling [Dif-Pradalier PRE 10; Nakata NF 13; Villard PPCF 13; Norscini Varenna 14; Cartier-Michaud Varenna 14; Imadera, this conf.]
- because near-critical self-organisation may be key to **next-gen. tokamaks**
- GFD staircase “**the vein and arteries of the weather system**” [McIntyre, Festival 09]
- **because it exists!** prediction leads to exp. discovery [Dif-Pradalier PRL 15; Hornung, in prep.]



- may point towards important **changes in current turbulence models**



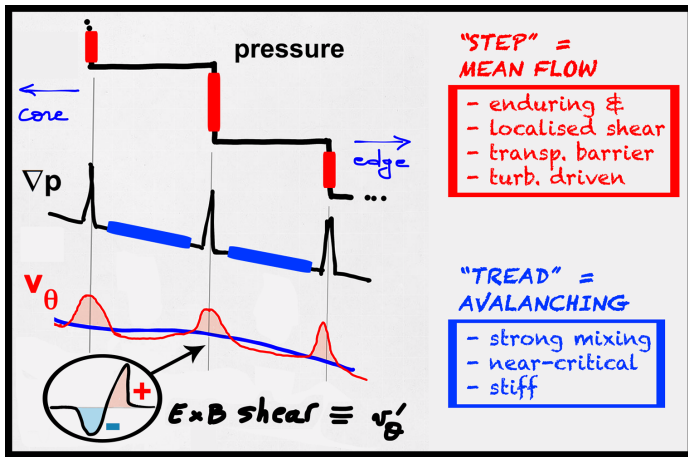
- ▶ similar turbulence level
- ▶ mean shear pattern is lost : isotropisation of the fluctuations.  
↳ ongoing : global impact on transport

## what we do know

Where staircases are observed in GYSELA

[Dif-Pradalier PRL 15]

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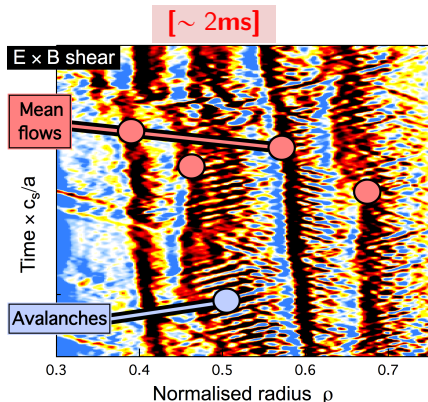
- ▶ avalanches interspersed between mean zonal flow layers
- ▶ shear flow layers associated to mean profile corrugations [ $\equiv \nabla p$  or  $\nabla f_s$ ]

## global & dynamic pattern

- ➔ impacts all turbulent fields
- ▶  $\mathbf{E} \times \mathbf{B}$  shear  $v'_{E,\theta} = r\partial_r(E_r/rB)$
- ▶ Heat flux;  $v_{E,\theta}$ ;  $\langle v_{Er}v_{E\theta} \rangle, \dots$

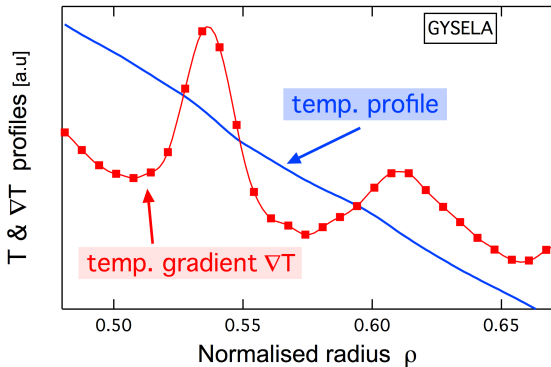
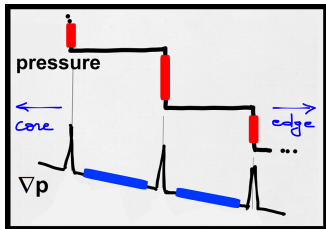
predicted figures:

- radial res.  $\delta^{\text{flow}} \sim 10\rho_i \sim 1\text{cm}$
- inter-shear layer: mesoscale  $\sim 30\rho_i$
- long-lived mean flows  $\geq 10\text{ms} \dots$
- ... yet fast meandering  $\sim 1\text{ms}$
- no link to low-order  $q$  rationals



## Legitimate concerns:

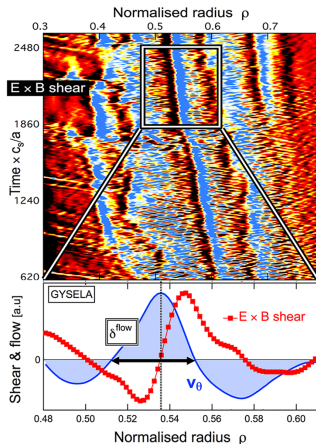
- ➔ does it really exist? [Dif-Pradalier '10; Nakata '13; Villard '13; Norscini '14; Cartier-Michaud '14; Imadera, this conf.]
- ➔ can we observe it? why never observed? [Dif-Pradalier PRL 15; Hornung, in prep.]



