

# Recent progress towards a physics-based understanding of the H-mode transition

G.R. Tynan<sup>1</sup>, I. Cziegler<sup>1</sup>, P.H. Diamond<sup>1,2</sup>, M. Malkov<sup>2</sup>

<sup>1</sup>Center for Momentum Transport & Flow Organization (CMTFO)

<sup>2</sup>Center for Astrophysics & Space Science

University of California San Diego

La Jolla CA USA

A. Hubbard, J.W. Hughes, J.L. Terry, J.H. Irby

MIT Plasma Science and Fusion Center,

Cambridge MA

Acknowledgements:

L. Schmitz, G. McKee, Z. Yan, M. Xu, G.S. Xu

# Motivation:

---

- The physics of the L-H transition has been a key open issue
- Empirical  $P_{th}$  scalings have large uncertainties
  - Can we do better w/ a physics-based model?
- Can we learn how to control access to other improved confinement regimes w/o problems of H-mode
  - E.g. I-mode in AUG, C-Mod

# A picture of origin of L-H Transition is Emerging

---

- Equilibrium shear flows at boundary exist in L-mode
- Results in finite Reynolds stress that reinforces shear flow
- Flow drive from turbulence increases with L-mode heating power
- When flow drive rate exceeds turbulence drive rate ( $R_T > 1$ ) turbulence collapses
- This then allows grad-P buildup (H-Mode pedestal)

# Modeling & experiment suggest turbulent-driven ExB flow plays key role

---

- 0-d Predator-prey models predicts limit cycle regime (Kim & Diamond PRL'03)
  - *Prey*: turbulence,
  - *Predators*: zonal flow, mean flow
- Evidence for limit-cycle between L-mode & H-mode:
  - ASDEX (Zohm'94), TJ-II (Estrada'09,'15), AUG (Conway PRL'11), DIII-D (Schmitz PRL'13), ...
- 1-d Predator-prey model (Miki & Diamond, '12, '13)
  - key role of ZF-triggered turbulence collapse leading to grad-P buildup
  - Interplay between ZF and mean (ion diamag) flows
  - 1-d front propagation effects (Schmitz'12, Estrada'15)

# ExB flow drive leads to turbulence collapse in simplified model

Turbulent Kinetic Energy:

$$\frac{\partial \tilde{K}}{\partial t} = (\gamma_{in} - \gamma_{corr}^{pl}) \tilde{K} - P - \partial_r \tilde{T}$$

Large-scale Shear Flow Kinetic Energy:

$$\frac{\partial \bar{K}}{\partial t} = P - \partial_r \bar{T} - v_{LF} \bar{K}$$

Turbulence Collapse Condition:

$$R_T \equiv \frac{P - \partial_r \bar{T}}{(\gamma_{in} - \gamma_{corr}^{pl}) \tilde{K}} > 1$$

Manz, PoP'12, Cziegler'15

See also Vianello, PPCF 2005 & 2006  
Hidalgo, Sanchez, JNM & PPCF 05,06

Definitions:

$$\tilde{K} = \frac{1}{2} \langle \tilde{v}^2 \rangle \quad \bar{K} = \frac{1}{2} \langle V^{LF} \rangle^2$$

$$P = \langle \tilde{v}_r \tilde{v}_\theta \rangle \frac{\partial \langle V^{LF} \rangle}{\partial r}$$

$$\tilde{T} = \langle \tilde{v}_r \tilde{v}_\theta^2 \rangle \quad \bar{T} = \langle \tilde{v}_r \tilde{v}_\theta \rangle \langle V^{LF} \rangle$$

$$\gamma_{in} = \gamma_{in}(\nabla n, \nabla T, \bar{V}_E')$$

$$\langle V^{LF} \rangle = \langle V_{ExB} \rangle + \langle V_i^{dia} \rangle$$

Ion Diamagnetic Flow  
**Survives** Turbulence  
Collapse & Locks In  
*H-mode State*

# Leads to closed form reduced predator-prey model

---

- Model closed with few key assumptions (Kim&Diamond PRL'03)

- $q \propto -\langle \tilde{v}_\perp^2 \rangle \tau_{corr} \nabla p_i$

- $\bar{V}_{i,dia} \propto \nabla p_i$

- $\langle \tilde{v}_r \tilde{v}_\theta \rangle \propto \frac{\langle \tilde{v}_\perp^2 \rangle \bar{V}'_{ExB, LF}}{1 + \alpha \bar{V}'_E{}^2}$

# Recent work on C-Mod validates key assumption in predator-prey model

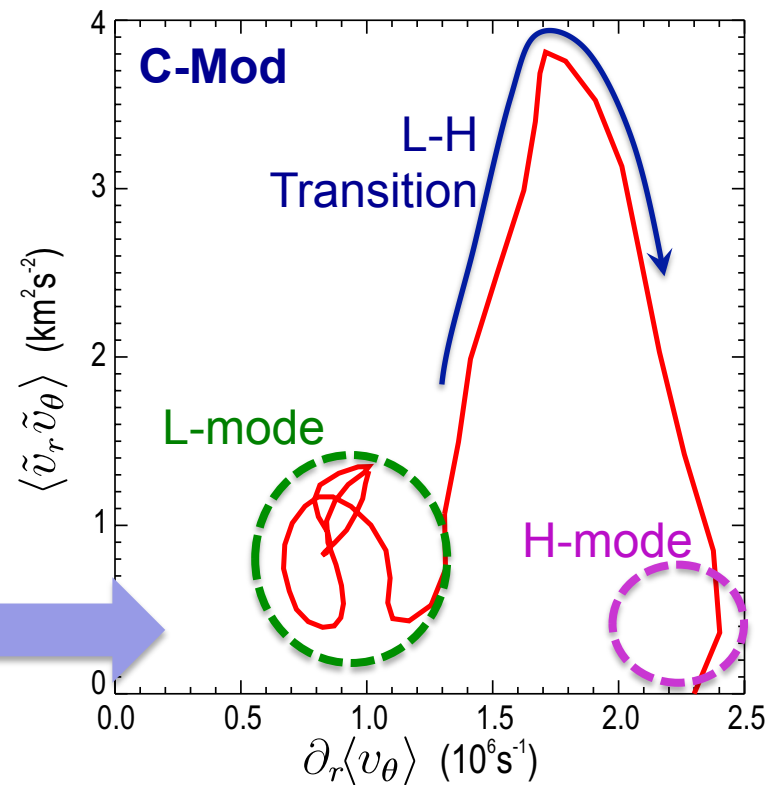
- Model (Kim&Diamond PRL'03) was closed with few key assumptions

- $q \propto -\langle \tilde{v}_\perp^2 \rangle \tau_{corr} \nabla p_i$

- $\bar{V}_{i,dia} \propto \nabla p_i$

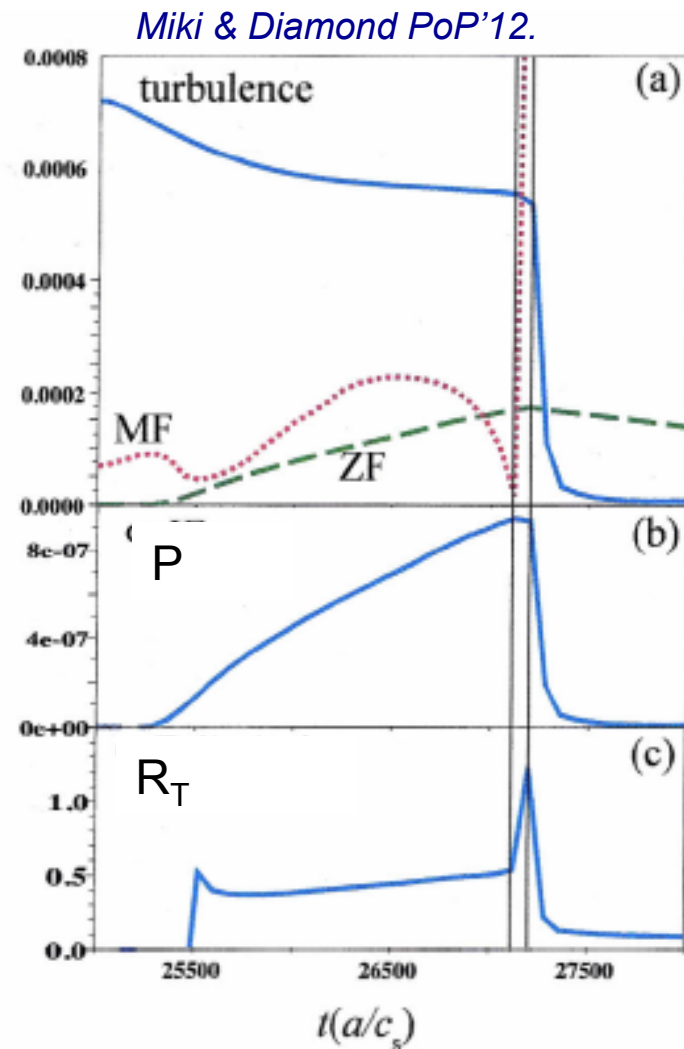
- $\langle \tilde{v}_r \tilde{v}_\theta \rangle \propto \frac{\langle \tilde{v}_\perp^2 \rangle \bar{V}'_{ExB, LF}}{1 + \alpha \bar{V}'_E{}^2}$

Experiments support this ansatz



# 1-D Reduced Model Shows L-H Transition Dynamics

- Turbulent-driven  $m,n=0,0$  ExB flow builds up & regulates turbulence
- ZF production,  $P$ , grows with heat flux
- Turbulence collapse when flow drive exceeds growth rate ( $R_T > 1$ )
- $\text{grad-}P_{\text{ion}}$  ExB flow then builds; turbulent-driven  $m,n=0$  ExB decays
- Strong  $\text{grad-}P_{\text{ion}}$  Mean Flow (MF) locks-in H-mode





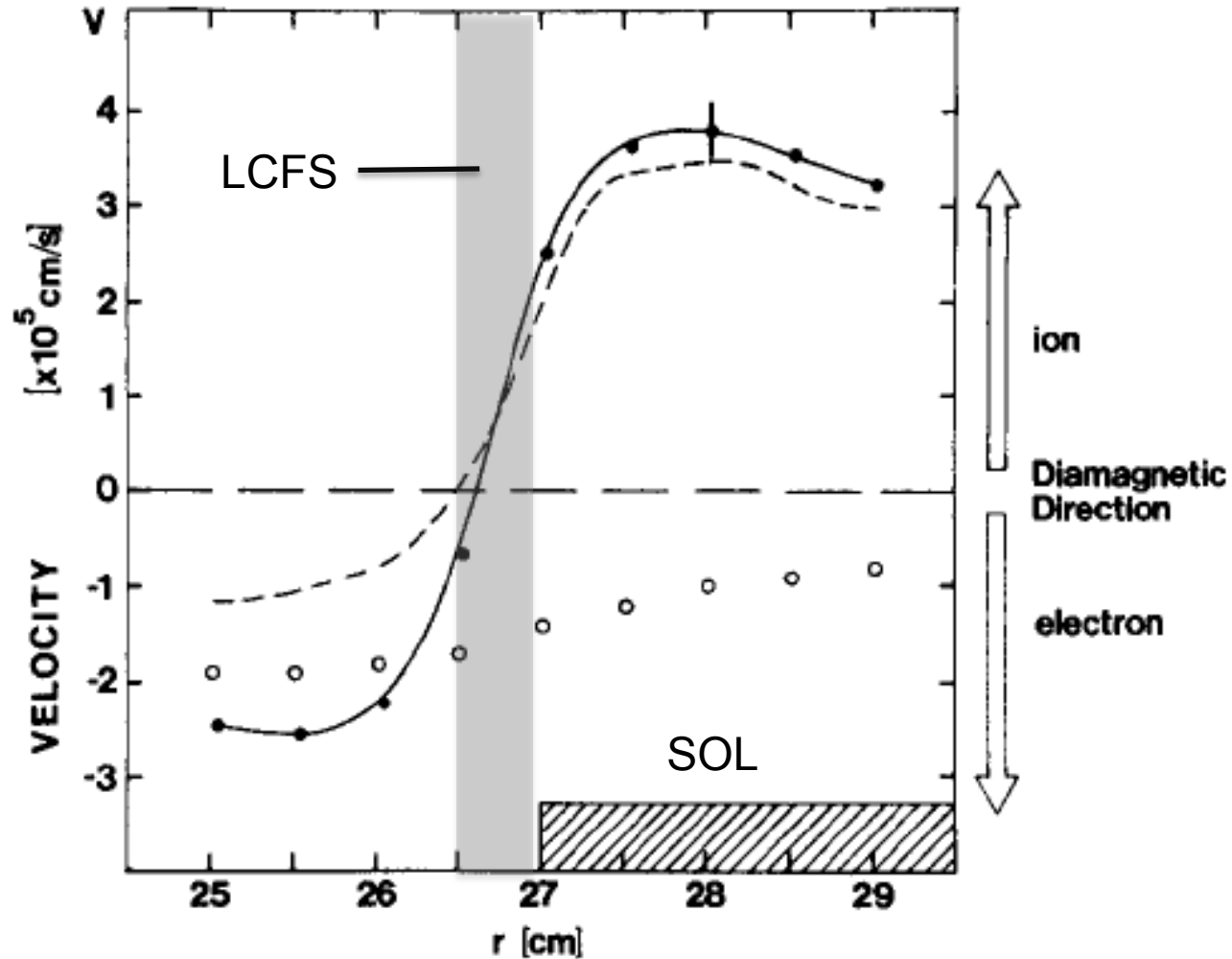
# A picture of origin of L-H Transition is Emerging

---

- **Equilibrium shear flows at boundary exist in L-mode**
- Results in finite Reynolds stress that reinforces shear flow
- Flow drive from turbulence increases with L-mode heating power
- When flow drive rate exceeds turbulence drive rate ( $R_T > 1$ ) turbulence collapses
- This then allows grad-P buildup (H-Mode pedestal)

