

# The field line map approach to plasma turbulence simulations

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für Plasmaphysik

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# Table of Contents

- 1 Introduction
- 2 Field line map
- 3 Application to simple turbulence model: Hasegawa-Wakatani
- 4 Outlook and summary

# Motivation

## Ultimate goal

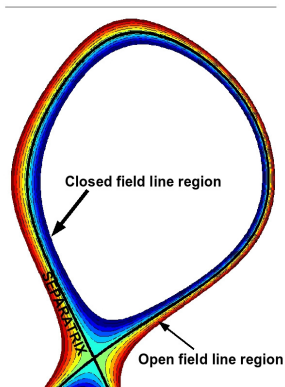
Prediction/Computation of anomalous transport/turbulence in edge and scrape-off layer

## Why?

- Boundary region may have high influence on overall performance of reactor [Stangeby90,McCracken93]
- Many phenomena occurring in edge/SOL not yet fully understood
- Prediction of heat loads on divertor plates

## Major challenges

- 1 Complex physical model
- 2 **Complex geometry**



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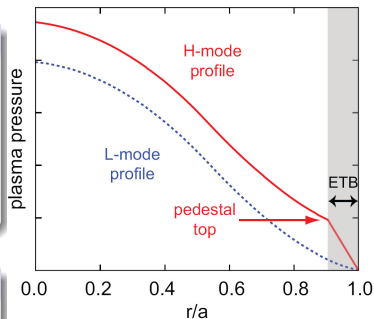
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- 1 Complex physical model
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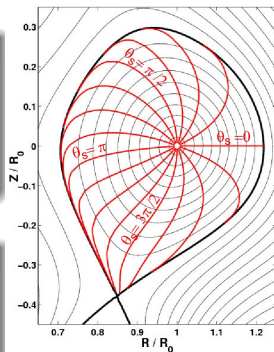
# Coordinates

## Field-aligned coordinates [D'haeseleer et al.90]

- Structures strongly elongated along field lines  $k_{\parallel} \ll k_{\perp}$
- Scale separation via transformation to field-aligned coordinates/grids
- **BUT: ILL DEFINED ON SEPARATRIX**

## Flux-aligned coordinates

- Ill defined on X/O-point [Mattor95], can be circumvented numerically
- However: X-point remains 'special', resolution imbalance



## Goal

- Development and implementation of numerical concept, which avoids field/flux-aligned coordinates  $\rightarrow$  applicable to separatrix
- Application to simplified plasma turbulence model

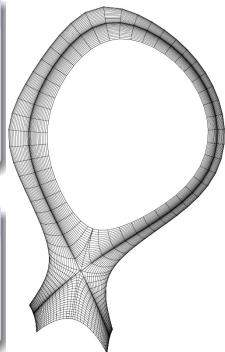
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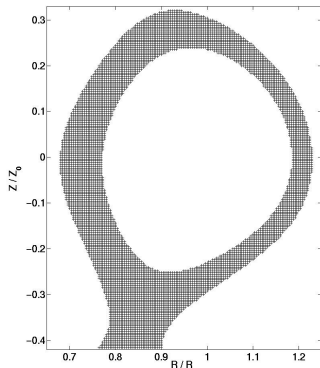
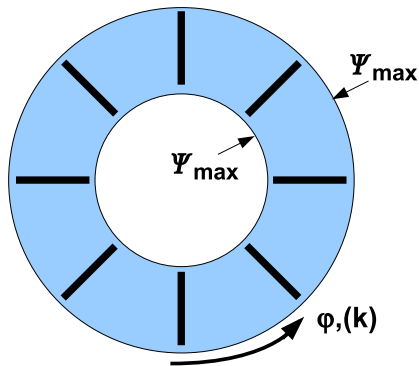
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## Overview

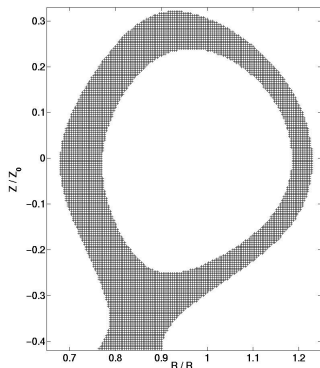
- See Flux Coordinate Independent approach (FCI) [Ottaviani11, Hariri and Ottaviani13, Hariri et al.14], F. Hariri's talk on Tuesday
- Cylindrical grid  $R_i, Z_j, \varphi_k \rightarrow$  no singularities<sup>1</sup> or special points
- Fieldline-following discretisation for parallel operators
- Grid sparsification in toroidal direction ( $k_{\parallel} \ll k_{\perp}$ )



<sup>1</sup>in relevant region



## Perpendicular operators



Assumption  $B_{tor} \gg B_{pol}$

→ Stencil remains within poloidal plane

→ Use of standard finite-difference methods

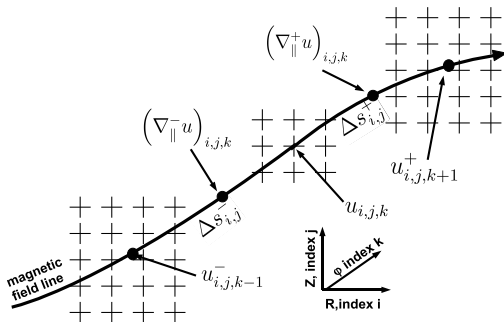
## Parallel operators

- Stencil covers neighbouring poloidal planes
- Discretisation via finite difference along magnetic field lines
- Fieldline tracing and interpolation

$$R^\pm = \int_0^{\pm\Delta\varphi} \frac{B^R}{B^\varphi} d\varphi, \quad Z^\pm = \int_0^{\pm\Delta\varphi} \frac{B^Z}{B^\varphi} d\varphi,$$

$$\left(\nabla_{\parallel}^\pm u\right)_{i,j,k} = \pm \frac{u_{i,j,k\pm 1}^\pm - u_{i,j,k}^\pm}{\Delta S_{i,j}^\pm},$$