- ▶ unambiguous on the mean gradients
- ▶ difficult to characterise on the profile itself
  - ➡ characterisation through **fluctuation measurements**

## 1<sup>st</sup> choice: directly measure the mean flows

- ▶ regularly-spaced poloidal  $\mathbf{E} \times \mathbf{B}$  shear flow layers
  - ➡ local deviation from neoclassical prediction  
[Bell PPCF 98, Crombé PRL 05, Dif-Pradalier PRL 09]
  - ➡ experimentally: a difficult measurement  
[C impurities, Doppler reflectometry, HIBP, ...]
- radial resolution  $\delta^{\text{flow}} \sim 10\rho_i \sim 1\text{cm}$   
[~ 1.5cm, Fujisawa PRL 04]
- radial extension: 2 – 3 layers  $\sim 60 - 100\rho_i$
- long-lived mean flows  $\geq 10\text{ms} \dots$   
[~ 1.5ms, Fujisawa PRL 04]
- ... yet fast meandering  $\sim 1\text{ms}$



need **fast, high-res.** flow measur. ➡ **direct charact.** through flows **impractical**



## 2<sup>nd</sup> choice: characterise avalanching/transport

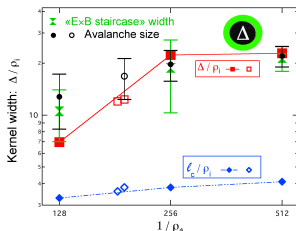
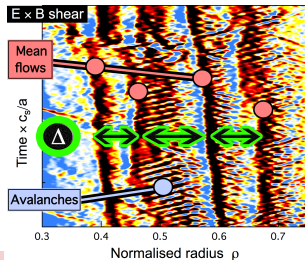
- ▶  $\rho_*$  dependence  $\implies$  break gyroBohm scaling?
  - **naive:** avalanching  $\equiv$  scale-independent  $\rightarrow$  gBohm breaking  $\rightarrow \rho_*$  dependence
  - **computation:** [Dif-Pradalier PRE 10]
    - staircase  $\implies$  mesoscale avalanching

$$Q = -\int \mathcal{K}(r, r') \nabla T(r') dr' ; \mathcal{K} = \frac{1}{\pi} \frac{\Delta/2}{(\Delta/2)^2 + (r-r')^2}$$

$\implies$  beneficial for conf.: spontaneous & mesoscale permeable/weak array of transp. barriers

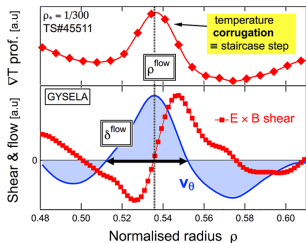
- ▶ direct obs. on fluctuations [Politzer PRL 00, Hornung PPCF 13]
  - high-res. temporal data
  - phase fluctuation  $\neq$  dens. fluctuations
  - **localisation:** poloidal & radial

$\implies$  experimentally: a difficult measurement



## 3<sup>rd</sup> choice: through profile corrugations

Mean Flow  $\leftrightarrow$  staircase step  $\leftrightarrow$  mean prof. corrugation



clear & systematic:  
medium-large machines

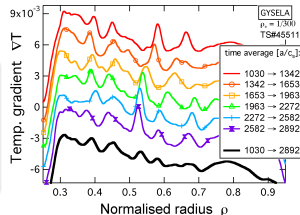
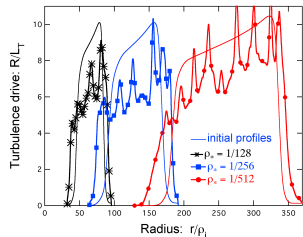
- locally, large  $\nabla p$
- $\rho^{\text{flow}} \sim \text{constant}$
- $\Delta \sim \text{constant}$
- near critical & stiff

[avalanches, SOC, ...]