# A Shear Layer Sustained by Equilibrium Exists at LCFS

TEXT Ritz PF'84



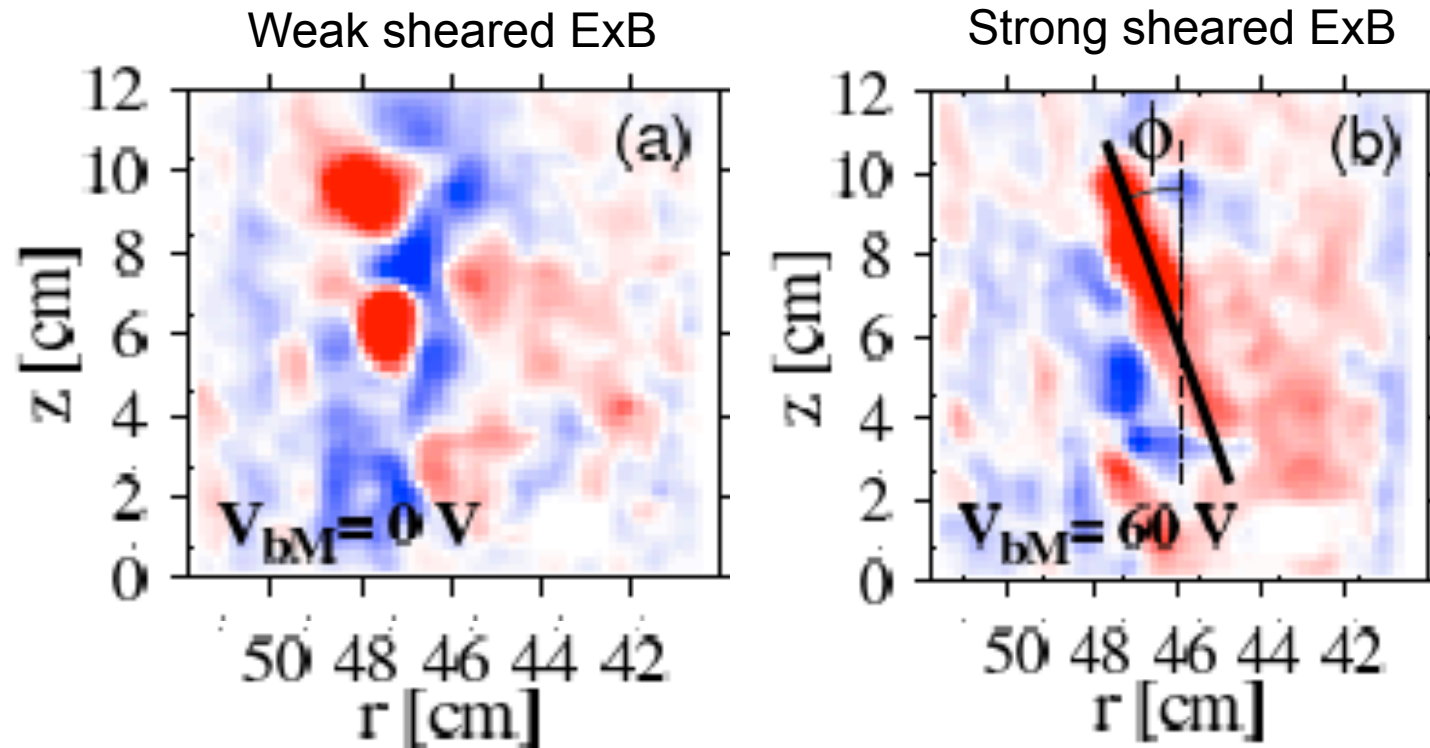
# A picture of origin of L-H Transition is Emerging

---

- Equilibrium shear flows at boundary exist in L-mode
- **Results in finite Reynolds stress that reinforces shear flow**
- Flow drive from turbulence increases with L-mode heating power
- When flow drive rate exceeds turbulence drive rate ( $R_T > 1$ ) turbulence collapses
- This then allows grad-P buildup (H-Mode pedestal)

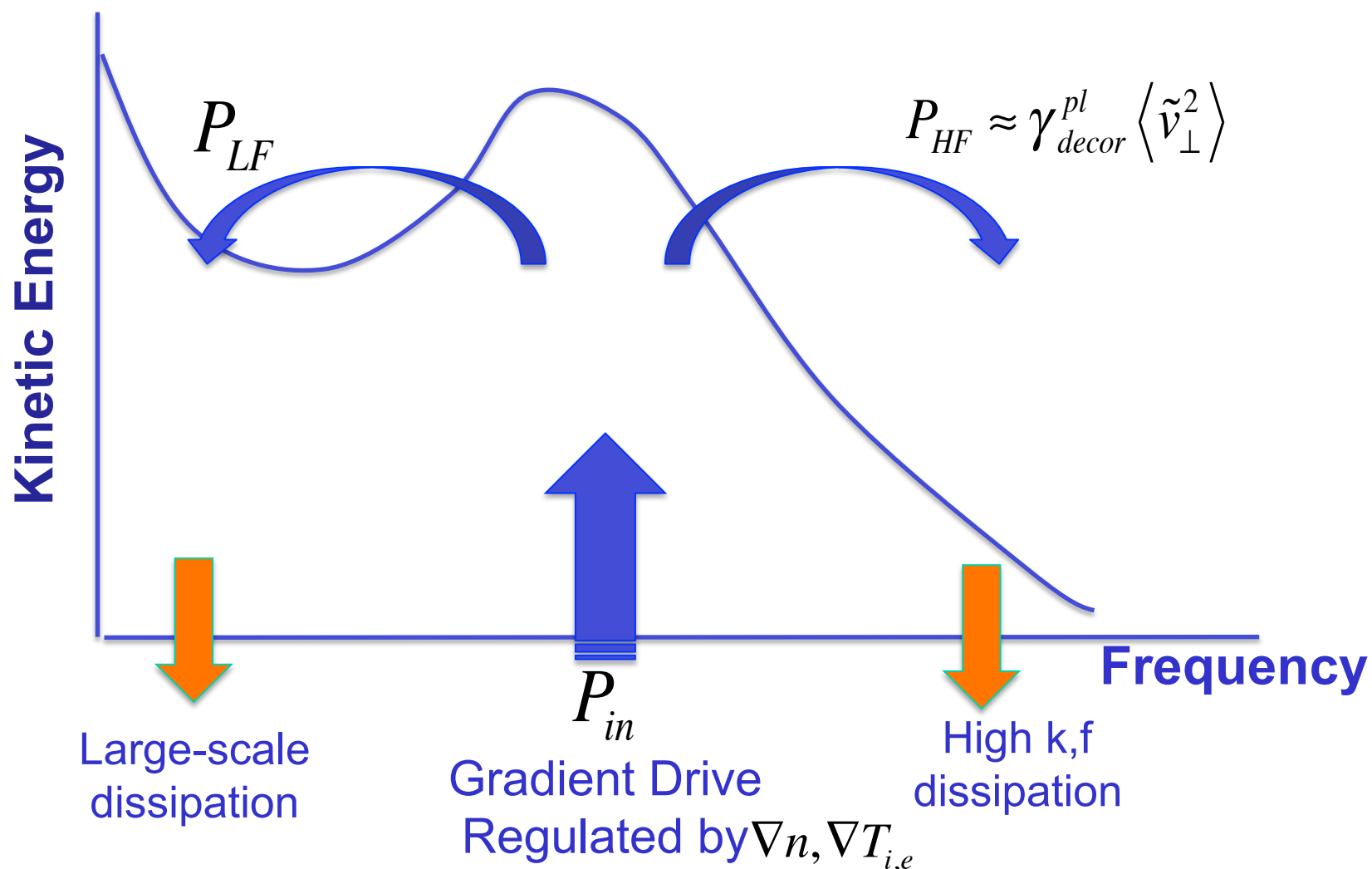
# Sheared ExB flow tilts & stretches eddies

TEXTOR Shesterikov, Xu et al PRL 2013



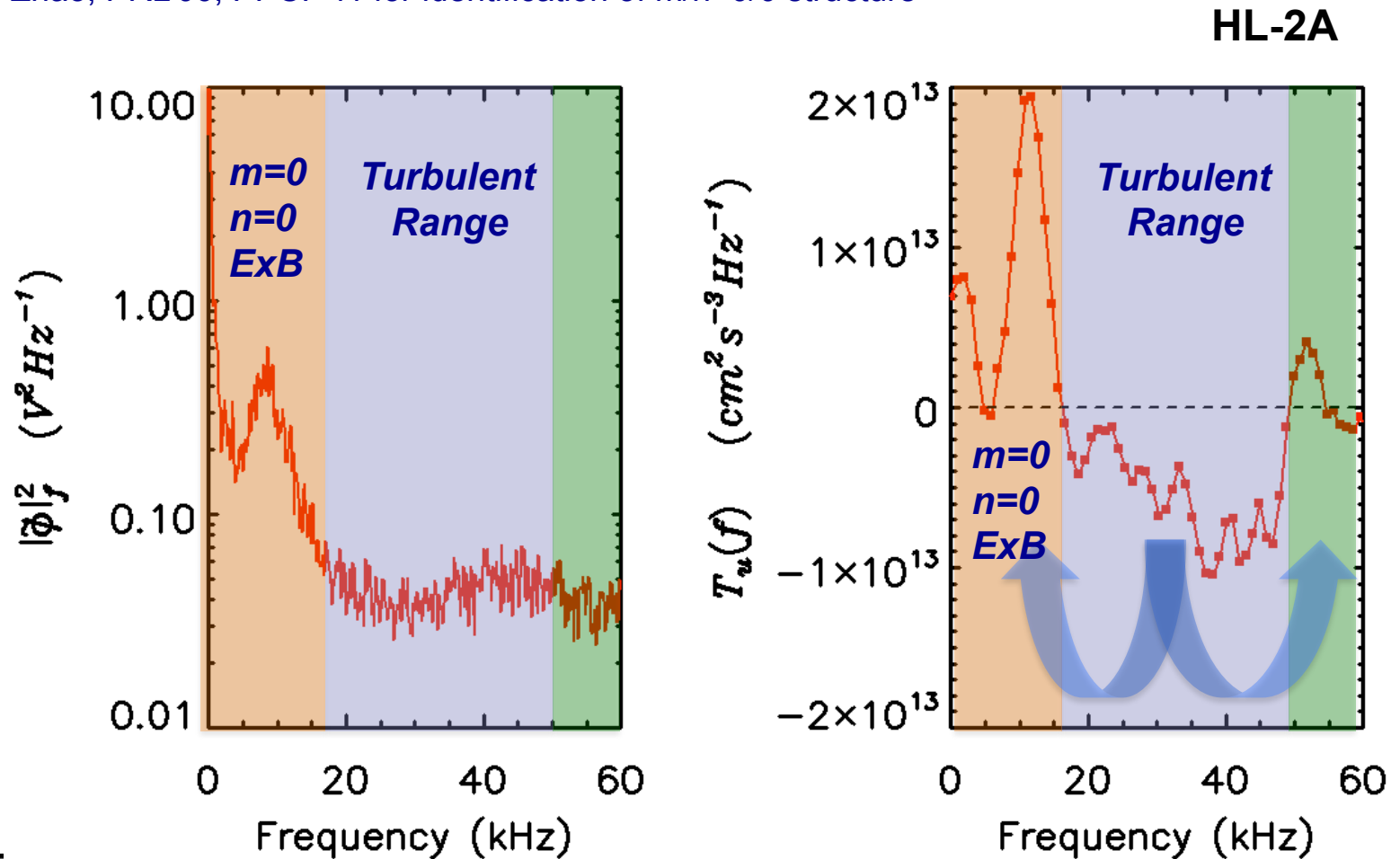
Eddy tilting correlates  $V_r$  and  $V_\theta$   
→ Non-zero Reynolds Stress  
→ Turbulence can drive sheared flow

# Spectral representation: L-mode turbulence established from energy input, nonlinear transfer & dissipation



# Weakly heated L-mode experiments confirm this picture

See M. Xu PRL '12 for flow drive physics & M. Xu PoP'10 for technique  
See K. Zhao, PRL'06, PPCF'11 for identification of  $m/n=0/0$  structure