$u_{i,j,k}^\pm$  via 2D interpolation



- Express via matrices  $\mathbf{Q}^\pm$ , i.e.:

$$\mathbf{q}^\pm = \mathbf{Q}^\pm \mathbf{u}$$

$$\mathbf{q}^\pm := \left( \left( \nabla_{\parallel}^\pm u \right)_{1,1,1}, \left( \nabla_{\parallel}^\pm u \right)_{2,1,1}, \dots \right), \quad \mathbf{u} := (u_{1,1,1}, u_{2,1,1}, \dots)$$

## Parallel diffusion

$$\mathcal{D}_{\parallel} u := \nabla \cdot [\mathbf{b} (\nabla_{\parallel} u)]$$

### Naive scheme

- Assume:  $\mathcal{D}_{\parallel} \approx \nabla_{\parallel}^2$
- Discretisation via further finite difference along magnetic field lines:

$$\left( \mathcal{D}_{\parallel}^{\text{naive}} u \right)_{i,j,k} = \frac{2}{\Delta s_{i,j}^+ + \Delta s_{i,j}^-} \left[ \nabla_{\parallel}^+ u_{i,j,k} - \nabla_{\parallel}^- u_{i,j,k} \right]$$

- However: Better scheme possible

## Parallel diffusion: Support operator method

- Motivated from [Günter et al. (2005 and 2007)]
- Construct scheme which mimics 'good' property on discrete level, i.e. self-adjointness ( $u, v = 0$  at boundaries):

$$\begin{aligned}\langle u, \mathcal{D}_{\parallel} v \rangle &= \int_V u \nabla \cdot [\mathbf{b}(\nabla_{\parallel} v)] dV = - \int_V \nabla_{\parallel} u \nabla_{\parallel} v dV \\ \rightarrow \nabla_{\parallel}^{\dagger} &= -\nabla \cdot [\mathbf{b} \circ], \quad \mathcal{D}_{\parallel}^{\dagger} = \mathcal{D}_{\parallel}\end{aligned}\tag{1}$$

### Support operator method [Shaskov (1996)]

Mimic integral equality 1 on discrete level

- 1 Define discrete space for scalars  $SG$  and fluxes (gradients)  $FG^{\pm}$  with corresponding inner product
- 2 Define and discretise prime operator:  $\nabla_{\parallel} \rightarrow \mathbf{Q}^{\pm} : SG \rightarrow FG^{\pm}$
- 3 Use integral equality 1 to construct derived operator

## Parallel diffusion: Support operator method

Discrete spaces:

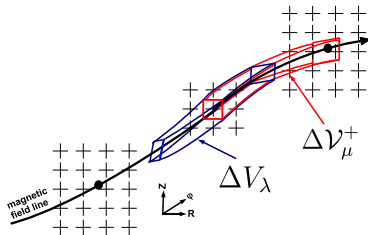
- Scalars  $SG$  : cylindrical grid
- Fluxes  $FG^\pm$  : staggered along field lines
- Inner products:

$$\langle u, v \rangle_{SG} := \sum_{\lambda} u_{\lambda} v_{\lambda} \Delta V_{\lambda}$$

$$\langle q^{\pm}, p^{\pm} \rangle_{FG^{\pm}} := \sum_{\mu} q_{\mu}^{\pm} p_{\mu}^{\pm} \Delta V_{\mu}$$

Greek indices denote summation over all grid points

- Two possibilities ' $\pm$ ' (will be combined later)



## Parallel diffusion: Support operator method

Discrete parallel gradient, discrete inner products, integral equality



Discrete parallel diffusion operator

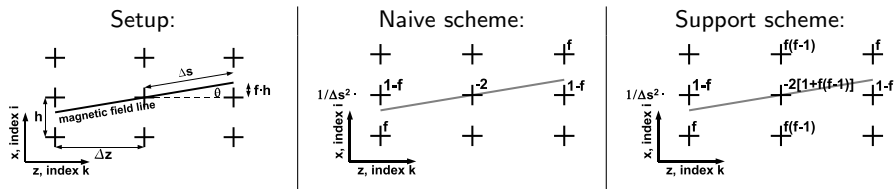
$$\mathbf{D}_{\parallel, \sigma, \lambda}^{\pm, \text{supp}} = - \sum_{\mu} \mathbf{Q}_{\mu, \lambda}^{\pm} \mathbf{Q}_{\mu, \sigma}^{\pm} \frac{\Delta V_{\mu}^{\pm}}{\Delta V_{\lambda}}$$

### Notes

- if volumes equal:  $\mathbf{D}_{\parallel}^{\pm, \text{supp}} = - (\mathbf{Q}^{\pm})^T \mathbf{Q}^{\pm}$
- Fusion of ' $\pm$ ' choice:  $\mathbf{D}_{\parallel}^{\text{supp}} = \frac{1}{2} \left( \mathbf{D}_{\parallel}^{+, \text{supp}} + \mathbf{D}_{\parallel}^{-, \text{supp}} \right)$
- Conserves L1-norm (energy)
- Decreases L2-norm  $\rightarrow$  numerically stable regardless of interpolation method

## Simple 2D model problem

- $x \sim$  coordinate within poloidal plane,  $z \sim$  toroidal coordinate
- Uniform magnetic field inclined by  $\tan \theta = \frac{f \cdot h}{\Delta z}$  with respect to grid
- Linear interpolation along  $x$  (1d)



- Stencil for support scheme larger
- For  $f = 0, 1$ , i.e. no displacement, both schemes yield standard second order finite-difference expression

## Numerical diffusion

### Problem

- Interpolation introduces erroneous numerical coupling among distinct field lines  
→ Leads to numerical perpendicular 'diffusion', which is dependent on resolution
- Might overwhelm real (slow) perpendicular dynamics  $\chi_{\parallel}/\chi_{\perp} \sim 1 \cdot 10^{10}$
- Extremely high resolutions might be needed