in a nutshell: 3 – 5 steps in medium size tokamak  
resolve:  $\sim \text{ms}$  range &  $\sim \text{cm}$  range over  $\sim 10\text{s}$  cm

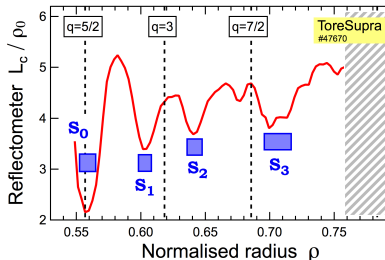
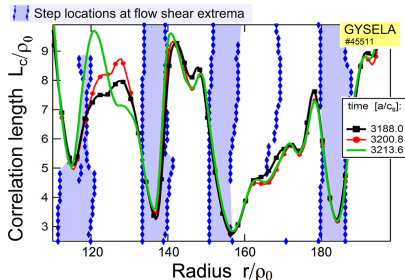
“straightforward” seems impractical

$\hookrightarrow$  “higher order”  $\equiv$  correlation measurements



- ▶ GYSELA: local minima turb. correlation lengths track staircase steps

[Ghendrih EPJD 14]

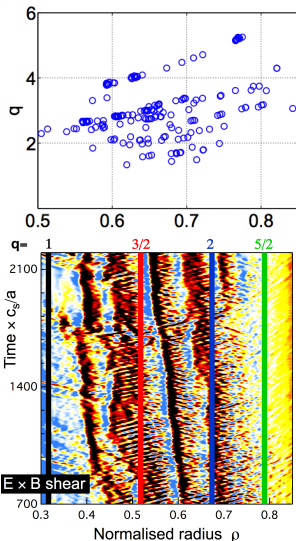
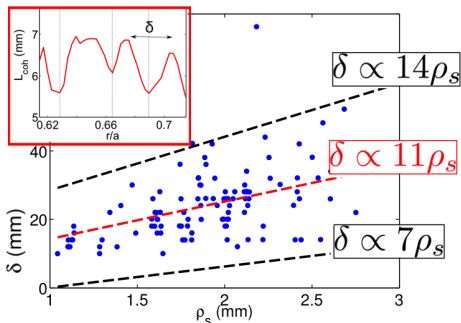


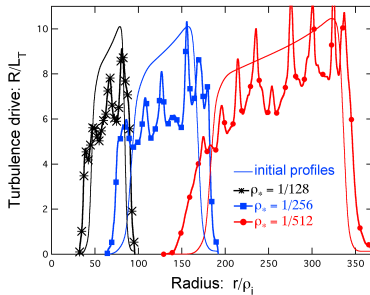
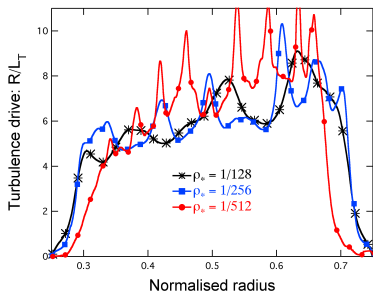
- ▶ Experiment:  $L_c \equiv$  proxy for fluct. size [Dif-Pradalier PRL 15 & Hornung, in prep.]



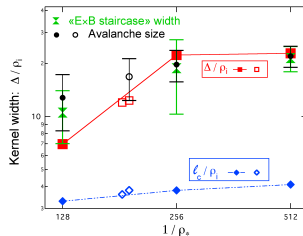
- $L_c$  minima quasi-regularly spaced radially
- reproducible: not random & robust w.r.t. definition of  $L_c$
- tilt consistent with flow shear around minima

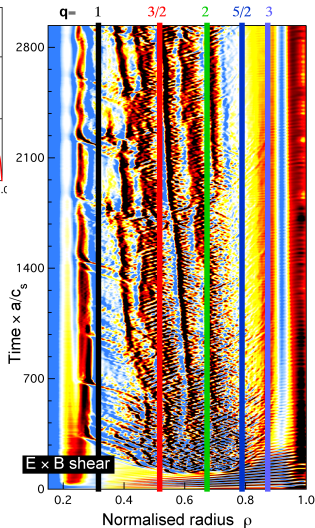
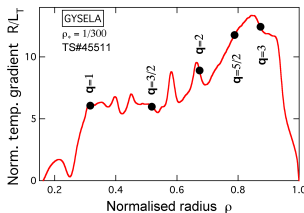
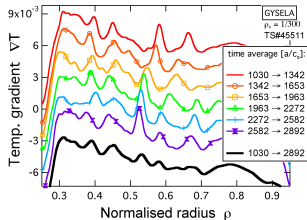
- ▶ flow width  $\delta \sim 11\rho_i$ ; consistent with GYSELA obs. & ZF measurements  
[Fujisawa PRL 04]
- ▶ turbulence-borne  $\Rightarrow$  not MHD  
[Dif-Pradalier PRL 15 & Hornung, in prep.]