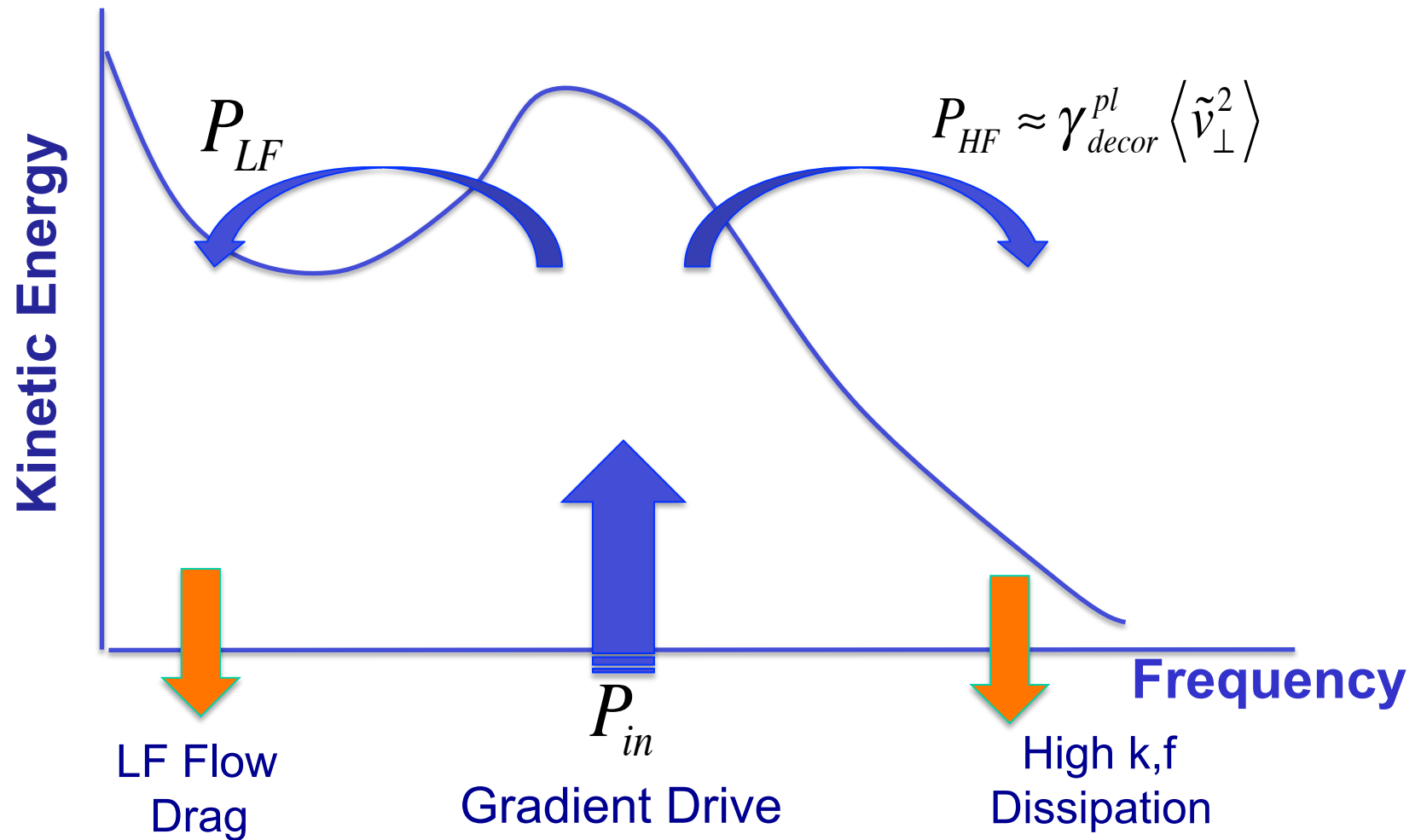


# A picture of origin of L-H Transition is Emerging

---

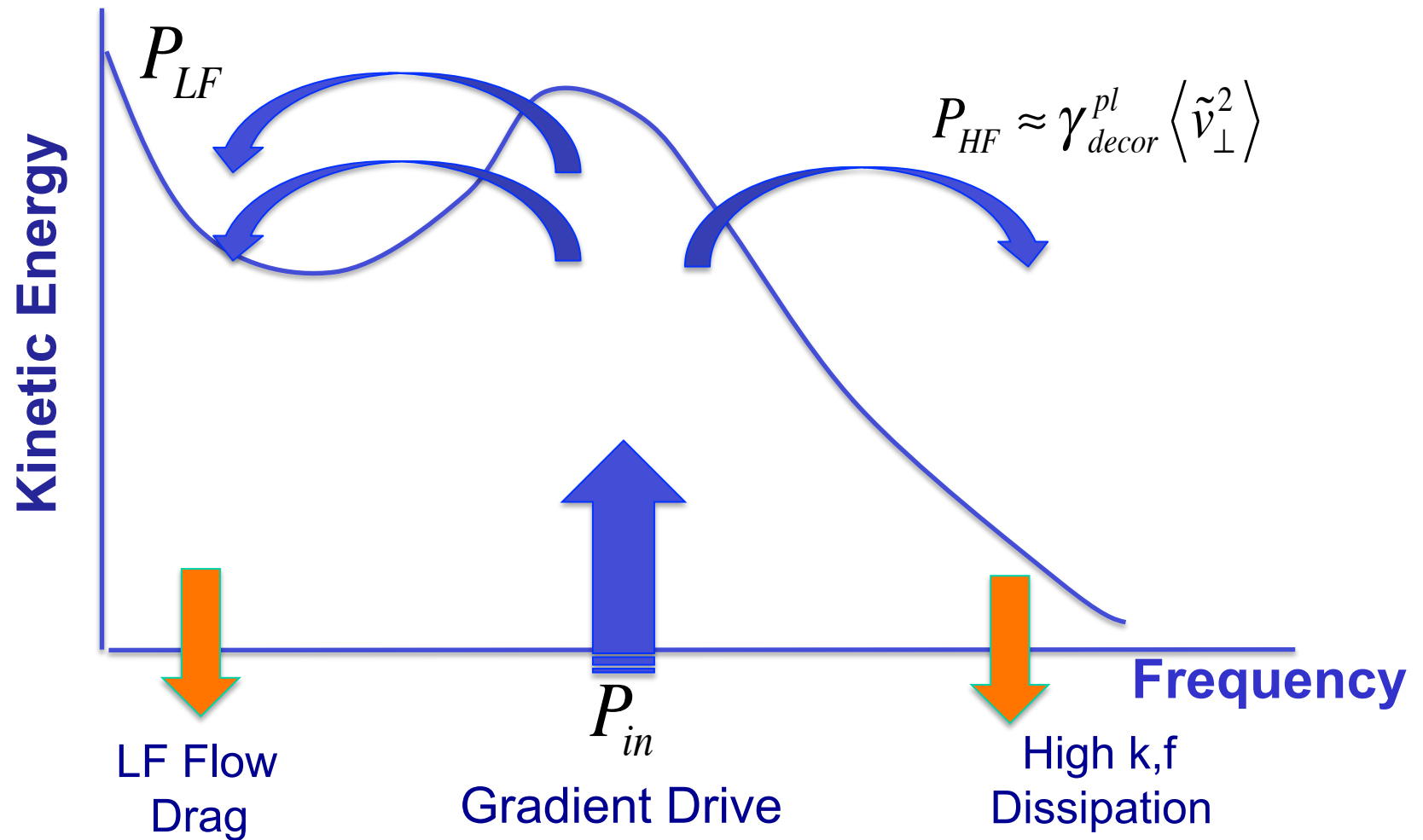
- Equilibrium shear flows at boundary exist in L-mode
- Results in finite Reynolds stress that reinforces shear flow
- **Flow drive from turbulence increases with L-mode heating power**
- When flow drive rate exceeds turbulence drive rate ( $R_T > 1$ ) turbulence collapses
- This then allows grad-P buildup (H-Mode pedestal)

# Spectral representation: L-mode turbulence established from energy input, nonlinear transfer & dissipation





# Power transfer to shear flow increases with plasma heating

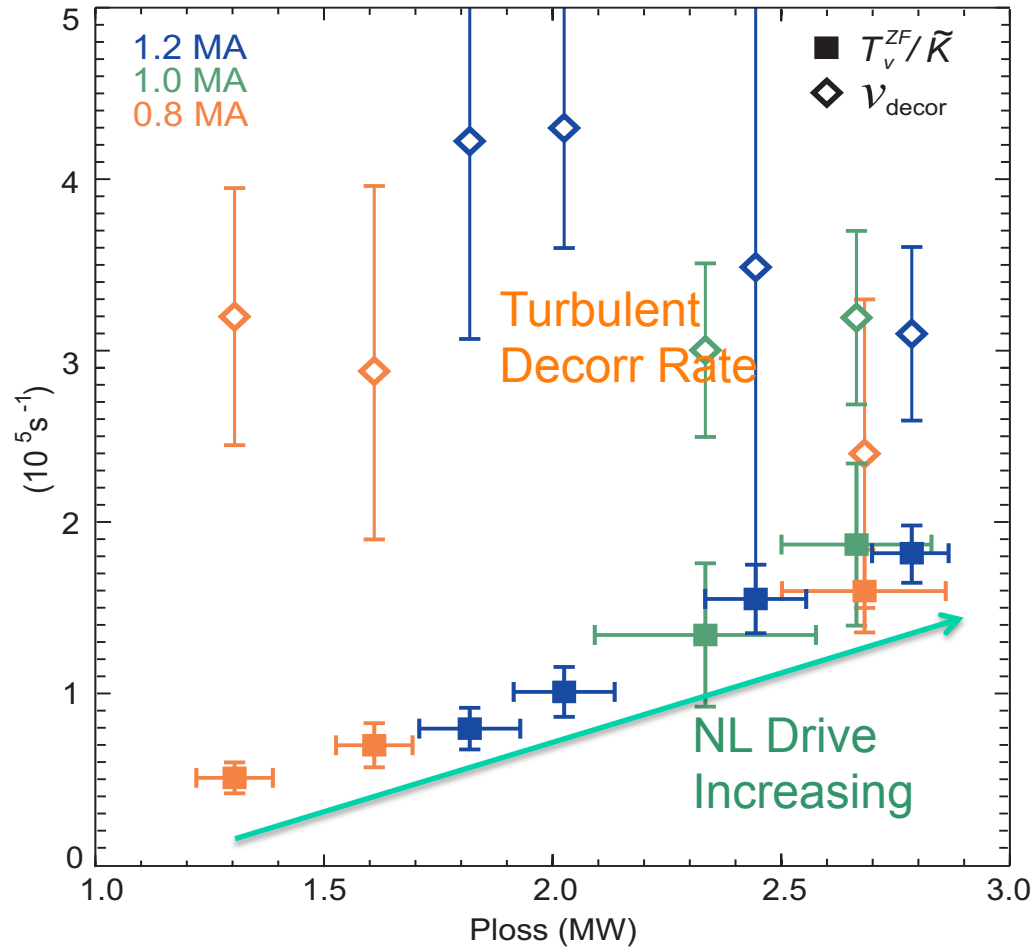


# Nonlinear shear flow drive becomes important at higher heat flux

Cziegler, et al, NF 2015 (accepted),

C-Mod

$$T_v^{ZF} = -\text{Re} \left[ \sum_{f_1; f=f_{ZF}} \langle \tilde{v}_\theta(f) \tilde{v}_r(f-f_1) \partial_r \tilde{v}_\theta(f_1) \rangle \right]$$

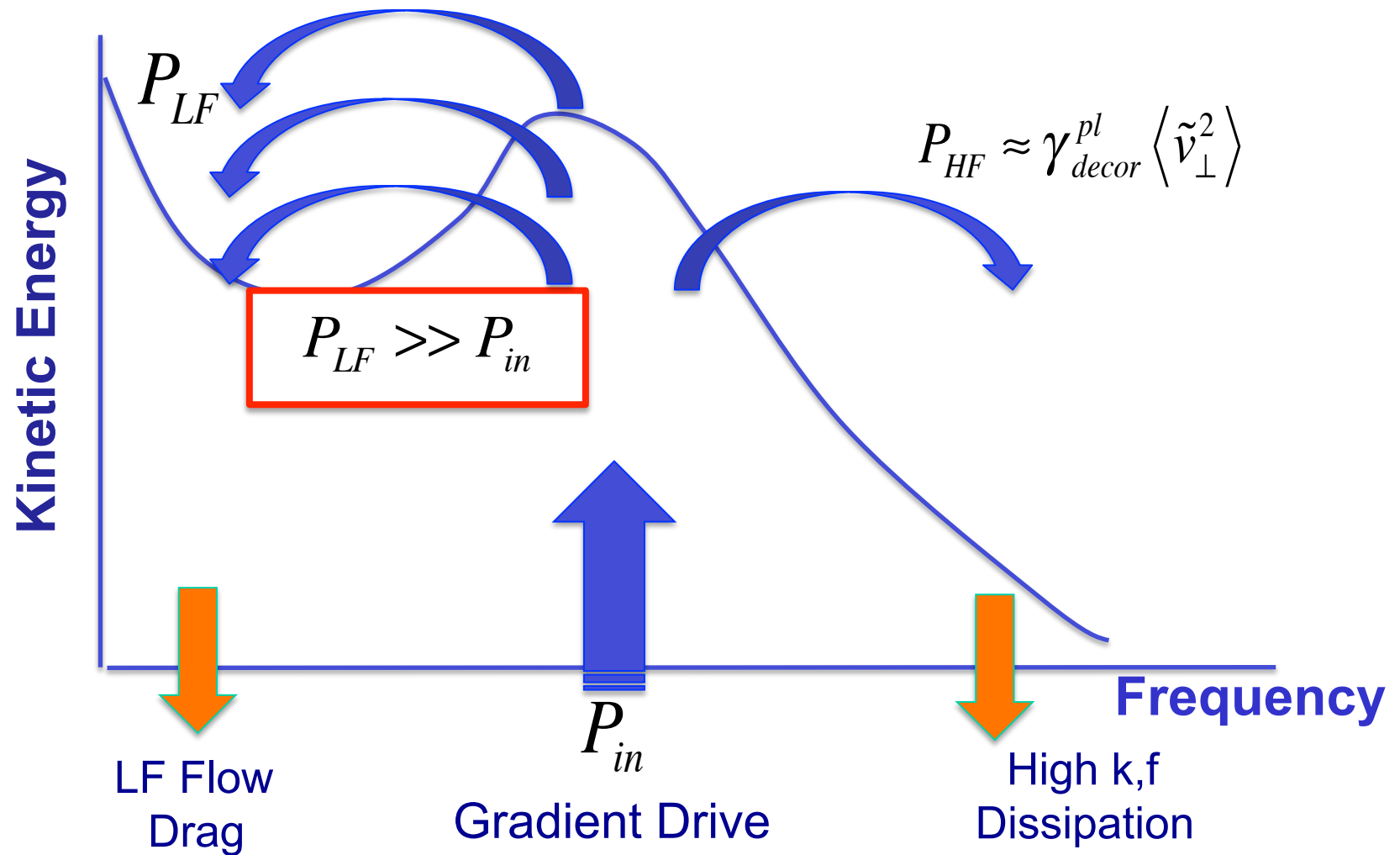


# A picture of origin of L-H Transition is Emerging

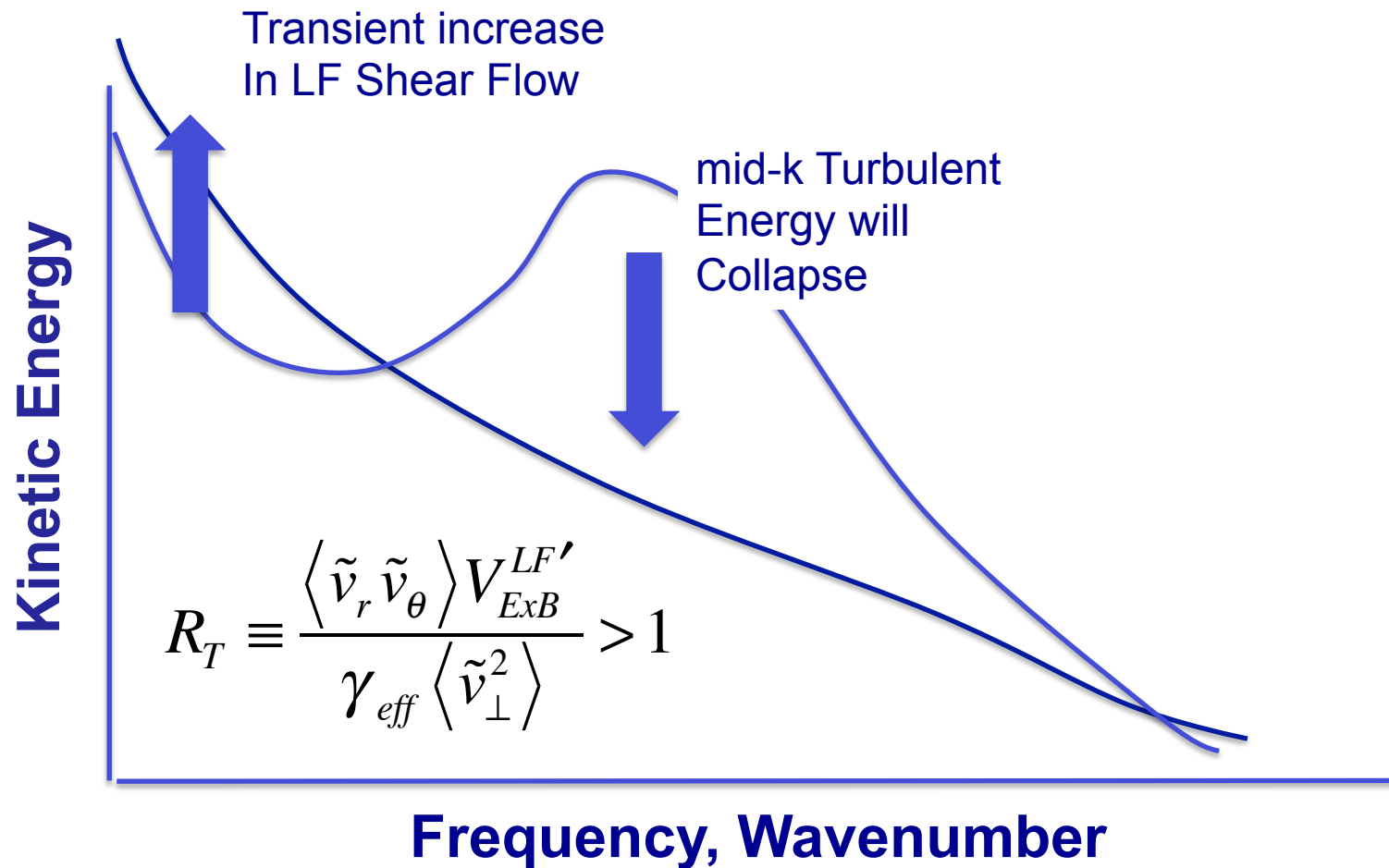
---

- Equilibrium shear flows at boundary exist in L-mode
- Results in finite Reynolds stress that reinforces shear flow
- Flow drive from turbulence increases with L-mode heating power
- **When flow drive rate exceeds turbulence drive rate ( $R_T > 1$ ) turbulence collapses**
- This then allows grad-P buildup (H-Mode pedestal)