### Scaling of numerical diffusion

Action of discrete parallel diffusion operator on mode  $u = \exp(ik_x x + ik_z z)$

$$\mathbf{D}_{\parallel}^{naive} u = \left[ -k_{\parallel}^2 - \frac{f(1-f)(k_x h)^2}{\Delta s^2} + \mathcal{O}\left(\frac{(k_x h, k_z \Delta z)^4}{\Delta s^2}\right) \right] u$$

$$\mathbf{D}_{\parallel}^{supp} u = \left[ -k_{\parallel}^2 + \mathcal{O}\left(\frac{(k_x h, k_z \Delta z)^4}{\Delta s^2}\right) \right] u$$

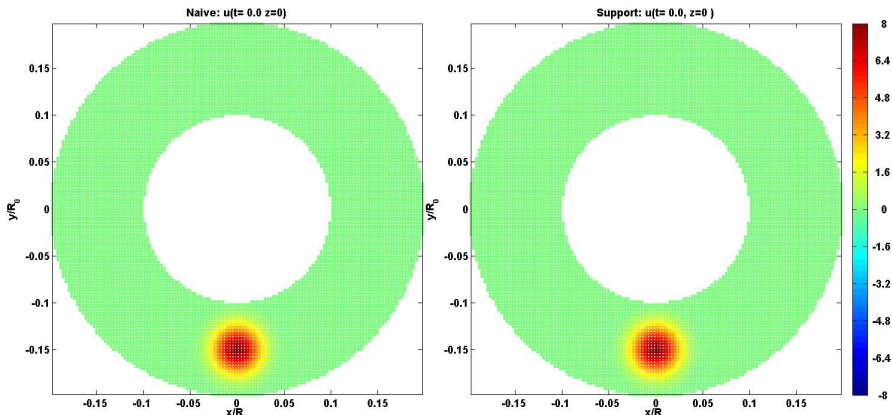
**Support scheme exhibits better scaling → numerical diffusion drastically reduced**



# Numerical diffusion: Example

## Example

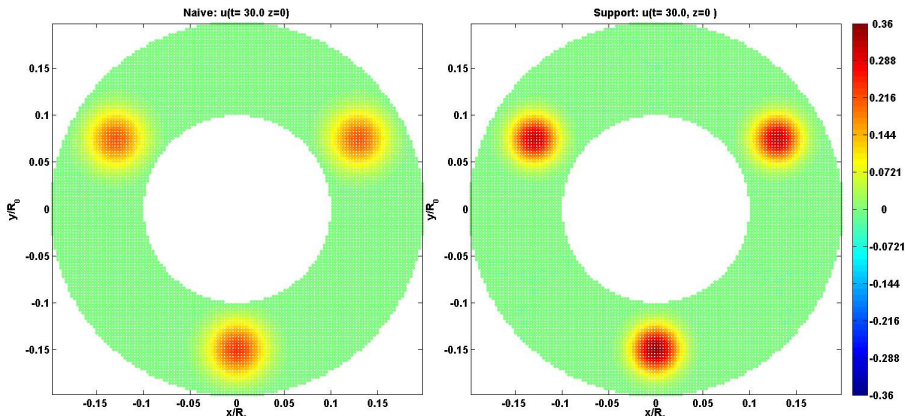
- Parallel diffusion equation  $\partial_t u = \chi_{\parallel} \mathcal{D}_{\parallel} u$
- Axial circular configuration with  $q = 3$
- Initial state:  $u(t = 0, x, y, z) = \exp\left(-\frac{(x-x_0)^2 + (y-y_0)^2}{\sigma^2}\right) \cdot \delta(z)$
- resolution:  $N_z = 8$ ,  $\frac{\sigma}{h} \approx 8.3$



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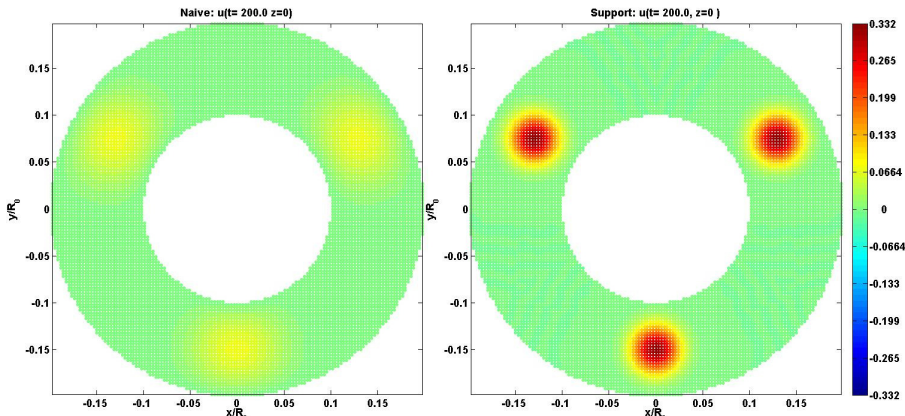
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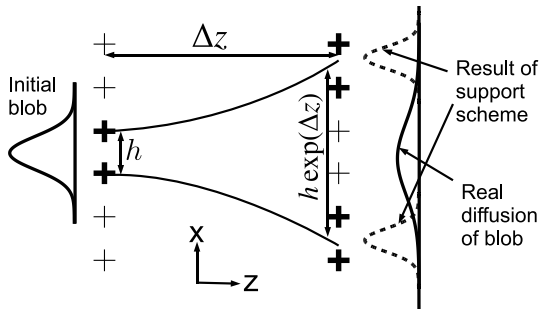
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## Map distortion

### Magnetic field lines around X-point

- Lowest order expansion:  $\mathbf{B} = B_0 [\mathbf{e}_z + \alpha (x\mathbf{e}_x - y\mathbf{e}_y)]$
- Distance between field lines diverges exponentially, i.e.  $\delta(z) \propto \exp \alpha z$

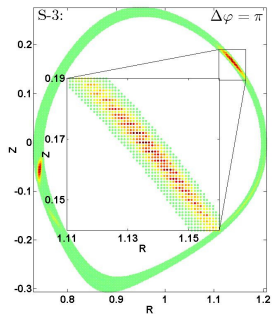


- With support scheme information is not only 'taken' but also 'sent' to neighbouring planes
- If points not properly connected erroneous wiggles may arise

## Map distortion: Example

Result of diffusion of blob in flux surfaces close to separatrix:

Support scheme  
Interpolation  
 $N_\varphi = 2$



# Map distortion: Remedy 1

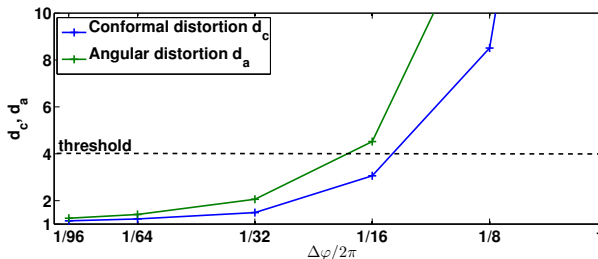
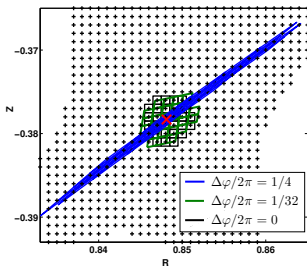
- Quantify distortion via mapped quads:

$$d_c = \max_{i,j} \frac{\text{longest side of mapped quad } i,j}{\text{shortest side of mapped quad } i,j},$$

$$d_a = \max_{i,j} \frac{\text{largest angle of mapped quad } i,j}{\text{smallest angle of mapped quad } i,j}.$$

- Require enough toroidal resolution, such that map distortion remains below threshold

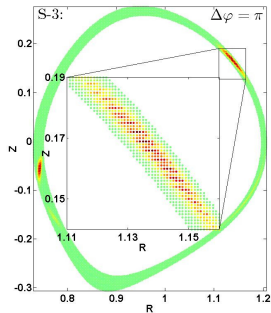
$$d_c, d_a \leq 4$$



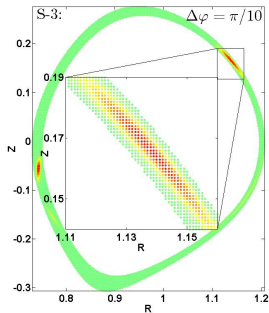
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Result of diffusion of blob in flux surfaces close to separatrix:

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Interpolation  
 $N_\varphi = 2$



Support scheme  
Interpolation  
 $N_\varphi = 20$



Note: Tiny oscillations on grid scale might still arise due to change of interpolation stamp [*Held15 et al., submitted*], could be cured by small amount of high order perpendicular dissipation.

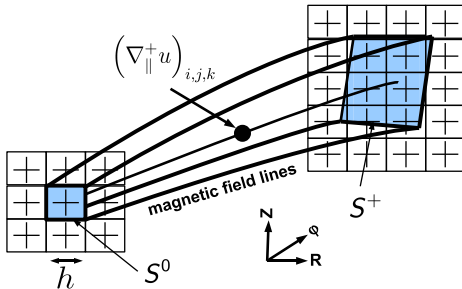
## Map distortion: Remedy 2

Change parallel gradient to account for distorted field lines:

### Coordinate free representation (Integration)

$$\nabla_{\parallel} u = \frac{1}{B} \nabla \cdot (u \mathbf{B}) = \frac{1}{B} \lim_{V \rightarrow 0} \frac{1}{V} \int_{\partial V} u \mathbf{B} \cdot d\mathbf{S}.$$

- Mimic surface integral on discrete level
- Integration over toroidal ends of flux box  $\nabla_{\parallel}^{\pm} u = \pm \frac{1}{B_V \Delta V} \left[ \int_{S^{\pm}} u B_{\text{tor}} dS - \int_{S^0} u B_{\text{tor}} dS \right]$
- Combination with interpolation

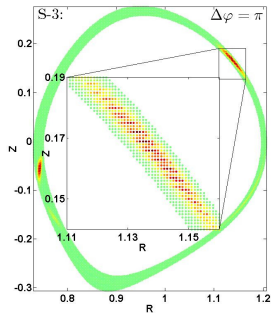




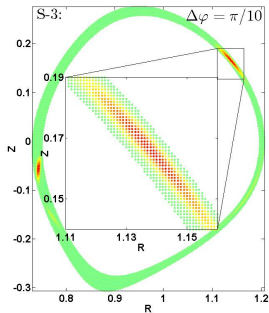
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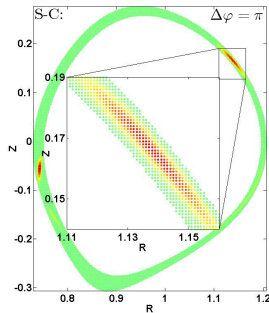
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Interpolation  
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Support scheme  
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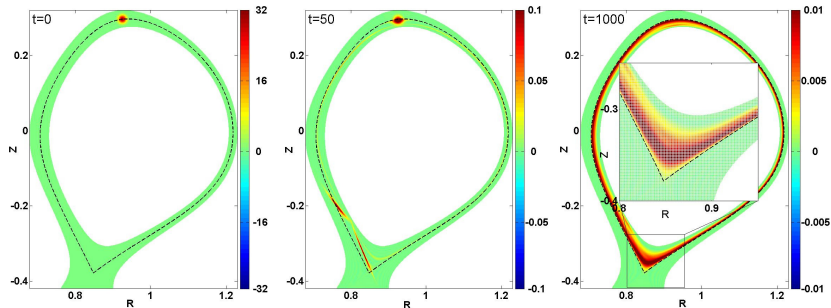
Support scheme  
Integration  
 $N_\varphi = 2$



## Application to separatrix

### Example

- Initial state: Gaussian( $R, Z$ )-delta( $\varphi$ ) blob located on separatrix
- resolution:  $N_z = 16$ ,  $\frac{\sigma}{h} = 10$
- Boundary conditions:  $u = 0$  at divertor plates



- Improvement of parallel boundary conditions is ongoing (e.g. Neumann conditions)

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## Hasegawa-Wakatani equations

Normalised equations [*Hasegawa and Wakatani83*]:

Simple self-consistent 3D turbulence model (resistive drift-wave turbulence)

$$\left( \frac{\partial}{\partial t} + \mathbf{v}_E \cdot \nabla \right) n = w_n \mathbf{v}_E \cdot \nabla \Psi(x_\perp) + \nu_n \nabla_\perp^6 n + \sigma \nabla \cdot [\mathbf{b} \nabla_\parallel (n - \phi)],$$
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- $\sigma$  : parallel conductivity, Ohm's law:  $\sigma J_\parallel = \nabla_\parallel (n - \phi)$



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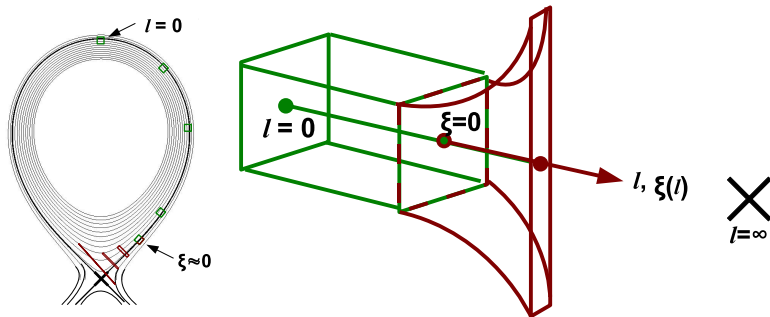
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- $\nu_n, \nu_\phi$  : viscosity coefficients
- $\sigma$  : parallel conductivity, Ohm's law:  $\sigma J_\parallel = \nabla_\parallel (n - \phi)$

## GRILLIX

- Model implemented for axisymmetric geometry with arbitrary poloidal cross section
- MPI + OpenMP parallelised

## X-Point: Theoretical background

Flux box around reference magnetic field line approaching X-point

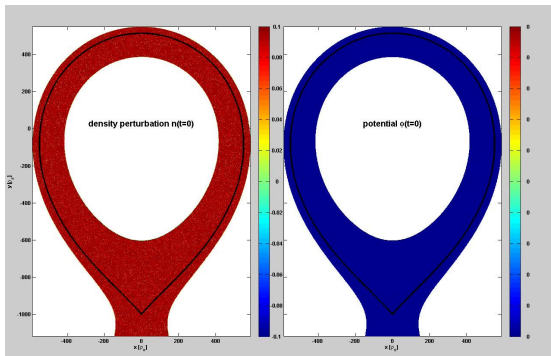


### Basic picture [Myra et al.]

- Strong distortion of field-aligned structures towards X-point
- Drastic increase of  $k_{\perp}$  towards X-point:  $k_{\perp}(l) \xrightarrow[\xi \gg \alpha^{-1}]{} k_{\perp 0} \exp(\alpha \xi(l))$   
→ Operators with highest  $k_{\perp}$  dependence dominant near X-point, i.e. dissipation
- X-point tends to disconnect structures → increase of  $k_{\parallel}$

# Axial diverted geometry<sup>2</sup>

Initial state

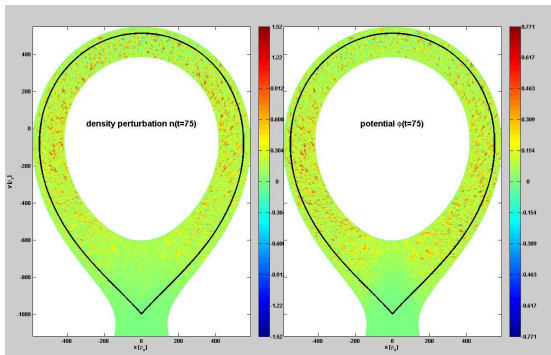


<sup>2</sup>Parameters reflect roughly:

$$T_e = 80\text{eV} \quad n_{e0} = 4.5 \cdot 10^{13} \text{cm}^{-3}, \quad B = 2.5\text{T}, \quad R_0 = 165\text{cm}, \quad a = 30\text{cm}, \quad L_n = 3.65\text{cm}, \quad M_i = 3670m_e, \quad q_{95} \approx 3/2$$

# Axial diverted geometry<sup>2</sup>

## Linear phase



## Observations

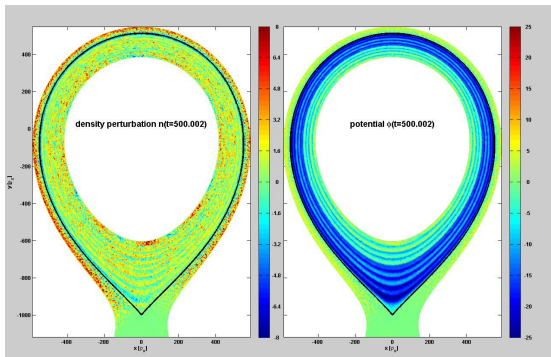
- Strongest effective drive at top
- Fluctuations towards sides driven by parallel currents
- No Fluctuations near X-point

<sup>2</sup>Parameters reflect roughly:

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# Axial diverted geometry<sup>2</sup>

Turbulent phase



## Observations

- Automatic development of flux aligned structures
- Clear and sharp separation between open and closed flux surfaces visible

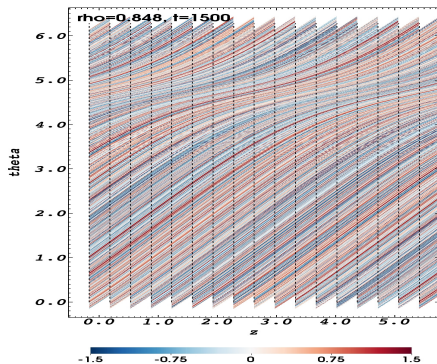
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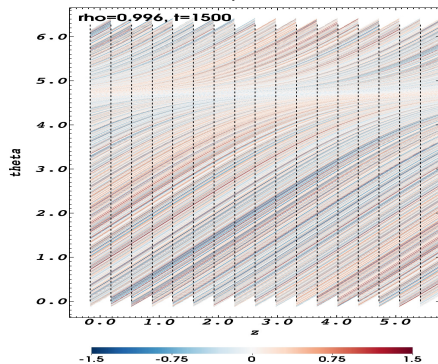
## View along magnetic field lines:

Nonadiabaticity  $n - \phi$  on flux surface...