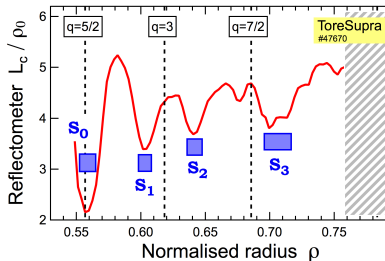
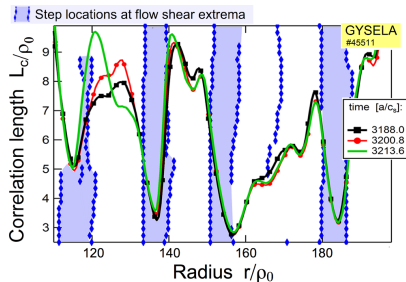
- ▶ how many shear layers?  $\Rightarrow$  depends on  $\rho_*$
- ▶ # shear layers  $\nearrow$  as machine size  $\nearrow$ 
  - $\hookrightarrow$  spacing saturates past  $\rho_* \sim 1/300$
  - $\hookrightarrow$  expect 3-5 in a medium-size tokamak



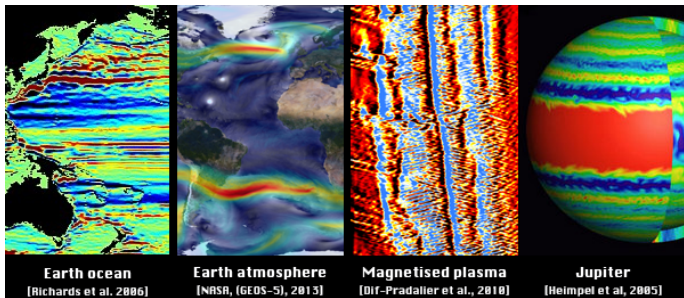


- ▶ spacing is regular, yet staircase is ...
- ▶ ...dynamical : not pinned to a precise location
  - ↳ meander, disappear, reform
  - ↳ not linked to low-order  $q$  rationals

[Dif-Pradaliér, PRL 15]



- ▶ structure difficult observe through direct profile measurements
- ▶ possible through turbulent fluctuation measurements [Hornung PPCF 13]
  - ↳ radial profiles of turb. correlation lengths [Ghendrih EPJD 14, Dif-Pradalier PRL 15]



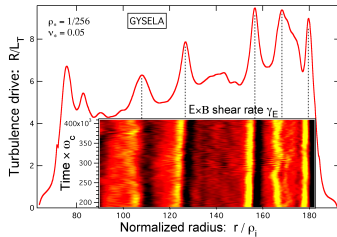
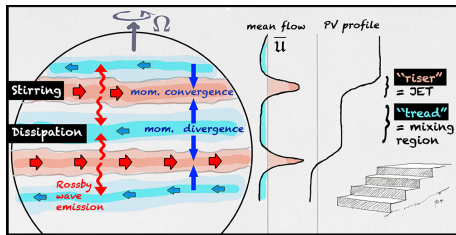
**Atmospheric fluids:** potential vorticity staircase, zonal flow structure from Rossby wave dynamics in quasigeostrophic turbulence ;

**Magnetised plasma:**  $\mathbf{E} \times \mathbf{B}$  staircase, quasi-regular pattern of zonal flow layers co-located with profile corrugations interspersed by avalanching ;

**Oceans:** thermohaline staircase, alternately sharpened and flattened convection gradients driven by competing heat & salt gradients



Why the name of “staircase”? analogy: “PV staircases” & atmospheric jets



[Festival de Théorie 2009 & 2011 ; Kavli Institute (KITP) 2014 & 2015]

Legitimate concern: **does it really exist?**  
**why never observed?**

[Dif-Pradalier '10; Nakata '13; Villard '13; Norscini '14; Cartier-Michaud '14; Imadera, this conf.]  
[Dif-Pradalier PRL 15; Hornung, in prep.]

- ▶ avalanches interspersed between mean zonal flow layers
- ▶ shear flow layers associated to mean profile corrugations [ $\equiv \nabla(PV)$  or  $\nabla p$ ]
- ▶ other means to **characterise & observe the staircase?**