# Further heating leads rate of flow drive exceeding rate of energy input into turbulence



# At this point, turbulence cannot be sustained...and thus turbulence amplitude collapses



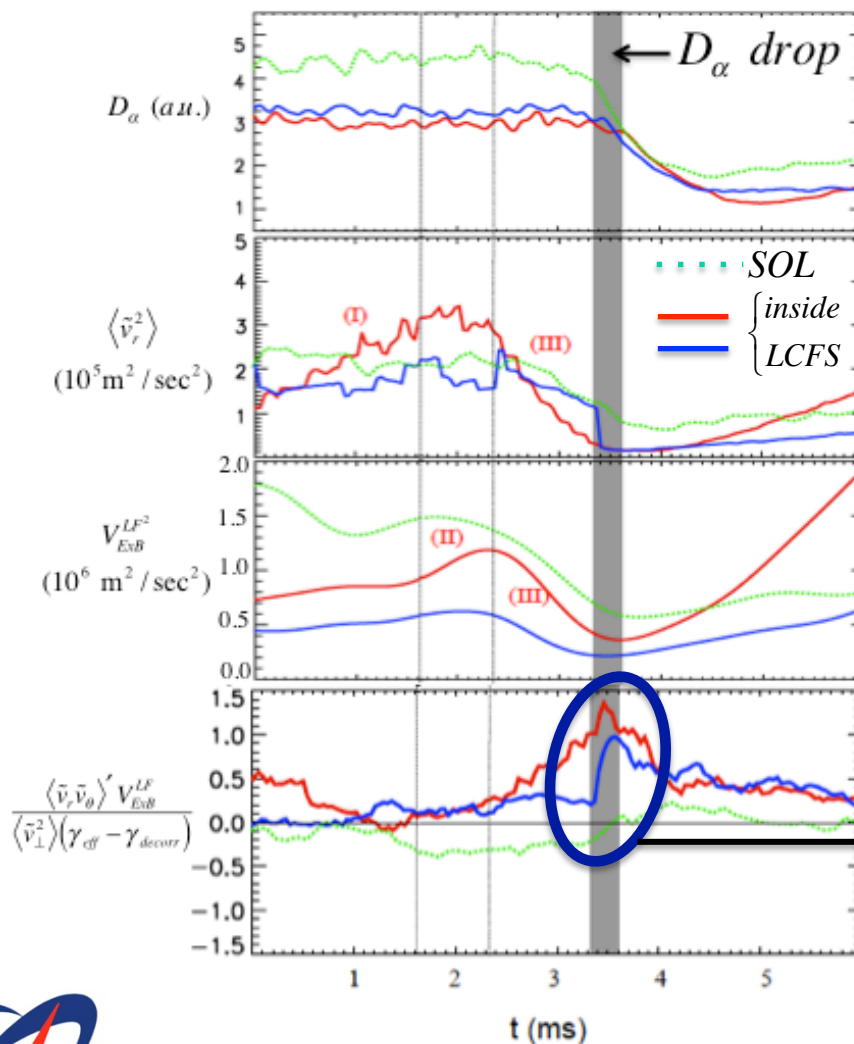
# RECENT EXPERIMENTS SHOW THIS BEHAVIOR

---

- EAST (Manz et al PoP'12, GS Xu et al NF'14)
- TEXTOR (Shesterikov et al PRL'13)
- DIII-D (Yan et al PRL'14)
- ALCATOR C-Mod (Cziegler, PPCF'14, NF'15)

# L-H Transition When m,n=0 LF ExB Drive Exceeds Energy Input Rate into Turbulence

Manz et al, PoP Aug 2012



- Turbulence Energy & LF ExB Energy Increase
- Power Transfer Increases
- Power Transfer Grows to ~Equal Turbulent Energy Input Rate
- L-H Transition Occurs

$$\langle \tilde{v}_r \tilde{v}_\theta \rangle V_{ZF}' \approx (\gamma_{eff} - \gamma_{decorr}) \tilde{v}_\perp^2$$

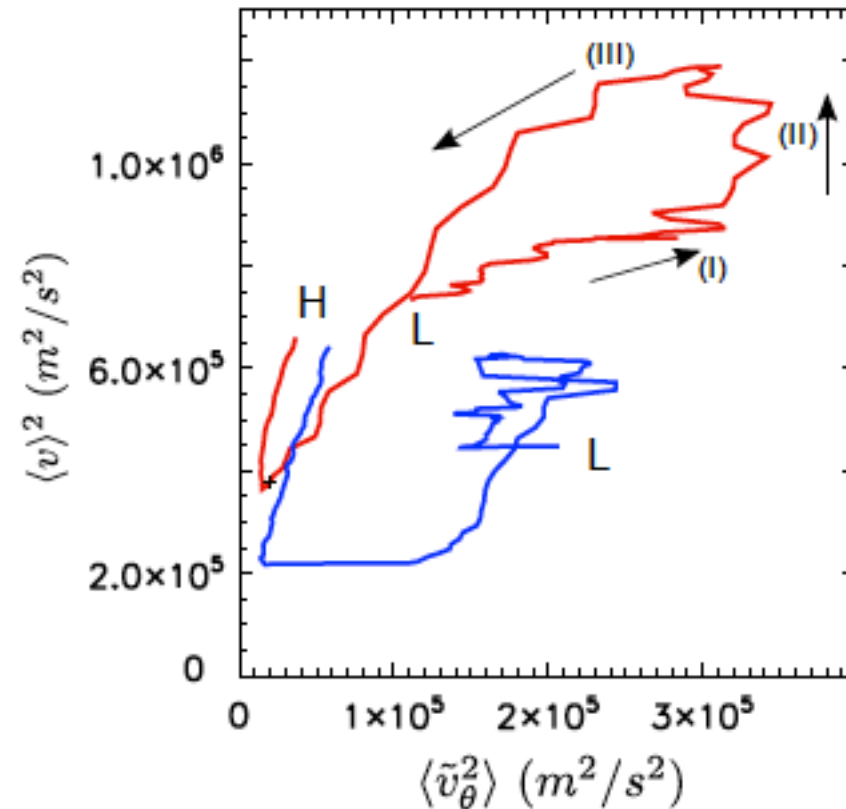


# Turbulence & LF $m,n=0$ ExB Energy Exhibit **One Orbit of Phase Space** Before the L-H Transition

Manz et al, PoP Aug 2012

**EAST**

- Turbulence Increases in L-mode (I)
- Power Transfer Increases the LF Shear Flow (I-II)
- Flow Saturates Turbulence (II-III)
- Transition to H-mode State Locks in LF Shear Flow
- Strong radial dependence



L-H Transition Appears to be Degenerate  
Case of Multi-orbit LCO Regime



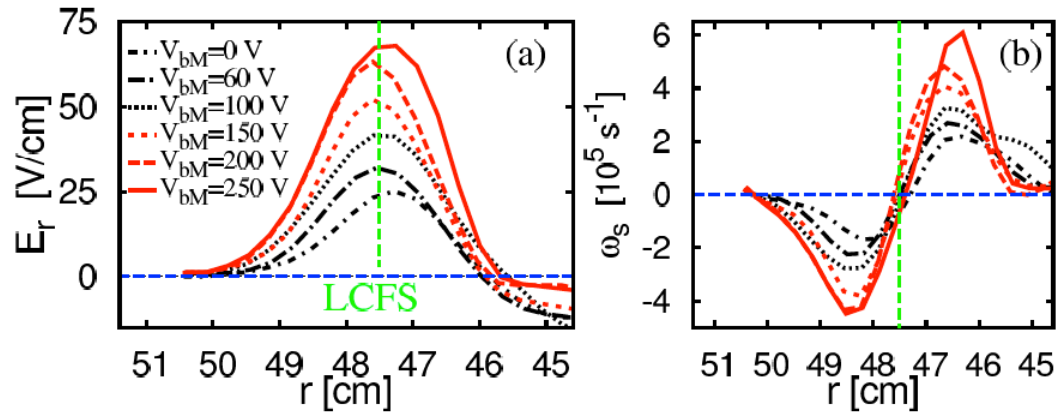


---

# **TEXTOR:** Energy Transfer During Biased L-H Transition

# Turbulent-driven ExB Flow Plays Roll in Biased H-mode Transition

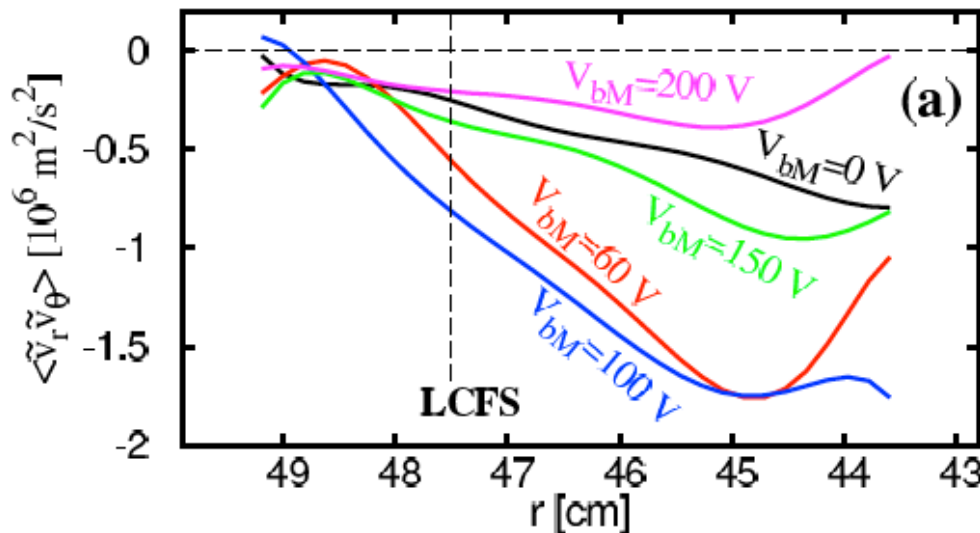
**TEXTOR** Shesterikov, Xu et al, PRL 2013