...well inside:



...near separatrix



## Observations

- structures are numerically smooth along field lines
- Fluctuations damped near X-Point  $\rightarrow$  X-point disconnects structures  $\rightarrow$  increase of  $k_{\parallel}$
- Similar to resistive X-point mode [Myra et al.00]

# Table of Contents

- 1 Introduction
- 2 Field line map
- 3 Application to simple turbulence model: Hasegawa-Wakatani
- 4 Outlook and summary

## Recent developments and Outlook

- Extension to full-f:  
geometric multigrid solver for nonlinear polarisation ✓

$$\nabla \cdot (n \nabla_{\perp} \phi) = rhs$$

- Additional moments (parallel transport model ✓, thermal fluctuations) towards electromagnetic drift reduced Braginskii model
- Realistic boundary conditions (ongoing)
- Experimental equilibria (eqdsk) ✓
- Extension to 3D geometries, i.e. stellarators (maybe)
- Study of blob propagation in realistic geometry (ongoing → *M.Siccinio*)
- Cross verification against similar codes (GBS, FENICIA, TOKAM3D, FELTOR, BOUT++)
- Investigate effect of geometry on turbulent transport and heat exhaust, relevant for ITER and DEMO



## Summary

### Field line map

- Cylindrical/Cartesian grid, sparsified in toroidal/axial direction
- FCI: Not based on field/flux aligned coordinates → applicable to separatrix, X/O-points
- Discretisation of perp. operators straight forward (2nd order finite differences)
- Parallel operators via field line tracing and interpolation/integration
- Reduction of numerical diffusion via self-adjoint discretisation with support operator method
- Effects of map distortion identified and resolved (recommendation: use integration-interpolation)

### Application to simple turbulence model

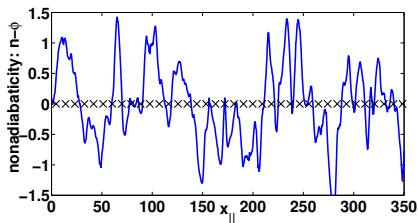
- Hasegawa-Wakatani model implemented in parallel code GRILLIX
- Automatic development of field aligned and zonal structures (with Cartesian grid!!)
- Application to X-point geometry:
  - ▶ Strong shear enhances dissipation and leads to disconnection of structures
  - ▶ Increase of  $k_{\parallel}$
  - ▶ Confirmed with GRILLIX

# Backup slides

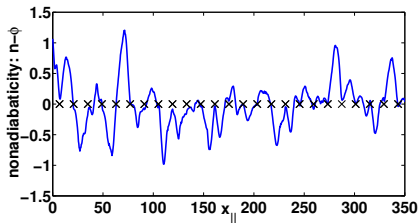
## View along magnetic field lines:

...along sample magnetic field line running in flux surfaces...

...well inside:



...near separatrix



crosses (x) denote positions where field line passes bottom, i.e. where field line has its closest approaches to X-point

### Observations

- Fluctuations strongly damped near X-point
- X-point disconnects structures  $\rightarrow$  increase of  $k_{\parallel}$

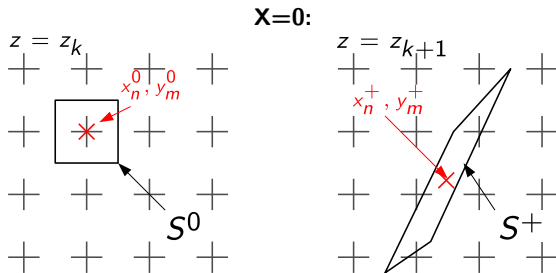
$\rightarrow$  Similar to resistive X-point mode [Myra et al.00], though different physical model  $\rightarrow$  Generic property of X-point

## Integration-Interpolation method (detailed)

$$\nabla_{\parallel} u = \frac{1}{B} \lim_{V \rightarrow 0} \frac{1}{V} \int u \mathbf{B} \cdot d\mathbf{S} \implies \mathbf{Q}^+ u = \frac{1}{B \Delta V^+} \underbrace{\left[ \int_{S^+} u B_{\text{tor}} dS_{\text{tor}} - \int_{S^0} u B_{\text{tor}} dS_{\text{tor}} \right]}_{\text{mimic}}$$

$$\int_{S^0} u B_{\text{tor}} dS_{\text{tor}} = \sum_{n,m=1}^{2^X} u(x_n^0, y_m^0, z = z_k) B_{\text{tor}}(x_n^0, y_m^0) \frac{h^2}{2^{2X}}, \quad \int_{S^+} u B_{\text{tor}} dS_{\text{tor}} = \sum_{n,m=1}^{2^X} u(x_n^+, y_m^+, z = z_{k+1}) B_{\text{tor}}(x_n^+, y_m^+) \Delta S^+,$$

$$u(x_n^0, y_m^0, z_k), u(x_n^+, y_m^+, z_{k+1}) \text{ via interpolation, } B_{\text{tor}} \text{ analytically available, } \Delta S^+ = \frac{h^2}{2^{2X}} \frac{B_{\text{tor}}(x_n^0, y_m^0)}{B_{\text{tor}}(x_n^+, y_m^+)}$$



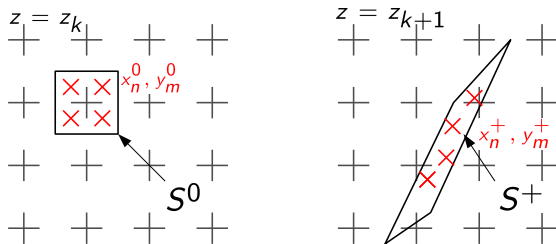
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**X=1:**



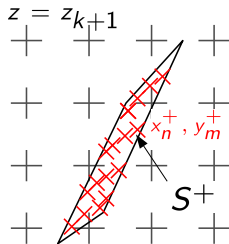
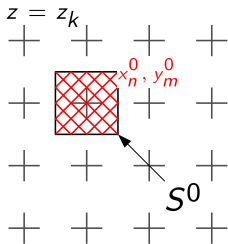
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**X=2:**



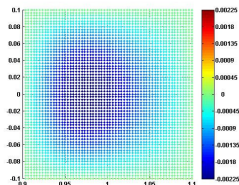
# Consistency of parallel diffusion operator in general geometry

Considered case [Held15]:

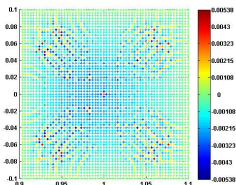
$$\Psi(R, Z) = \cos\left(\frac{\pi}{2}(R - R_0)\right) \cos\left(\frac{\pi}{2}Z\right), \quad u(R, Z, \varphi) = -\Psi(R, Z) \cos(\varphi)$$

Difference to analytic result:

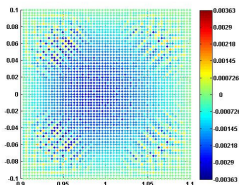
**D-3**



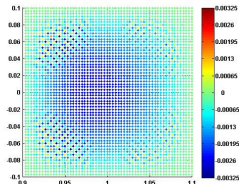
**S-3X0**



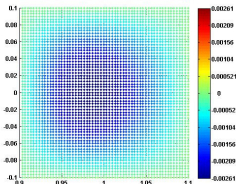
**S-3X1**



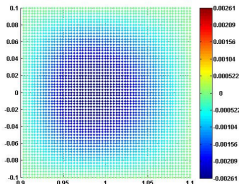
**S-3X2**



**S-3X4**



**S-3X6**

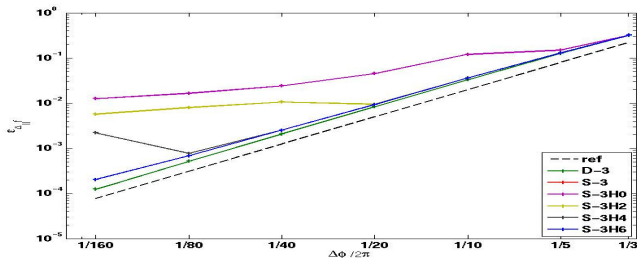


## Consistency of parallel diffusion operator in general geometry

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Quantify numerical error  $\epsilon_{\Delta_{\parallel} u} := \frac{|(\mathcal{D}_{\parallel} u)_{\text{numeric}} - (\mathcal{D}_{\parallel} u)_{\text{analytic}}|_2}{|(\mathcal{D}_{\parallel} u)_{\text{analytic}}|_2}$



- Second order accuracy in toroidal direction can be achieved
- Prerequisite: sufficiently high perp. resolution, sufficiently high  $X$  to account for map distortion

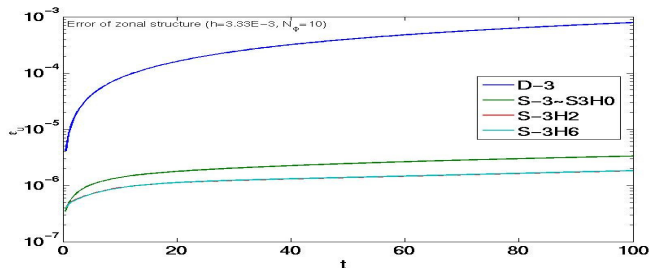


## Numerical diffusion in general geometry

Considered case [Held15]:

$$\Psi(R, Z) = \cos\left(\frac{\pi}{2}(R - R_0)\right) \cos\left(\frac{\pi}{2}Z\right), \quad u(R, Z, \varphi, t = 0) = 0.1\Psi(R, Z)^2$$

Measure numerical diffusion:  $\epsilon_u(t) = \frac{|u(t) - u(t=0)|_2}{|u(t=0)|_2}$

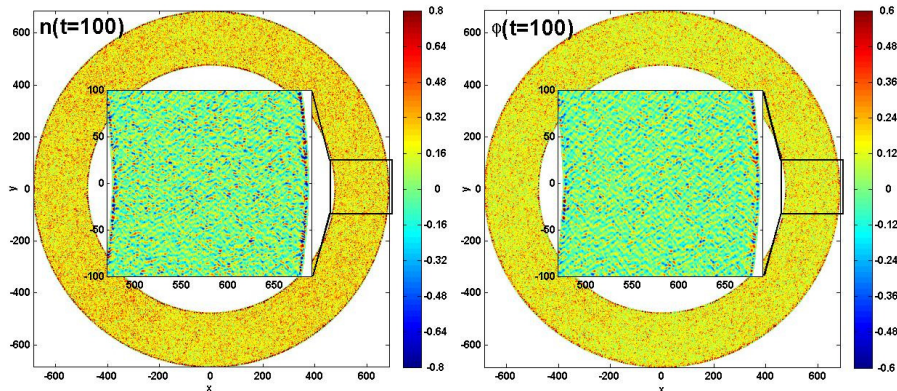


- Lower numerical diffusion with support scheme compared to naive scheme also in general geometries

# Axial circular geometry

## Demonstration simulation

$$T_e = 80\text{eV}, \quad n_{e0} = 4.5 \cdot 10^{13}\text{cm}^{-3}, \quad B = 2.5\text{T}, \quad R_0 = 165\text{cm},$$
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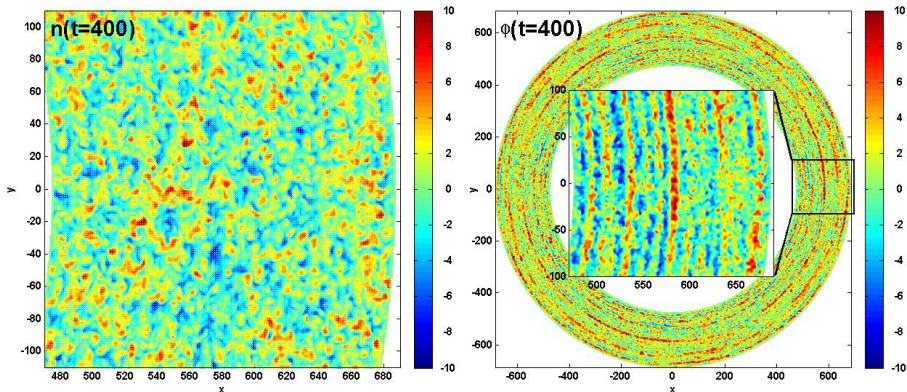