For lower values of shearing

$$\langle \tilde{v}_r \tilde{v}_\theta \rangle \propto \bar{V}'_E$$

→ Negative Viscosity

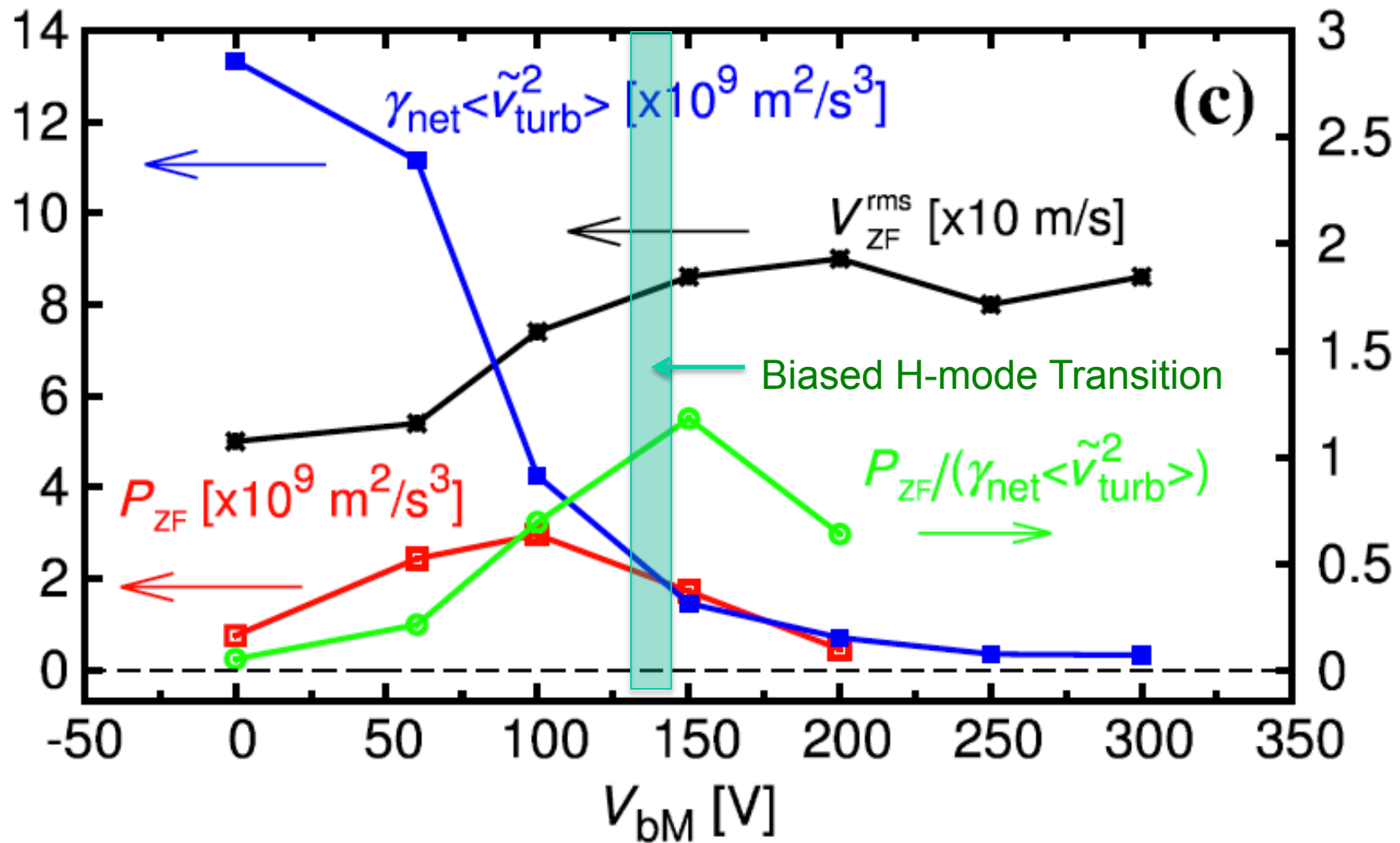


Stress reinforces ExB flow

Stress saturates and decreases for larger values of stress

# Biased H-mode Onset Occurs When Power Transfer to LF ExB Shear Flow Exceeds Threshold

TEXTOR Shesterikov, Xu et al, PRL 2013

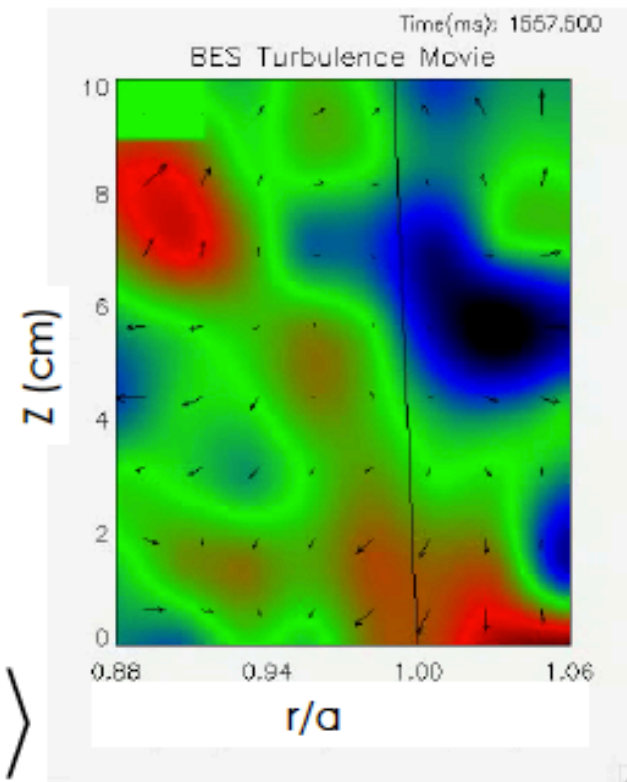


---

## **DIII-D:** L-H Transition Studies via BES velocimetry

# Turbulent Velocity Fluctuation Measured from Image-based Velocimetry

- **Vector-matching frame by frame to infer short time scale velocity fluctuation [1]**
  - Orthogonal Dynamic Programming
- **With measured  $\tilde{V}_\theta$  and  $\tilde{V}_r$** 
  - Analyze velocity fluctuation spectrum
  - Infer Reynolds stress, in principal proportional to the electrostatic Reynolds stress:  $\langle \tilde{V}_r \tilde{V}_\theta \rangle$   
(Requires further validation)

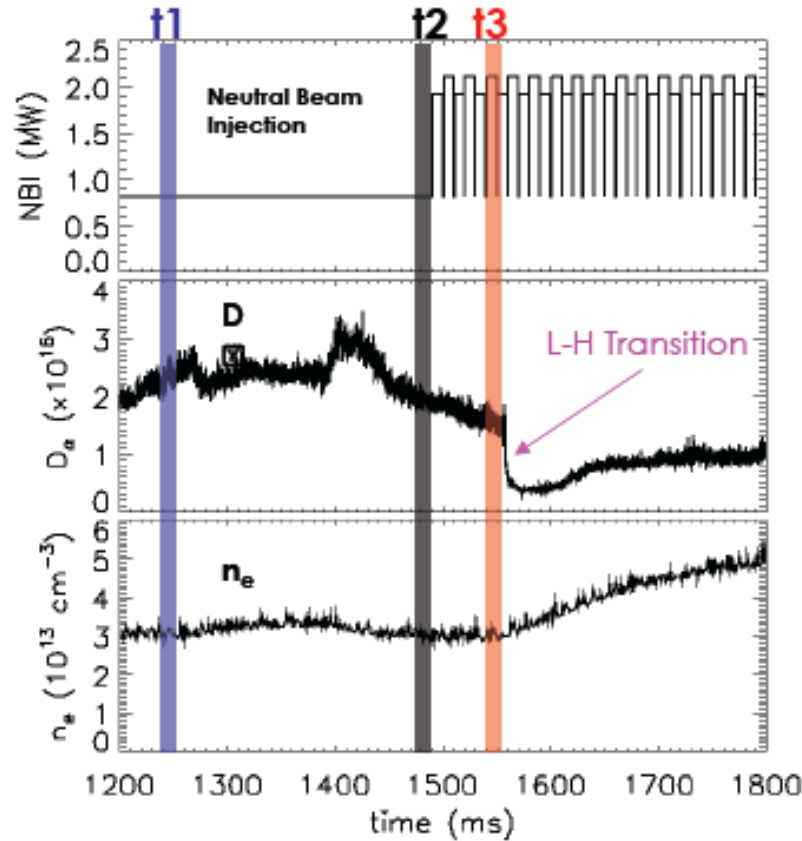


Z. Yan et al, PRL 2014

[1] G. McKee, et al., RSI, 75, 3490, 2004

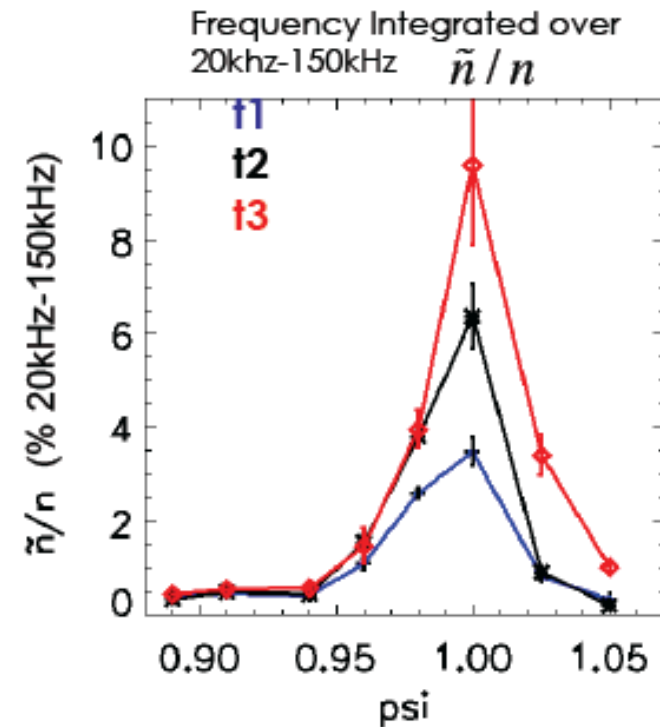
# Turbulence Increases with Time and Power Approaching L-H Transition

Z. Yan et al, PRL 2014



$t_1, t_2: P_{\text{INPUT}} < P_{\text{L-H}}$

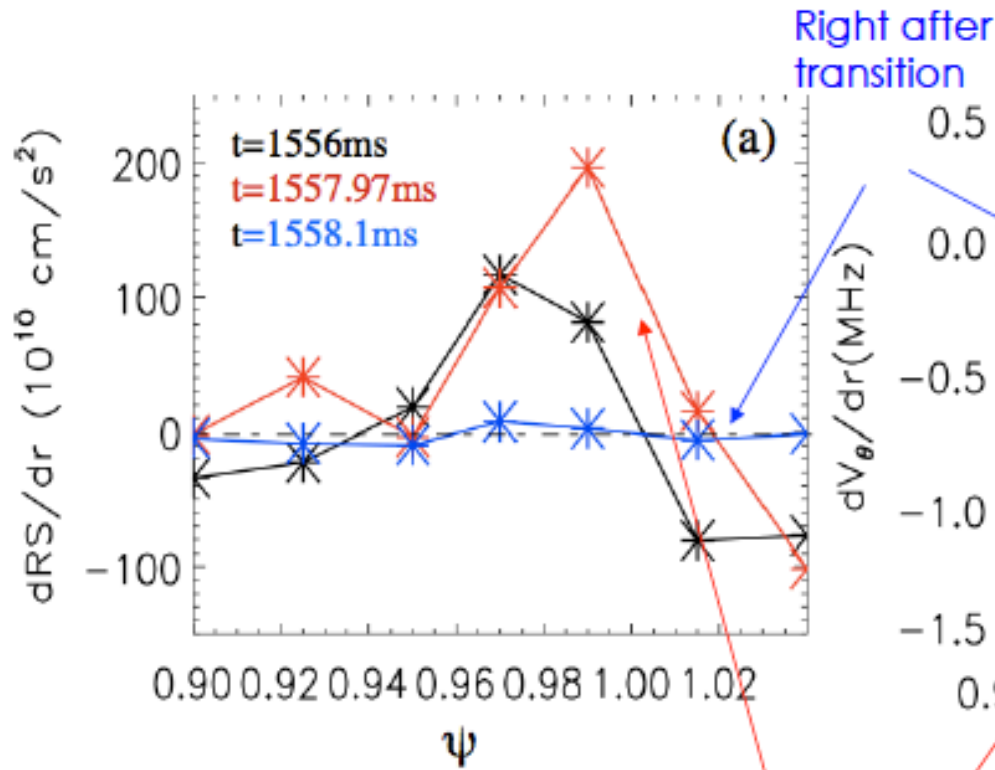
$t_3: P_{\text{INPUT}} \geq P_{\text{L-H}}$



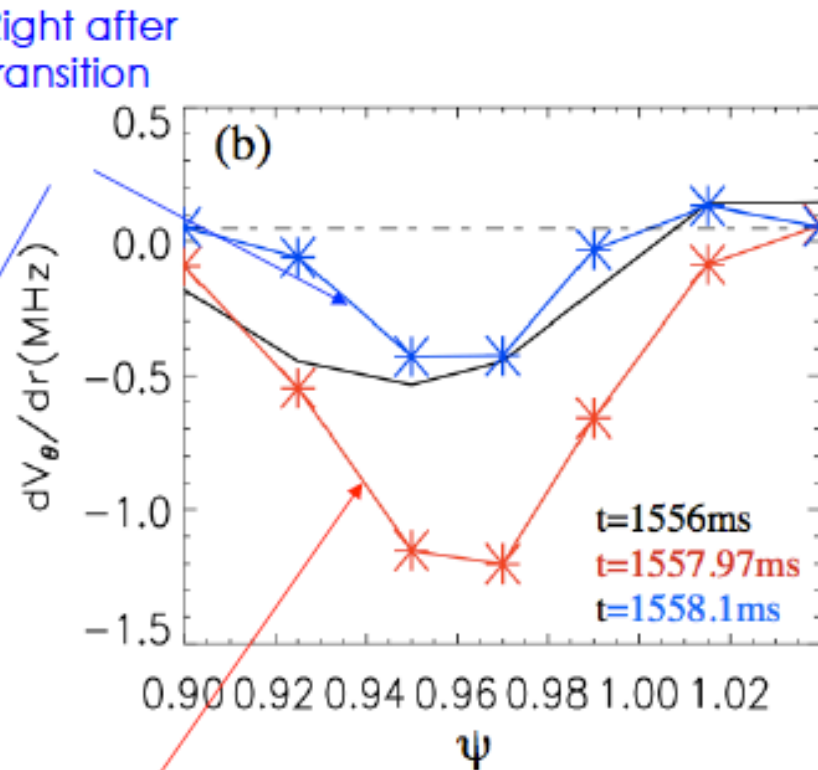
Increased turbulence provides increased drive for zonal flow

# The Final Increases Are Localized in The Pedestal Region

Radial profile of Reynolds stress gradient



Radial profile of velocity shear



Right before transition

Z. Yan et al, PRL 2014

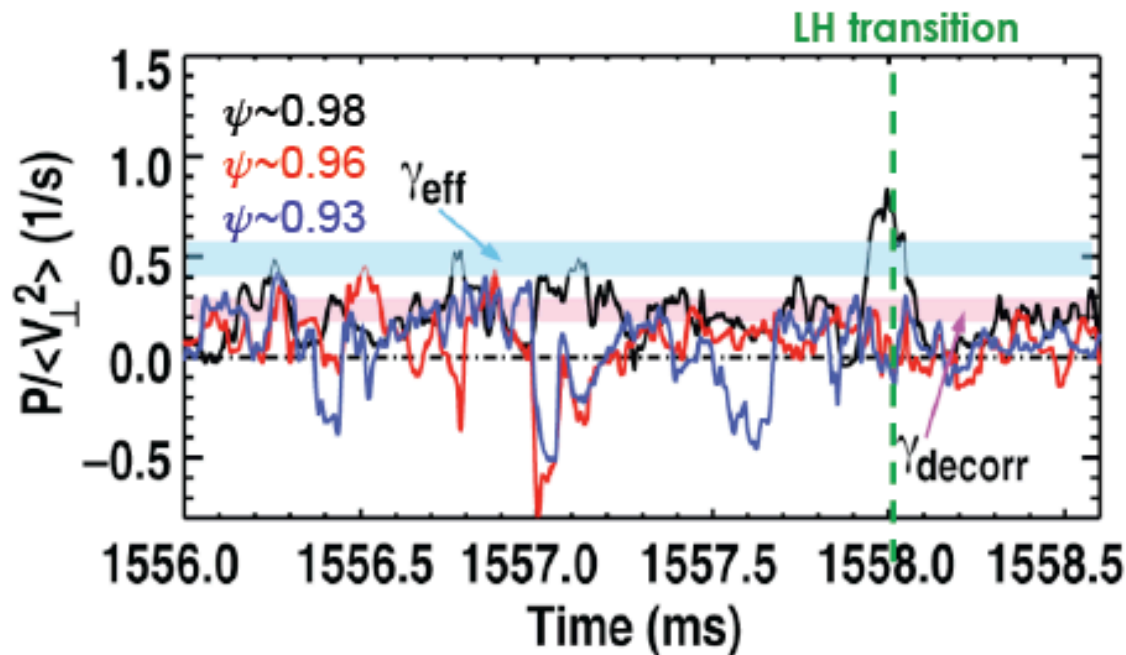
# Energy Transfer Jumps Up Right Before the Transition and Localized at the Plasma Edge Region

Evolution of zonal flow and turbulence [2]

Z. Yan et al, PRL 2014

$$\frac{1}{2} \frac{\partial \langle \tilde{v}^2 \rangle}{\partial t} = \gamma_{\text{eff}} \langle \tilde{v}^2 \rangle - \langle \tilde{v}_r \tilde{v}_\theta \rangle \frac{\partial V_E^{LF}}{\partial r}$$

$$\frac{1}{2} \frac{\partial \langle V_E^{LF^2} \rangle}{\partial t} = + \langle \tilde{v}_r \tilde{v}_\theta \rangle \frac{\partial V_E^{LF}}{\partial r} - \gamma_{LF} \langle V_E^{LF^2} \rangle$$



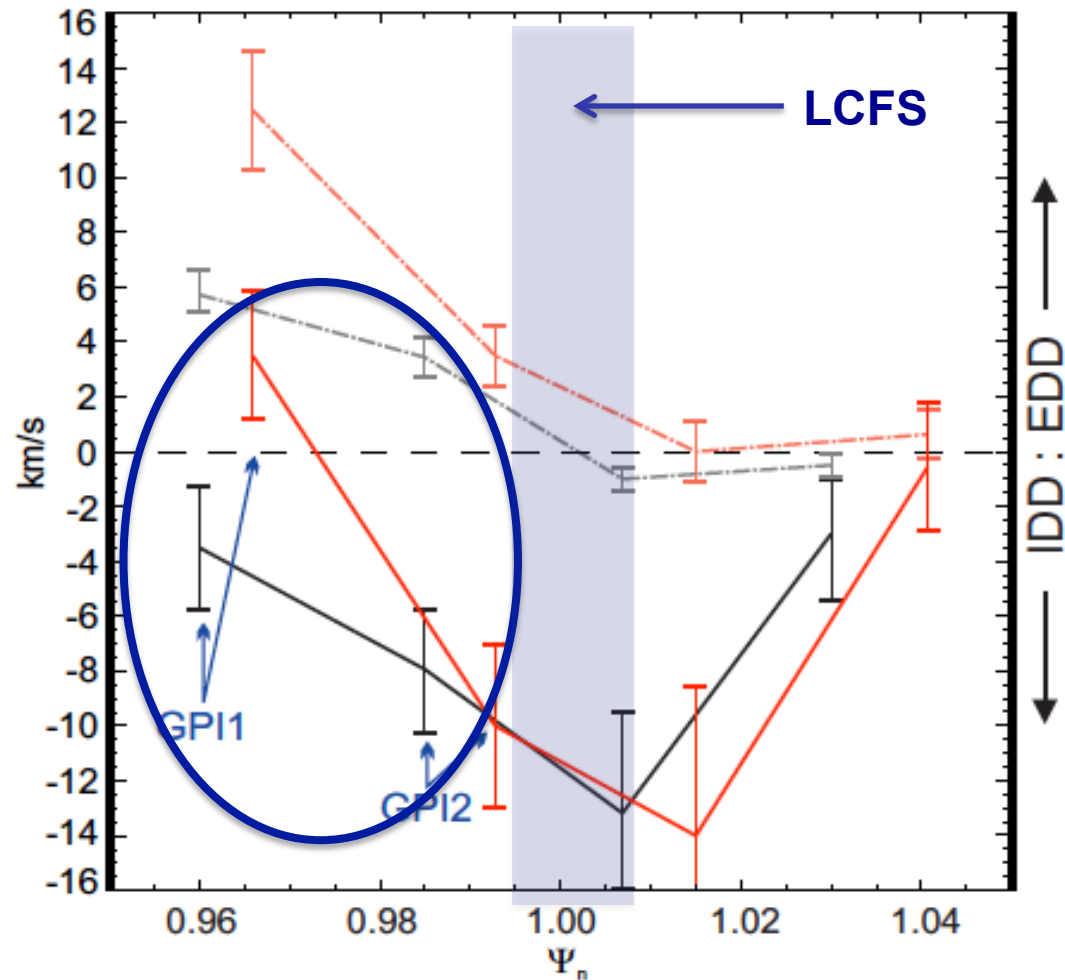


---

# **C-Mod:** L-H Transition Studies via He-GPI velocimetry

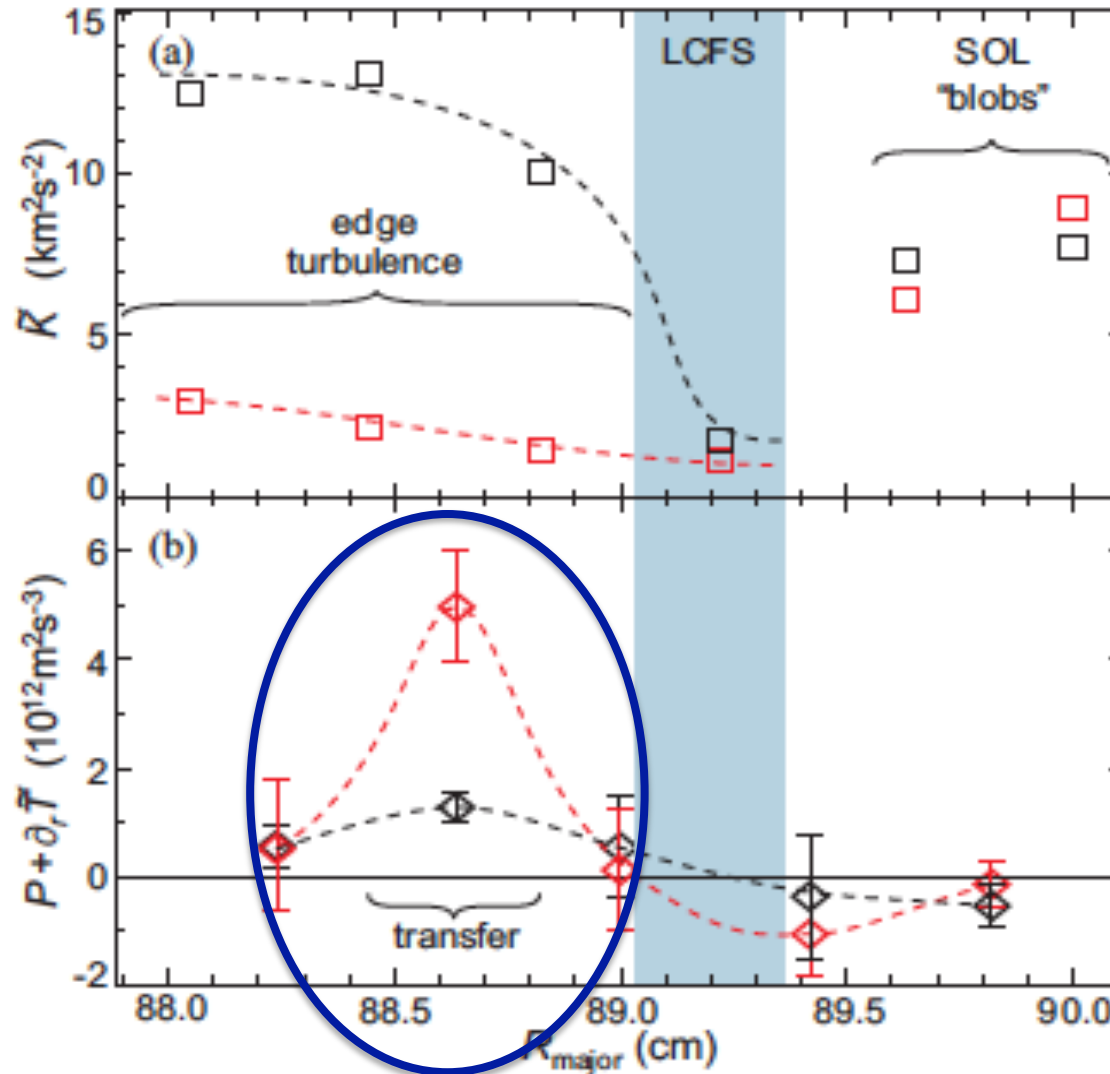
# Observe a “kick” in $V_{\text{pol}}$ at L-H transition

Cziegler et al PPCF'14



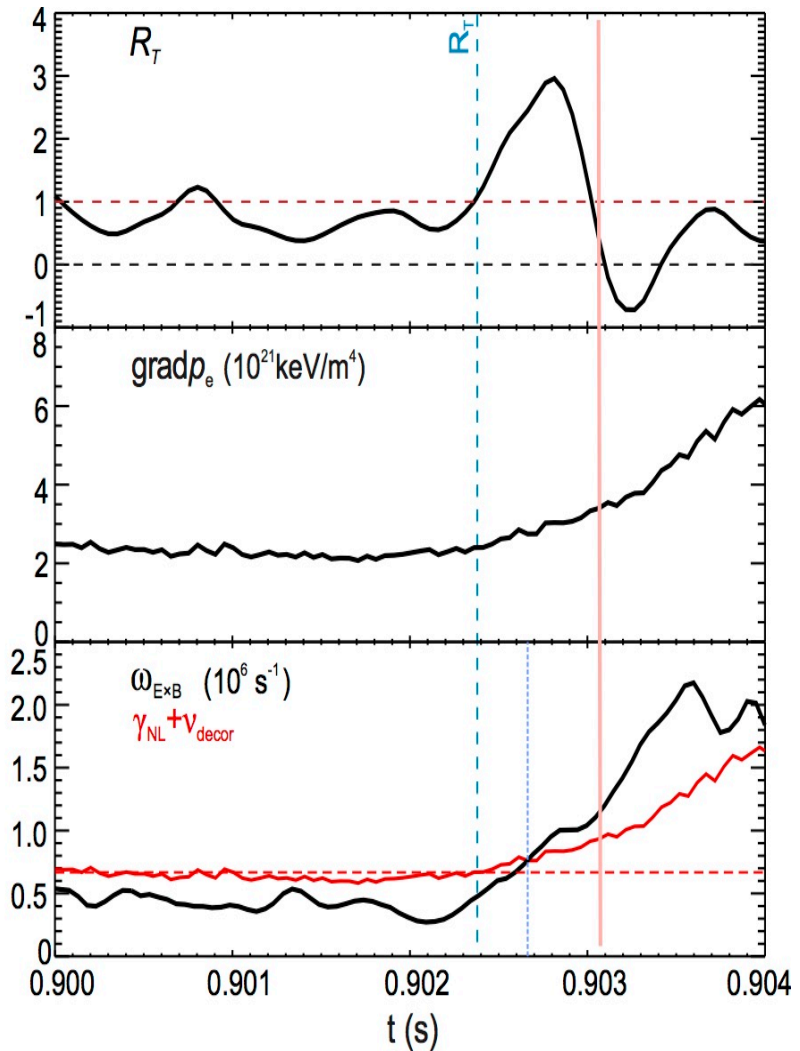
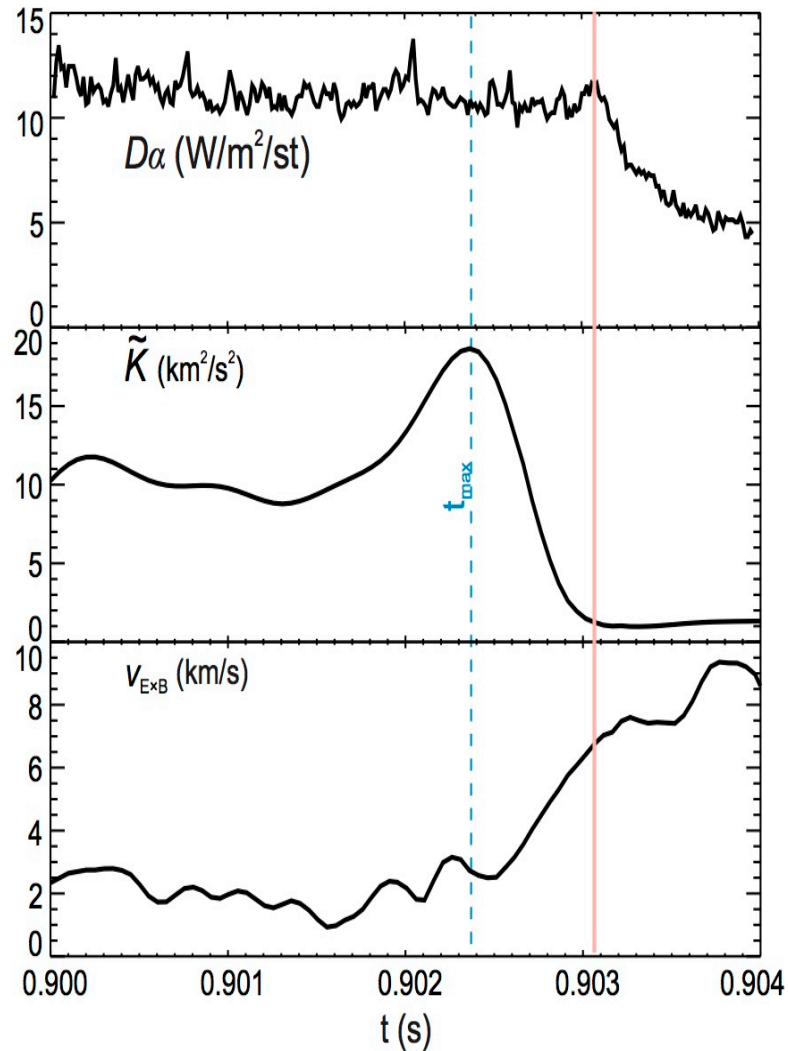
# Spatial localization of flow drive & turbulence quenching

Cziegler, NF'15



# Time sequence of transition

Tynan, Cziegler et al, submitted



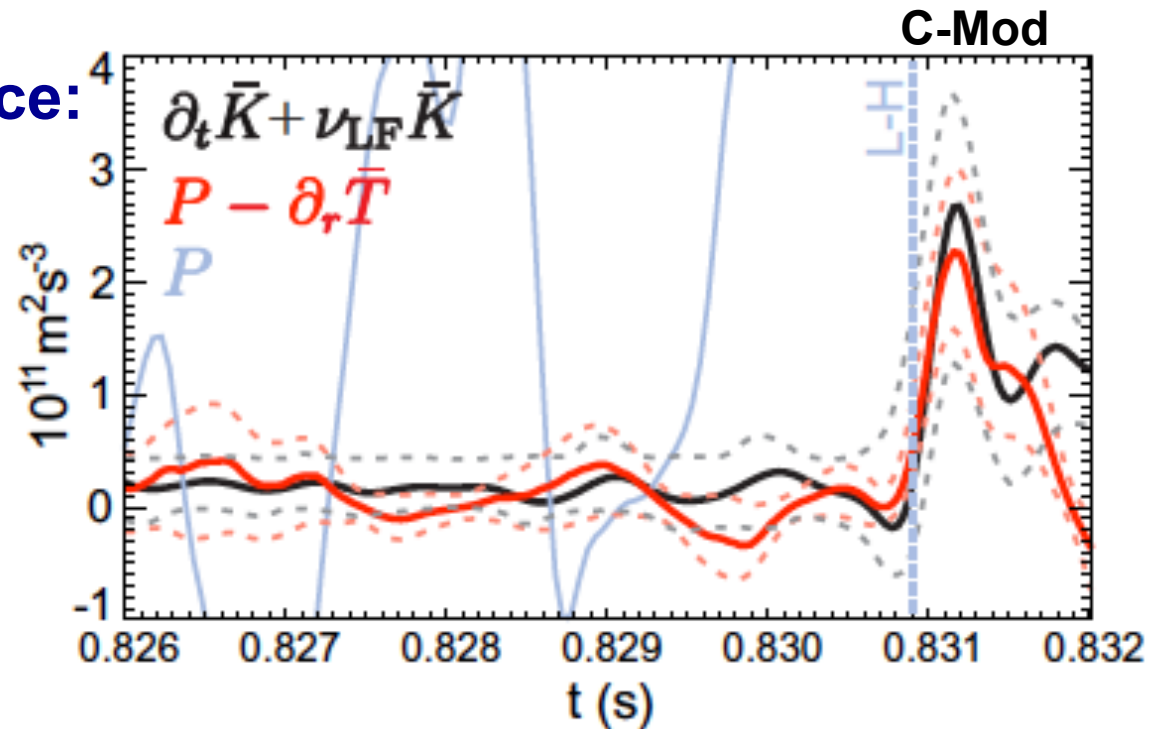
# Flow drive is consistent with observed flow

Cziegler, PPCF '14

## ExB Flow Power Balance:

$$\frac{\partial}{\partial t} \bar{K} + v_{ZF} \bar{K} = P - \partial_r \bar{T}$$

Flow Damping  
Determined  
From Steady-state  
L-mode Experimental  
Value



**Flow drive also consistent with estimated damping (including toroidal effects) in EAST (GS Xu NF '14)**



# Turbulence simulations now available...

---

- Chone' et al PoP'14
- Park et al PoP'15
- Nielsen et al, ArchivX, submitted '15

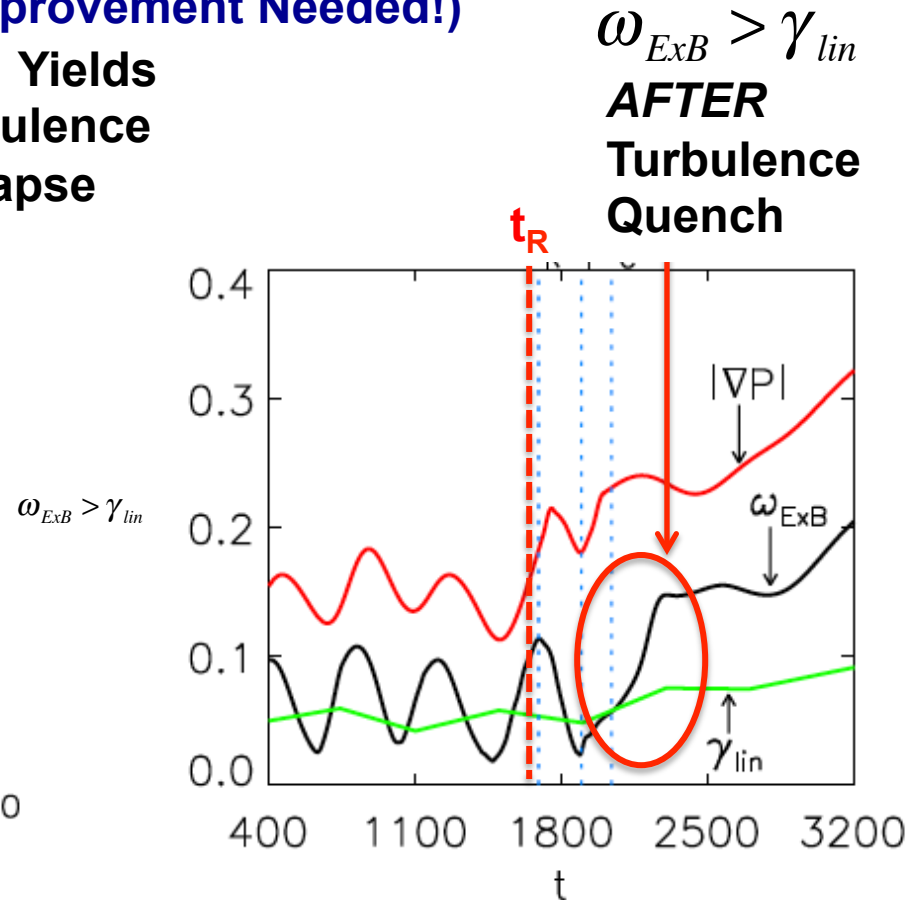
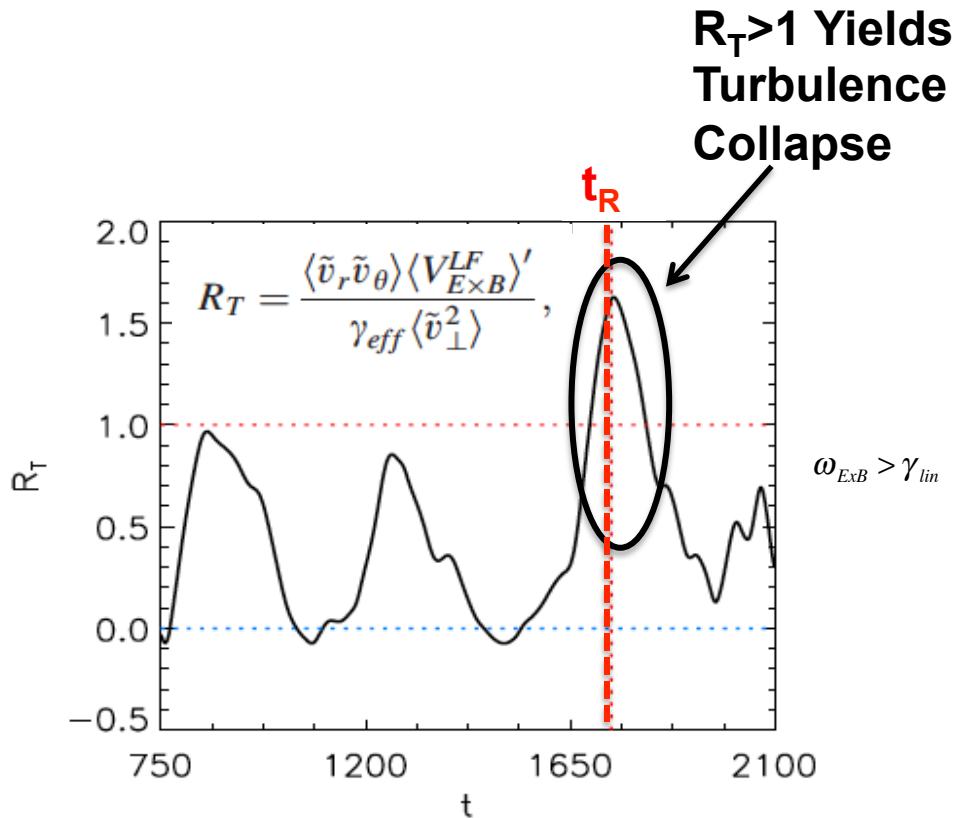
# Turbulence simulations now available...

---

- Chone' et al PoP'14
  - “The transition scenario from one regime to another is the following: increasing the input power leads to more violent avalanches, triggering strong stabilising ZF. The resulting steepening of the pressure gradient further generates a sheared mean flow via the neoclassical friction. If the shear becomes strong enough to prevent a new burst, the barrier locks on because of this positive feedback loop between pressure gradient and poloidal flow, mediated by neoclassical terms.”

# Turbulence simulations show same sequence

Park et al, PoP 2015, RBM Model (Improvement Needed!)



- **Chone' et al PoP'14** state that turbulent-driven ExB flow, together with mean flow, plays key role in edge barrier formation



# A picture of L-H Transition is Emerging

---

**SAME PICTURE IN**

**MULTIPLE EXPERIMENT**

**REDUCED MODELS, AND**

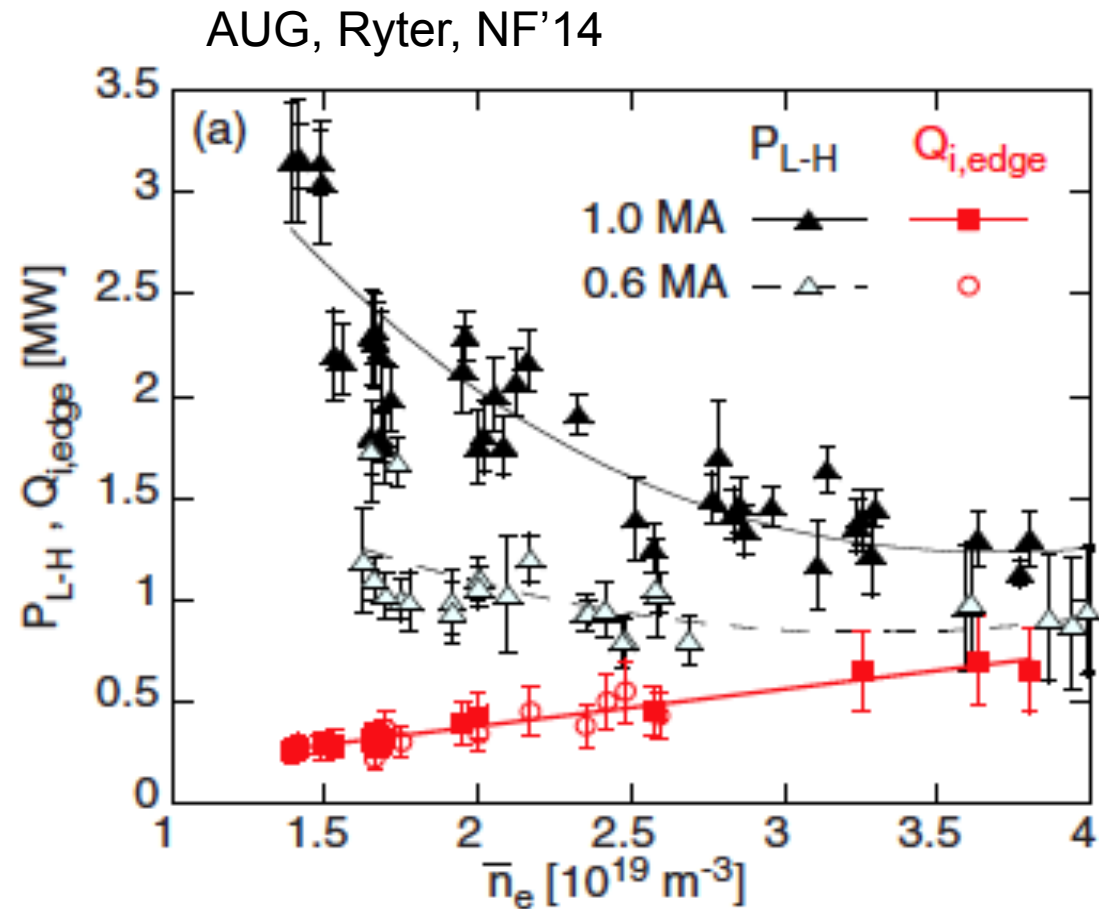
**TURBULENT SIMULATIONS**

---

# Linking Microscopic Turbulence Physics to Macroscopic System Behavior

- AUG Results Show Ion Heat Flux is Key
- Motivates Modified 1-D Predator-Prey Model (Malkov et al '15)

# L-H transition appears related ion heat flux at LCFS



# Extend Predator-prey to Separate e-ion Channels

Malkov, PoP'15

- Based on 1-D numerical 5-field model (Miki & Diamond 2012,13+)
- Currently operates on 6 fields (+ $P_e$ ) with self-consistently evolved transport coefficients
- Heat transport, + Two species with coupling, i,e:

$$\frac{\partial P_i}{\partial t} + \frac{1}{r} \frac{\partial}{\partial r} r \Gamma_i = \frac{2m}{M\tau} (P_e - P_i) + Q_i + \text{Anom coupl. to } T_e + \text{ZF dissip}$$

$$\frac{\partial P_e}{\partial t} + \frac{1}{r} \frac{\partial}{\partial r} r \Gamma_e = -\frac{2m}{M\tau} (P_e - P_i) + Q_e + \text{An cpl to } T_i \text{ (both thru turb)}$$

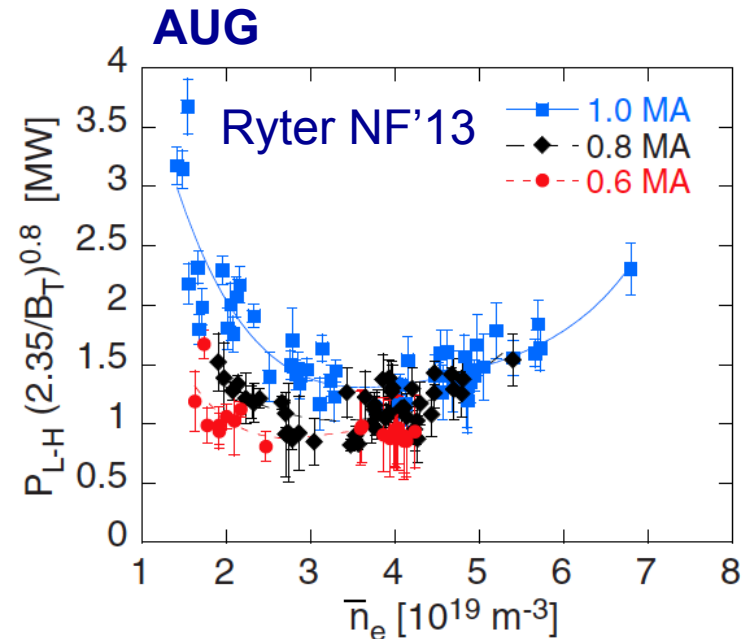
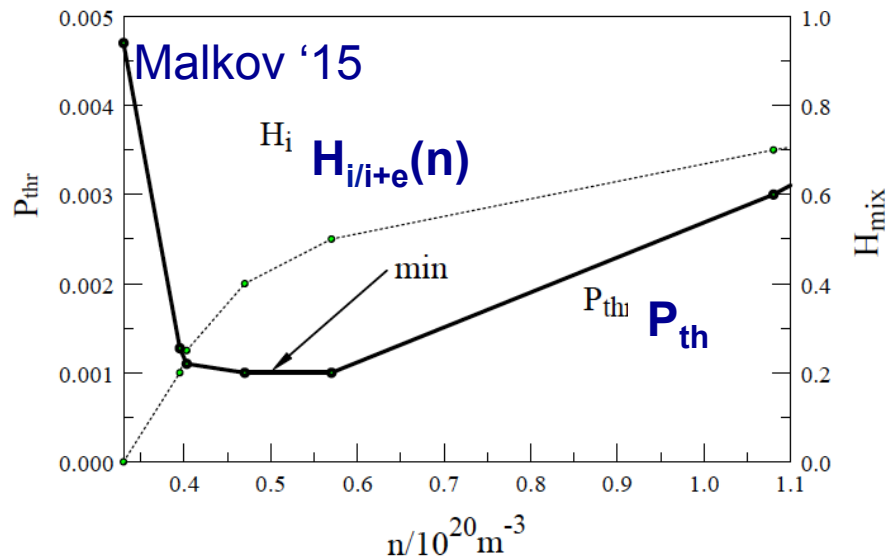
$$\Gamma = -(\chi_{neo} + \chi_t) \frac{\partial P}{\partial r}$$

- Density and the mean flow, as before

# Model with $H_{i/i+e}(n) \sim n$ gives minimum in $P_{th}(n)$

Key Parameter:  
e-i heating mix

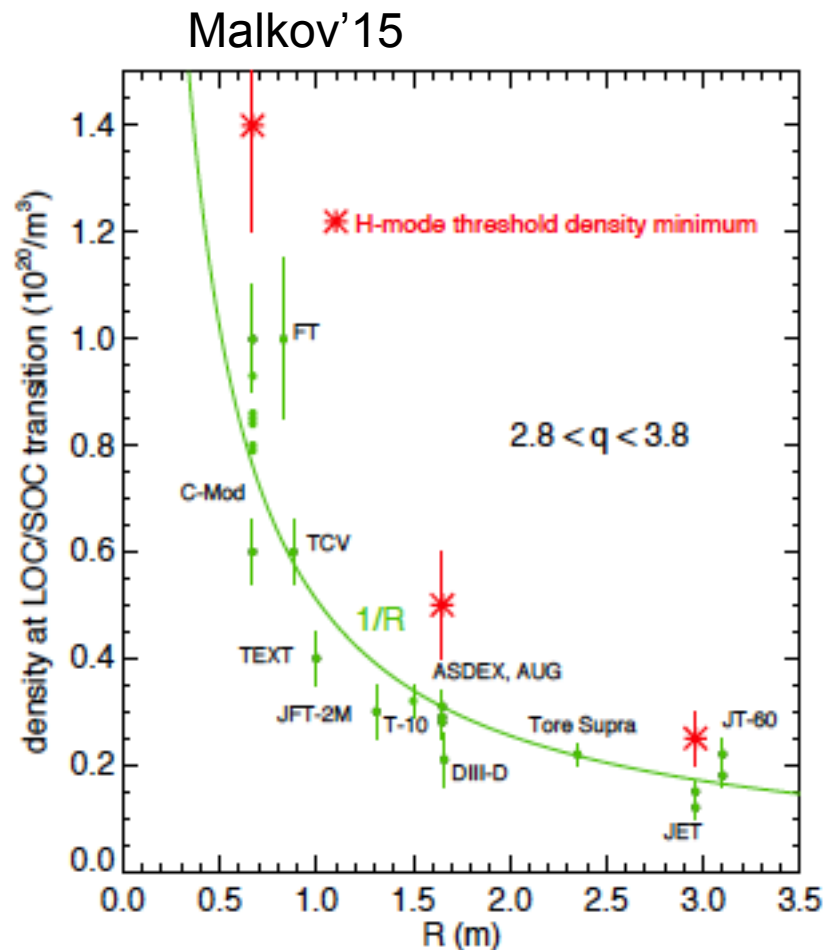
$$H_{i/(i+e)} = \frac{Q_i}{Q_i + Q_e}$$



- Result requires  $H_{i/i+e}(n) \sim n$  to yield minimum in  $P_{th}(n)$
- Normalized e-i equilibration time  $\tau_{e-i}^{coll} / \tau_{E_e}$  important
- Ion heat flux at LCFS is key; linked to transport across plasma column

(*n.b.* see also X. Wu et al, NF '15 for very recent similar work)

# Link between e-i equilibration and minimum in $P_{th}$ ?



- $P_{th}$  minimum correlated with LOC-SOC transition
- LOC-SOC due to increased e-i collisional coupling
- Suggests  $P_{th}$  minimum related to e-i coupling

# Summary of collisionally coupled model results

---

## Malkov PoP15

- $P_{\text{thr}}(n)$  grows monotonically in both pure ion ( $H_{i/i+e}=1$ ) and pure electron ( $H_{i/i+e}=0$ ) heating regimes with collisional coupling
- The descending low-density branch with distinct minimum, results from combined increased electron-to-ion collisional heat transfer and growing fraction  $H_{i/i+e}$  of heat deposited directly to ions
- Upturn in  $P_{\text{thr}}$  due to increase of shear flow damping (trapped-passing ions, neutrals,...)
- Finite heating mix,  $H_{i/i+e}$ , essential for core heat transport to build edge grad- $P_{\text{ion}}$
- Many mechanisms to give  $H_{i/i+e} > 0$

# Open Questions

---

- How to Reconcile w/ lack of ZF in AUG (Ryter, EPS'15)?
- Access to more attractive improved confinement regimes (e.g. I-mode) (Cmod, AUG)
- Significance (or lack thereof) of GAMs in turbulent-driven L-H transition model (Scott'04, Hallatschek'07, Kobayashi PRL'13, Itoh'14)
- H-mode Access in Collisionless (ITER) regime
- Origin of grad-B x B favorable power threshold (Fedorczak & Diamond'12)
- Effect of divertor configuration on access (Hughes TTF'15)
- Isotope effect (e.g. Ryter'13)
- Origin of co/cntr rotation effect on  $P_{th}$  (McKee'09)

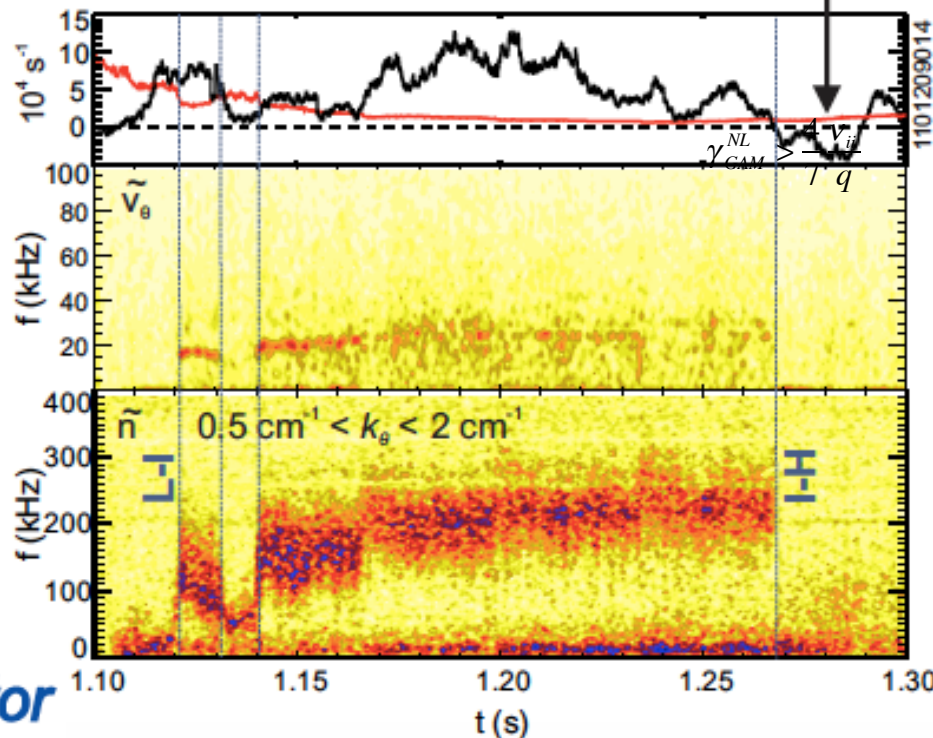


# Access to I-mode related to nonlinear GAM drive

Cziegler et al, PoP'13

Drive term compared to collisional damping rate from neoclassical theory:

$$\gamma = 4/7(\nu_{ii}/q)$$



$$\gamma_{NL}^{GAM} = \frac{\sum_{f_1} T_v(f_1, f)}{|v_{\perp}|^2(f)}$$

L-mode – I\_mode transition related to GAM drive exceeding GAM damping rate, i.e.

$$\gamma_{GAM}^{NL} > \frac{4}{7} \frac{\nu_{ii}}{q}$$

Reminiscent of L-H Transition!

# What about GAM damping effects on $V_\theta$ ?

---

Neoclassical Damping Rate given by:  $\nu_{ZF} \sim (1 + 2q^2)\nu_{ii}$

Results in a Reynolds Force That is too Small

To explain Observed  $V_\theta$  Transient; Kobayashi PRL'13, Itoh'14

## RESPONSE:

- ZFs grow/GAMs die as L-H is Approached (McKee', M.Xu PRL'12)
- EAST, HT-6M Show That Reynolds Stress **IS** Consistent w/ ExB Shear Flow & neoclassical Damping (GS Xu, IAEA'04, GS Xu, PRL'11)
- C-Mod shows that flow transient consistent with stress & empirically derived flow damping (Cziegler PPCF'14)

Further Quantitative Experiment Study, Models/  
Simulations with Toroidal Flow Needed to Resolve

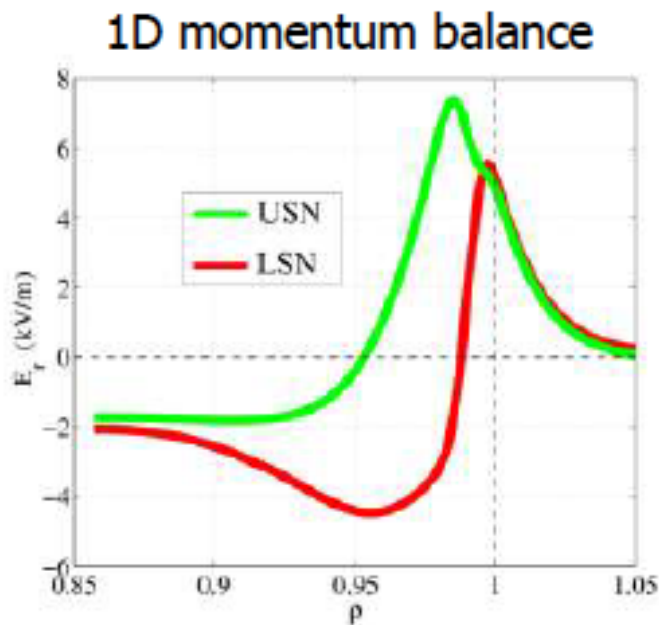
# Extending reduced model to collisionless regime (ITER relevant)

---

- Predator-prey model: ZF damping is collisional
  - Need ZF instability → turbulence → ion heating?
- Anomalous heat exchange important (Zhao&Diamond'13)
  - Heat exchange linked to fluctuation intensity-dependent coupling
  - Multiple pathways possible
    - CTEM →  $\langle J.E \rangle$  dissipation → ion heating
    - ITG-driven turbulent dissipation → ion heating
    - ZF dissipation → ion heating

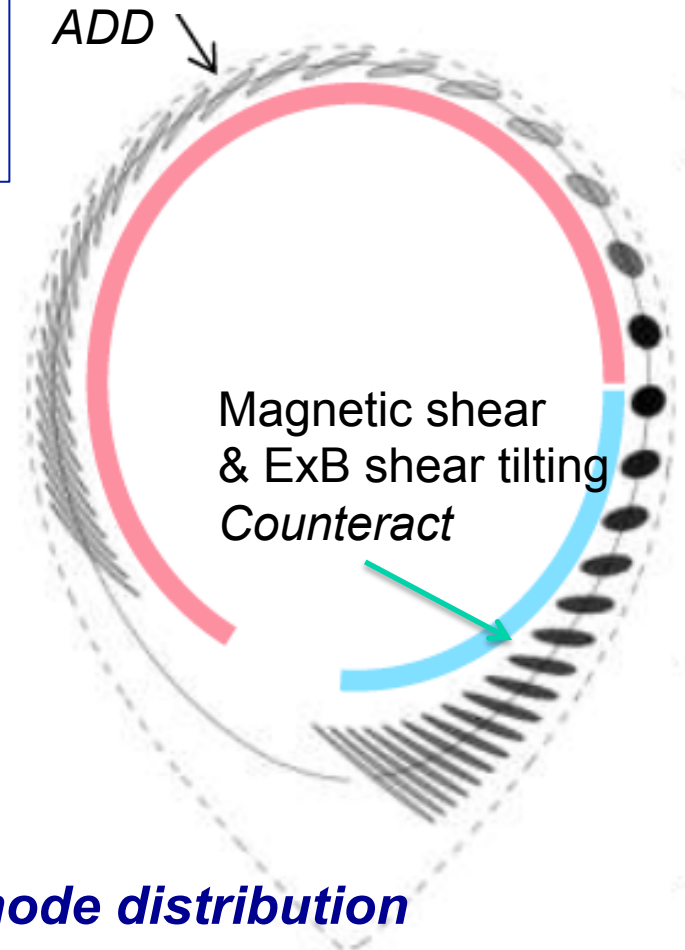
# Possible Origin of grad-BxB Effect

- Magnetic-shear induced tilting competes/ complements ExB shearing
- Up-down asymmetry determines dominant effect
- RESULT: Stronger sheared ExB in favorable configuration



Magnetic shear  
& ExB shear tilting

ADD



Magnetic shear  
& ExB shear tilting  
*Counteract*

**...But result depends on detailed poloidal mode distribution**

Fedorczak, Diamond et al '13

# Conclusions

---

- Microturbulence and reduced model studies point to origin of H-mode:
  - Increased heating leads to enhanced turbulent Reynolds stress driven sheared ExB flows
  - Transition initiated when rate of turbulent-ExB shear flow drive exceeds turbulent energy input rate
  - Results in turbulence collapse that allows edge pedestal to grow, locking in H-mode
- Turbulence simulations capture same picture
- Macroscale studies suggest ion heat flux at LCFS key
- Motivated modified predator-prey model that captures  $P_{\text{thr}}(n)$  dependence
- Many Open Issues: ITER Collisionless regime, grad-BxB effect, Isotope effect, role of GAMs, I-mode vs H-mode, ....