Linear phase: poloidally propagating drift waves

# Axial circular geometry

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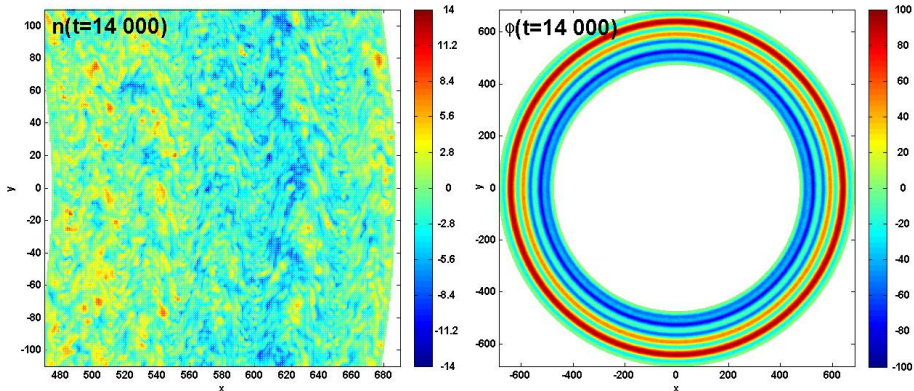


Onset of turbulent phase: Eddy mitosis

# Axial circular geometry

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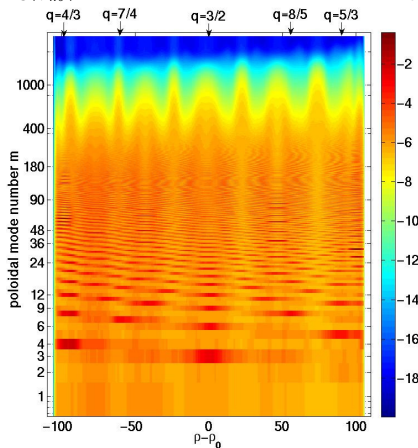


Saturated phase: Zonal flow

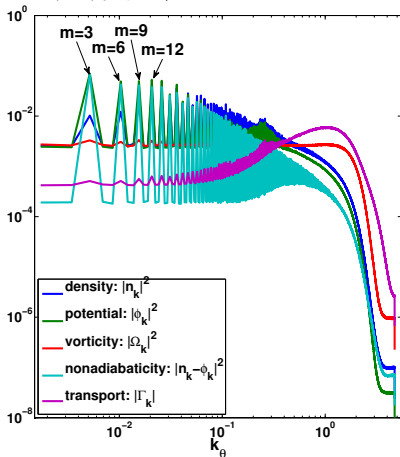
# Axial circular geometry

Poloidal spectra:

$$\log |\phi_{k\theta}|^2$$



spectra at  $\rho = \rho_0$ ,  $q = 3/2$



$k_{\parallel} \approx 0$  structures on rational surfaces in accordance with geometry

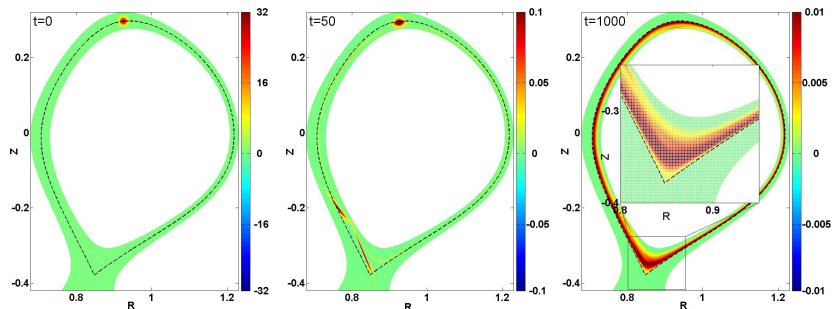
# Toroidal Geometry

## GRILLIX (A.Stegmeir, see talk on Thursday)

- FCI applied to toroidal geometries
- Discretisation of parallel diffusion
- Based on integral representation for parallel gradient to cope with map distortion

$$\nabla_{\parallel} u = \frac{1}{B} \nabla \cdot (u\mathbf{B}) = \frac{1}{B} \lim_{V \rightarrow 0} \frac{1}{V} \int_{\partial V} u\mathbf{B} \cdot d\mathbf{S}.$$

*Simulation of temperature blob in realistic toroidal geometry (parallel diffusion):*



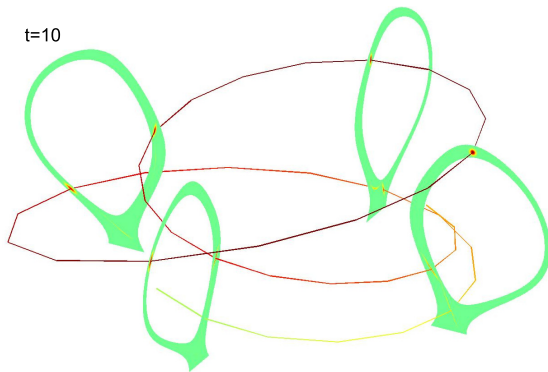
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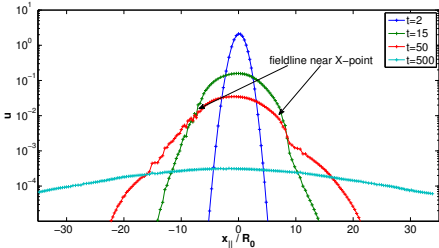
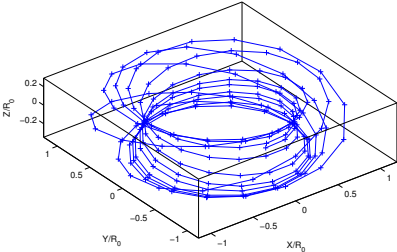
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*Simulation of temperature blob in realistic toroidal geometry (parallel diffusion):*



# Structure along sample magnetic field line



of Flux surface average:

