

Gaussian Radial-Basis-Function solution of the non-linear Fokker-Planck equation

Eero Hirvijoki

Jeff Candy

Emily Belli

Ola Embréus

European Fusion Theory Conference

Lissabon 2015



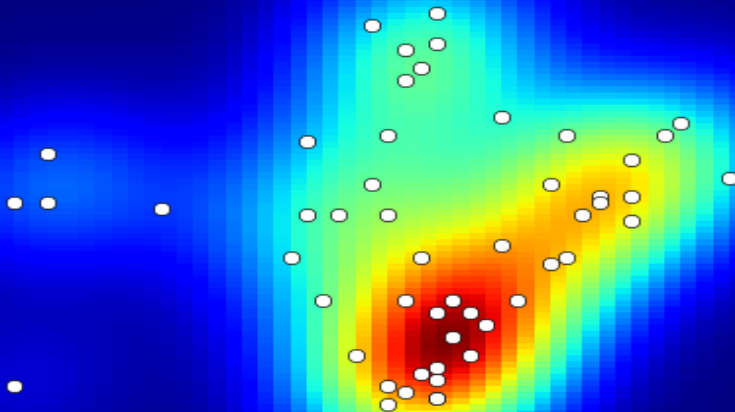
CHALMERS
UNIVERSITY OF TECHNOLOGY



Outline

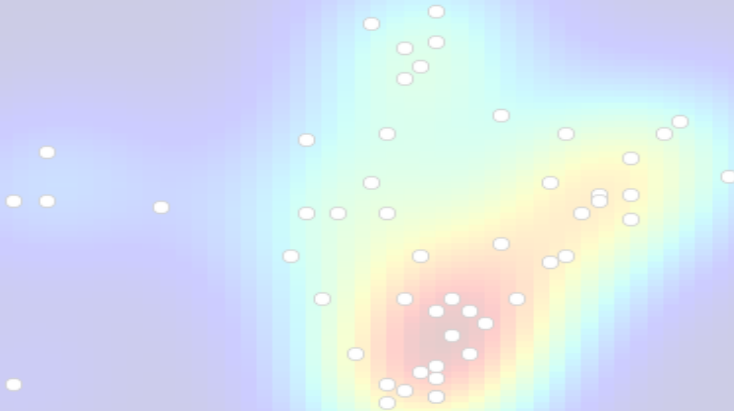
- ① What are Radial Basis Functions (RBFs)?
- ② Discretization of the Collision operator
- ③ Non-linear relaxation problem
- ④ What about the Vlasov-Maxwell part?
- ⑤ Summary and where to proceed

RBFs common in interpolation of scattered data



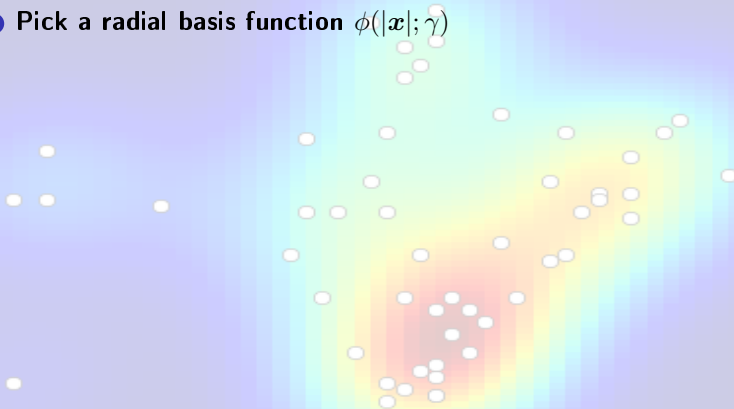
Recipe

- 1 Collect some scattered data $\{x_k, f_k\}$ for $k = 1, 2, \dots, N$.



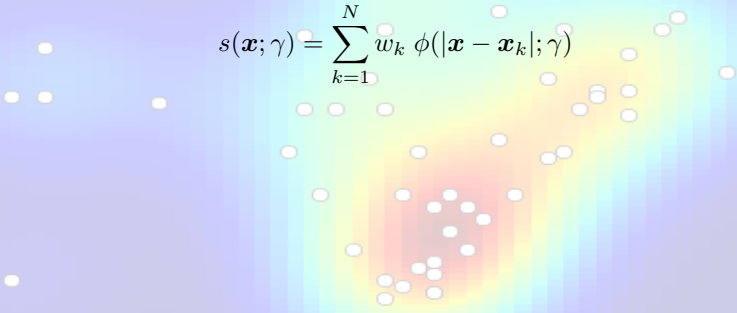
Recipe

- 1 Collect some scattered data $\{\mathbf{x}_k, f_k\}$ for $k = 1, 2, \dots, N$.
- 2 Pick a radial basis function $\phi(|\mathbf{x}|; \gamma)$



Recipe

- 1 Collect some scattered data $\{\mathbf{x}_k, f_k\}$ for $k = 1, 2, \dots, N$.
- 2 Pick a radial basis function $\phi(|\mathbf{x}|; \gamma)$
- 3 Define an RBF interpolant $f(\mathbf{x}) \approx s(\mathbf{x}; \gamma)$

$$s(\mathbf{x}; \gamma) = \sum_{k=1}^N w_k \phi(|\mathbf{x} - \mathbf{x}_k|; \gamma)$$


Recipe

- 1 Collect some scattered data $\{\mathbf{x}_k, f_k\}$ for $k = 1, 2, \dots, N$.
- 2 Pick a radial basis function $\phi(|\mathbf{x}|; \gamma)$
- 3 Define an RBF interpolant $f(\mathbf{x}) \approx s(\mathbf{x}; \gamma)$

$$s(\mathbf{x}; \gamma) = \sum_{k=1}^N w_k \phi(|\mathbf{x} - \mathbf{x}_k|; \gamma)$$

- 4 Compute the coefficients w_k as the solution of the linear system

$$A_{ij} w_j = f_i \quad \text{where} \quad A_{ij} = \phi(|\mathbf{x}_i - \mathbf{x}_j|; \gamma)$$

Recipe

- 1 Collect some scattered data $\{\mathbf{x}_k, f_k\}$ for $k = 1, 2, \dots, N$.
- 2 Pick a radial basis function $\phi(|\mathbf{x}|; \gamma)$
- 3 Define an RBF interpolant $f(\mathbf{x}) \approx s(\mathbf{x}; \gamma)$

$$s(\mathbf{x}; \gamma) = \sum_{k=1}^N w_k \phi(|\mathbf{x} - \mathbf{x}_k|; \gamma)$$

- 4 Compute the coefficients w_k as the solution of the linear system

$$A_{ij} w_j = f_i \quad \text{where} \quad A_{ij} = \phi(|\mathbf{x}_i - \mathbf{x}_j|; \gamma)$$

- 5 Evaluate the interpolant $s(\mathbf{x}; \gamma)$ where you wish

Many different choices for basis functions

Gaussian: $\phi(|\mathbf{x}|; \gamma) = e^{-(\gamma|\mathbf{x}|)^2}$

Polyharmonic Spline: $\phi(|\mathbf{x}|; \gamma) = \gamma|\mathbf{x}|, (\gamma|\mathbf{x}|)^3, \dots$

Multiquadric: $\phi(|\mathbf{x}|; \gamma) = \sqrt{1 + (\gamma|\mathbf{x}|)^2}$

Inverse Multiquadric: $\phi(|\mathbf{x}|; \gamma) = \frac{1}{\sqrt{1 + (\gamma|\mathbf{x}|)^2}}$

Inverse Quadratic: $\phi(|\mathbf{x}|; \gamma) = \frac{1}{1 + (\gamma|\mathbf{x}|)^2}$

Thin-Plate Spline: $\phi(|\mathbf{x}|; \gamma) = (\gamma|\mathbf{x}|)^2 \ln(\gamma|\mathbf{x}|)$

Many different choices for basis functions

Gaussian: $\phi(|\mathbf{x}|; \gamma) = e^{-(\gamma|\mathbf{x}|)^2}$

Polyharmonic Spline: $\phi(|\mathbf{x}|; \gamma) = \gamma|\mathbf{x}|, (\gamma|\mathbf{x}|)^3, \dots$

Multiquadric: $\phi(|\mathbf{x}|; \gamma) = \sqrt{1 + (\gamma|\mathbf{x}|)^2}$

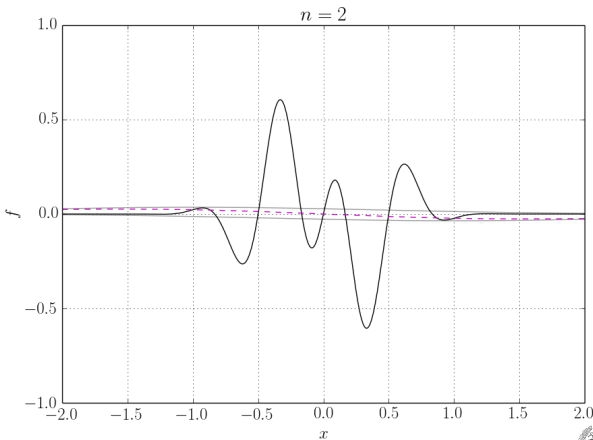
Inverse Multiquadric: $\phi(|\mathbf{x}|; \gamma) = \frac{1}{\sqrt{1 + (\gamma|\mathbf{x}|)^2}}$

Inverse Quadratic: $\phi(|\mathbf{x}|; \gamma) = \frac{1}{1 + (\gamma|\mathbf{x}|)^2}$

Thin-Plate Spline: $\phi(|\mathbf{x}|; \gamma) = (\gamma|\mathbf{x}|)^2 \ln(\gamma|\mathbf{x}|)$

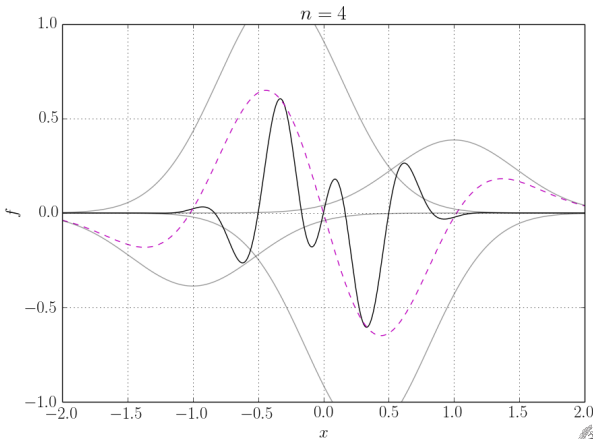
How good can the interpolant be?

$$f(x) \approx \sum_{k=1}^n w_k \exp[-\gamma(x - x_k)^2]$$



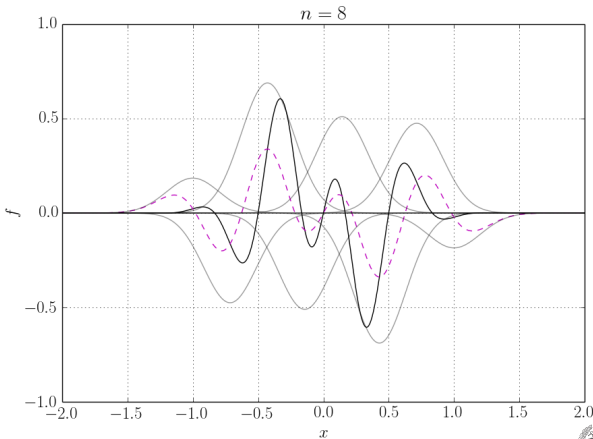
How good can the interpolant be?

$$f(x) \approx \sum_{k=1}^n w_k \exp[-\gamma(x - x_k)^2]$$



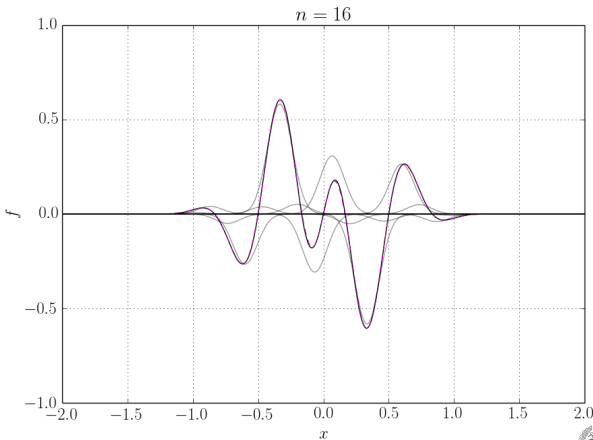
How good can the interpolant be?

$$f(x) \approx \sum_{k=1}^n w_k \exp[-\gamma(x - x_k)^2]$$



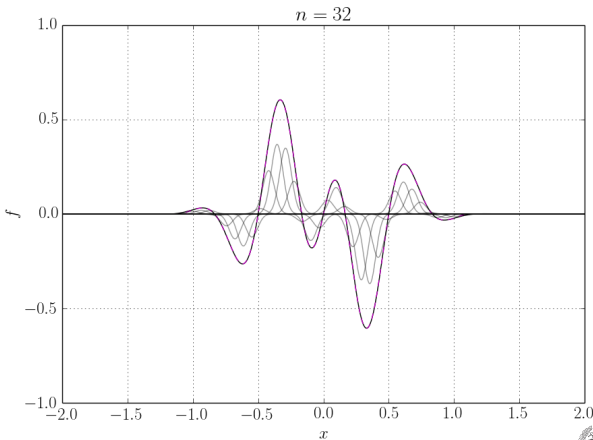
How good can the interpolant be?

$$f(x) \approx \sum_{k=1}^n w_k \exp[-\gamma(x - x_k)^2]$$



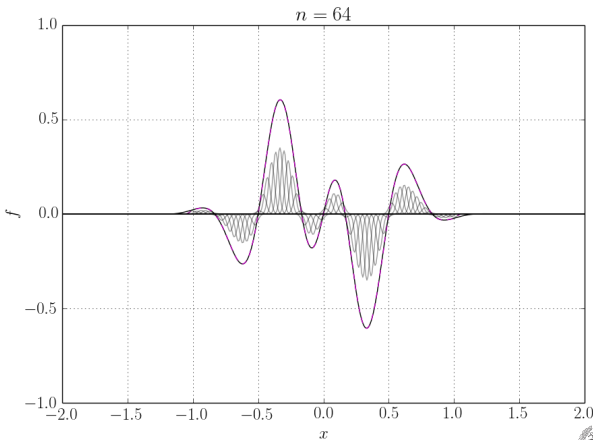
How good can the interpolant be?

$$f(x) \approx \sum_{k=1}^n w_k \exp[-\gamma(x - x_k)^2]$$



How good can the interpolant be?

$$f(x) \approx \sum_{k=1}^n w_k \exp[-\gamma(x - x_k)^2]$$



Outline

- ① What are Radial Basis Functions (RBFs)?
- ② Discretization of the Collision operator**
- ③ Non-linear relaxation problem
- ④ What about the Vlasov-Maxwell part?
- ⑤ Summary and where to proceed

Statistical description of collisions in plasmas

1 Fokker-Planck operator:

$$C_{s\bar{s}}[f_s, f_{\bar{s}}] = \frac{\partial}{\partial v^i} \left(D_{s\bar{s}}^{ij} \frac{\partial f_s}{\partial v^j} - K_{s\bar{s}}^i f_s \right)$$

Statistical description of collisions in plasmas

1 Fokker-Planck operator:

$$C_{s\bar{s}}[f_s, f_{\bar{s}}] = \frac{\partial}{\partial v^i} \left(D_{s\bar{s}}^{ij} \frac{\partial f_{\bar{s}}}{\partial v^j} - K_{s\bar{s}}^i f_s \right)$$

2 Friction and diffusion coefficients

$$K_{s\bar{s}}^i = -\gamma_{s\bar{s}} \frac{m_s}{m_{\bar{s}}} \frac{\partial \phi_{\bar{s}}}{\partial v^i}, \quad D_{s\bar{s}}^{ij} = -\gamma_{s\bar{s}} \frac{\partial^2 \psi_{\bar{s}}}{\partial v^i \partial v^j}, \quad \gamma_{s\bar{s}} = \left(\frac{e_s e_{\bar{s}}}{m_s \epsilon_0} \right)^2 \ln \Lambda.$$

Statistical description of collisions in plasmas

1 Fokker-Planck operator:

$$C_{s\bar{s}}[f_s, f_{\bar{s}}] = \frac{\partial}{\partial v^i} \left(D_{s\bar{s}}^{ij} \frac{\partial f_s}{\partial v^j} - K_{s\bar{s}}^i f_s \right)$$

2 Friction and diffusion coefficients

$$K_{s\bar{s}}^i = -\gamma_{s\bar{s}} \frac{m_s}{m_{\bar{s}}} \frac{\partial \phi_{\bar{s}}}{\partial v^i}, \quad D_{s\bar{s}}^{ij} = -\gamma_{s\bar{s}} \frac{\partial^2 \psi_{\bar{s}}}{\partial v^i \partial v^j}, \quad \gamma_{s\bar{s}} = \left(\frac{e_s e_{\bar{s}}}{m_s \epsilon_0} \right)^2 \ln \Lambda.$$

3 Rosenbluth potentials

$$\phi_{\bar{s}}(\mathbf{v}) = -\frac{1}{4\pi} \int d\bar{\mathbf{v}} f_{\bar{s}}(\bar{\mathbf{v}}) |\mathbf{v} - \bar{\mathbf{v}}|^{-1},$$

$$\psi_{\bar{s}}(\mathbf{v}) = -\frac{1}{8\pi} \int d\bar{\mathbf{v}} f_{\bar{s}}(\bar{\mathbf{v}}) |\mathbf{v} - \bar{\mathbf{v}}|,$$

Statistical description of collisions in plasmas

1 Fokker-Planck operator:

$$C_{s\bar{s}}[f_s, f_{\bar{s}}] = \frac{\partial}{\partial v^i} \left(D_{s\bar{s}}^{ij} \frac{\partial f_s}{\partial v^j} - K_{s\bar{s}}^i f_s \right)$$

2 Friction and diffusion coefficients

$$K_{s\bar{s}}^i = -\gamma_{s\bar{s}} \frac{m_s}{m_{\bar{s}}} \frac{\partial \phi_{\bar{s}}}{\partial v^i}, \quad D_{s\bar{s}}^{ij} = -\gamma_{s\bar{s}} \frac{\partial^2 \psi_{\bar{s}}}{\partial v^i \partial v^j}, \quad \gamma_{s\bar{s}} = \left(\frac{e_s e_{\bar{s}}}{m_s \epsilon_0} \right)^2 \ln \Lambda.$$

3 Rosenbluth potentials **are difficult to compute**

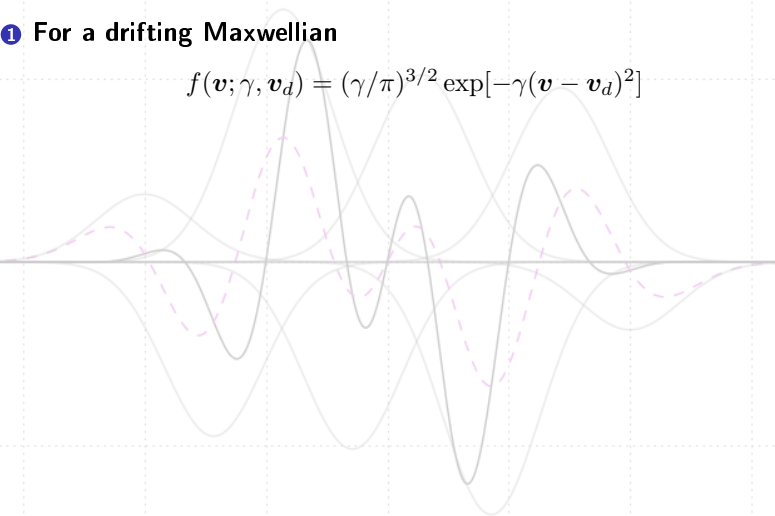
$$\phi_{\bar{s}}(\mathbf{v}) = -\frac{1}{4\pi} \int d\bar{\mathbf{v}} f_{\bar{s}}(\bar{\mathbf{v}}) |\mathbf{v} - \bar{\mathbf{v}}|^{-1},$$

$$\psi_{\bar{s}}(\mathbf{v}) = -\frac{1}{8\pi} \int d\bar{\mathbf{v}} f_{\bar{s}}(\bar{\mathbf{v}}) |\mathbf{v} - \bar{\mathbf{v}}|,$$

A special case relaxes difficulties

① For a drifting Maxwellian

$$f(\mathbf{v}; \gamma, \mathbf{v}_d) = (\gamma/\pi)^{3/2} \exp[-\gamma(\mathbf{v} - \mathbf{v}_d)^2]$$



A special case relaxes difficulties

- 1 For a drifting Maxwellian

$$f(\mathbf{v}; \gamma, \mathbf{v}_d) = (\gamma/\pi)^{3/2} \exp[-\gamma(\mathbf{v} - \mathbf{v}_d)^2]$$

- 2 we have the Rosenbluth potentials

$$\varphi(\mathbf{v}; \gamma, \mathbf{v}_d) = -\frac{\sqrt{\gamma}}{4\pi} \Phi(\sqrt{\gamma}|\mathbf{v} - \mathbf{v}_d|),$$

$$\psi(\mathbf{v}; \gamma, \mathbf{v}_d) = -\frac{1}{8\pi\sqrt{\gamma}} \Psi(\sqrt{\gamma}|\mathbf{v} - \mathbf{v}_d|)$$

A special case relaxes difficulties

- 1 For a drifting Maxwellian

$$f(\mathbf{v}; \gamma, \mathbf{v}_d) = (\gamma/\pi)^{3/2} \exp[-\gamma(\mathbf{v} - \mathbf{v}_d)^2]$$

- 2 we have the Rosenbluth potentials

$$\varphi(\mathbf{v}; \gamma, \mathbf{v}_d) = -\frac{\sqrt{\gamma}}{4\pi} \Phi(\sqrt{\gamma}|\mathbf{v} - \mathbf{v}_d|),$$

$$\psi(\mathbf{v}; \gamma, \mathbf{v}_d) = -\frac{1}{8\pi\sqrt{\gamma}} \Psi(\sqrt{\gamma}|\mathbf{v} - \mathbf{v}_d|)$$

- 3 where the Φ and Ψ are

$$\Phi(y) = \frac{\operatorname{erf}(y)}{y},$$

$$\Psi(y) = \left[y + \frac{1}{2y} \right] \operatorname{erf}(y) + \frac{\exp(-y^2)}{\sqrt{\pi}}$$

A special case relaxes difficulties

- 1 For a drifting Maxwellian

$$f(\mathbf{v}; \gamma, \mathbf{v}_d) = (\gamma/\pi)^{3/2} \exp[-\gamma(\mathbf{v} - \mathbf{v}_d)^2]$$

- 2 we have the Rosenbluth potentials

$$\varphi(\mathbf{v}; \gamma, \mathbf{v}_d) = -\frac{\sqrt{\gamma}}{4\pi} \Phi(\sqrt{\gamma}|\mathbf{v} - \mathbf{v}_d|),$$

$$\psi(\mathbf{v}; \gamma, \mathbf{v}_d) = -\frac{1}{8\pi\sqrt{\gamma}} \Psi(\sqrt{\gamma}|\mathbf{v} - \mathbf{v}_d|)$$

- 3 where the Φ and Ψ are **easy to compute and differentiate**

$$\Phi(y) = \frac{\operatorname{erf}(y)}{y},$$

$$\Psi(y) = \left[y + \frac{1}{2y} \right] \operatorname{erf}(y) + \frac{\exp(-y^2)}{\sqrt{\pi}}$$

Our idea for easier computations

- 1 Express the distribution function as a *convolution*

$$f_s(\mathbf{v}, t) = \int \left(\frac{\gamma}{\pi}\right)^{3/2} \exp[-\gamma(\mathbf{v} - \mathbf{u})^2] W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma,$$

$$n_s(t) = \int W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma$$

Our idea for easier computations

- 1 Express the distribution function as a *convolution*

$$f_s(\mathbf{v}, t) = \int \left(\frac{\gamma}{\pi}\right)^{3/2} \exp[-\gamma(\mathbf{v} - \mathbf{u})^2] W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma,$$

$$n_s(t) = \int W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma$$

- 2 Compute the Rosenbluth potentials

$$\varphi_s(\mathbf{v}, t) = -\frac{1}{4\pi} \int \gamma^{1/2} \Phi(\gamma^{1/2}|\mathbf{v} - \mathbf{u}|) W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma,$$

$$\psi_s(\mathbf{v}, t) = -\frac{1}{8\pi} \int \gamma^{-1/2} \Psi(\gamma^{1/2}|\mathbf{v} - \mathbf{u}|) W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma,$$

Our idea for easier computations

- 1 Express the distribution function as a *convolution*

$$f_s(\mathbf{v}, t) = \int \left(\frac{\gamma}{\pi}\right)^{3/2} \exp[-\gamma(\mathbf{v} - \mathbf{u})^2] W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma,$$

$$n_s(t) = \int W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma$$

- 2 Compute the Rosenbluth potentials

$$\varphi_s(\mathbf{v}, t) = -\frac{1}{4\pi} \int \gamma^{1/2} \Phi(\gamma^{1/2}|\mathbf{v} - \mathbf{u}|) W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma,$$

$$\psi_s(\mathbf{v}, t) = -\frac{1}{8\pi} \int \gamma^{-1/2} \Psi(\gamma^{1/2}|\mathbf{v} - \mathbf{u}|) W_s(\mathbf{u}, \gamma, t) d\mathbf{u} d\gamma,$$

- 3 Define the weight function

$$(3\text{-D}) W_s(\mathbf{u}, \gamma, t) = \sum_i w_s^i(t) \delta(\mathbf{u} - \mathbf{v}_s^i) \delta(\gamma - \gamma_s^i),$$

$$(2\text{-D}) W_s(\mathbf{u}, \gamma, t) = \sum_i w_s^i(t) \frac{1}{2\pi v_{s,\perp}^i} \delta(u_{\parallel} - v_{s,\parallel}^i) \delta(u_{\perp} - v_{s,\perp}^i) \delta(\gamma - \gamma_s^i)$$

We get a “mesh-free” approach

① Choice of weight function gives the RBF basis functions

$$(3\text{-D}) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} \exp[-\gamma_s^i(\mathbf{v} - \mathbf{v}_s^i)^2]$$

$$(2\text{-D}) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} I_0(2\gamma_s^i v_{s,\perp}^i v_\perp) e^{-\gamma_s^i(v_\parallel - v_{s,\parallel}^i)^2 - \gamma_s^i(v_\perp^2 + (v_{s,\perp}^i)^2)}$$

We get a “mesh-free” approach

- Choice of weight function gives the RBF basis functions

$$(3\text{-D}) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} \exp[-\gamma_s^i(\mathbf{v} - \mathbf{v}_s^i)^2]$$

$$(2\text{-D}) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} I_0(2\gamma_s^i v_{s,\perp}^i v_\perp) e^{-\gamma_s^i(v_\parallel - v_{s,\parallel}^i)^2 - \gamma_s^i(v_\perp^2 + (v_{s,\perp}^i)^2)}$$

- Discretized distribution function

$$f_s(\mathbf{v}, t) = \sum_i w_s^i(t) F_s^i(\mathbf{v})$$

We get a “mesh-free” approach

- Choice of weight function gives the RBF basis functions

$$(3\text{-D}) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} \exp[-\gamma_s^i(\mathbf{v} - \mathbf{v}_s^i)^2]$$

$$(2\text{-D}) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} I_0(2\gamma_s^i v_{s,\perp}^i v_\perp) e^{-\gamma_s^i(v_\parallel - v_{s,\parallel}^i)^2 - \gamma_s^i(v_\perp^2 + (v_{s,\perp}^i)^2)}$$

- Discretized distribution function

$$f_s(\mathbf{v}, t) = \sum_i w_s^i(t) F_s^i(\mathbf{v})$$

- Discretized potentials

$$\varphi_s = \sum_i w_s^i(t) \varphi_s^i(\mathbf{v})$$

$$\psi_s = \sum_i w_s^i(t) \psi_s^i(\mathbf{v})$$

We get a “mesh-free” approach

- Choice of weight function gives the RBF basis functions

$$(3\text{-D}) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} \exp[-\gamma_s^i(\mathbf{v} - \mathbf{v}_s^i)^2]$$

$$(2\text{-D}) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} I_0(2\gamma_s^i v_{s,\perp}^i v_\perp) e^{-\gamma_s^i(v_\parallel - v_{s,\parallel}^i)^2 - \gamma_s^i(v_\perp^2 + (v_{s,\perp}^i)^2)}$$

- Discretized distribution function

$$f_s(\mathbf{v}, t) = \sum_i w_s^i(t) F_s^i(\mathbf{v})$$

- Discretized potentials

$$\varphi_s = \sum_i w_s^i(t) \varphi_s^i(\mathbf{v})$$

$$\psi_s = \sum_i w_s^i(t) \psi_s^i(\mathbf{v})$$

What about time evolution?

Collisional evolution

1 From the continuous representation

$$\frac{\partial f_s}{\partial t} = \sum_{\bar{s}} C_{s\bar{s}}[f_s, f_{\bar{s}}]$$

Collisional evolution

1 From the continuous representation

$$\frac{\partial f_s}{\partial t} = \sum_{\bar{s}} C_{s\bar{s}}[f_s, f_{\bar{s}}]$$

2 to a bilinear expansion

$$\sum_i F_s^i(\mathbf{v}) \frac{\partial w_s^i}{\partial t} = \sum_{k, \ell, \bar{s}} w_s^k(t) w_{\bar{s}}^\ell(t) C_{s\bar{s}}^{k\ell}(\mathbf{v})$$

Collisional evolution

1 From the continuous representation

$$\frac{\partial f_s}{\partial t} = \sum_{\bar{s}} C_{s\bar{s}}[f_s, f_{\bar{s}}]$$

2 to a bilinear expansion

$$\sum_i F_s^i(\mathbf{v}) \frac{\partial w_s^i}{\partial t} = \sum_{k, \ell, \bar{s}} w_s^k(t) w_{\bar{s}}^\ell(t) C_{s\bar{s}}^{k\ell}(\mathbf{v})$$

3 with an analytical coefficient

$$C_{s\bar{s}}^{k\ell}(\mathbf{v}) = \gamma_{s\bar{s}} \left[\frac{m_s}{m_{\bar{s}}} F_s^k F_{\bar{s}}^\ell + \mu_{s\bar{s}} \frac{\partial \varphi_{\bar{s}}^\ell}{\partial \mathbf{v}} \cdot \frac{\partial F_s^k}{\partial \mathbf{v}} - \frac{\partial^2 \psi_{\bar{s}}^\ell}{\partial \mathbf{v} \partial \mathbf{v}} : \frac{\partial^2 F_s^k}{\partial \mathbf{v} \partial \mathbf{v}} \right],$$

where $\mu_{s\bar{s}} = m_s/m_{\bar{s}} - 1$

Equation for the weights

- 1 Apply center collocation, i.e., evaluate the expanded equation at the points v_s^i

$$\sum_j \mathcal{M}_s^{ij} \frac{\partial w_s^j}{\partial t} = \sum_{k, \ell, \bar{s}} w_s^k w_{\bar{s}}^{\ell} C_{s\bar{s}}^{ik\ell},$$

Equation for the weights

- 1 Apply center collocation, i.e., evaluate the expanded equation at the points v_s^i

$$\sum_j \mathcal{M}_s^{ij} \frac{\partial w_s^j}{\partial t} = \sum_{k,\ell,\bar{s}} w_s^k w_{\bar{s}}^{\ell} C_{s\bar{s}}^{ik\ell},$$

- 2 Define the “RBF matrix” and the “collision kernel”

$$\mathcal{M}_s^{ij} = F_s^j(v_s^i), \quad C_{s\bar{s}}^{ik\ell} = C_{s\bar{s}}^{k\ell}(v_s^i)$$

Equation for the weights

- 1 Apply center collocation, i.e., evaluate the expanded equation at the points v_s^i

$$\sum_j \mathcal{M}_s^{ij} \frac{\partial w_s^j}{\partial t} = \sum_{k, \ell, \bar{s}} w_s^k w_{\bar{s}}^{\ell} C_{s\bar{s}}^{ik\ell},$$

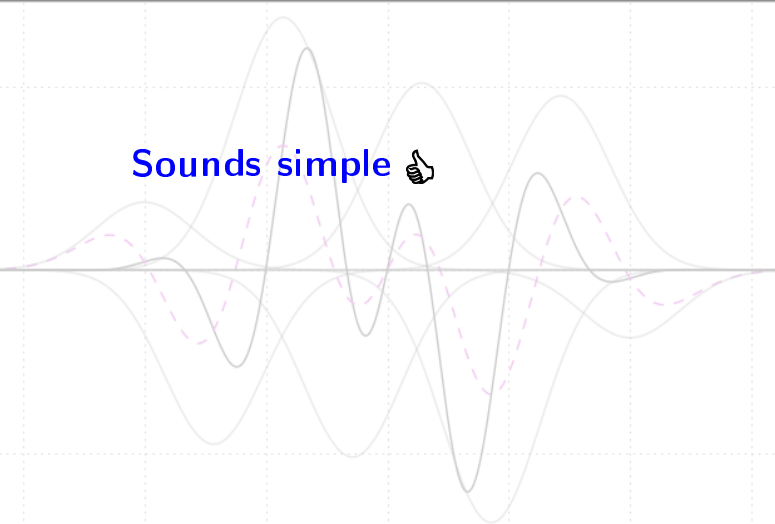
- 2 Define the “RBF matrix” and the “collision kernel”

$$\mathcal{M}_s^{ij} = F_s^j(v_s^i), \quad C_{s\bar{s}}^{ik\ell} = C_{s\bar{s}}^{k\ell}(v_s^i)$$

We chose center collocation because of its simplicity.
Galerkin projection could also be worth considering.

The RBF method

Sounds simple 👍



The RBF method

Sounds simple 👍

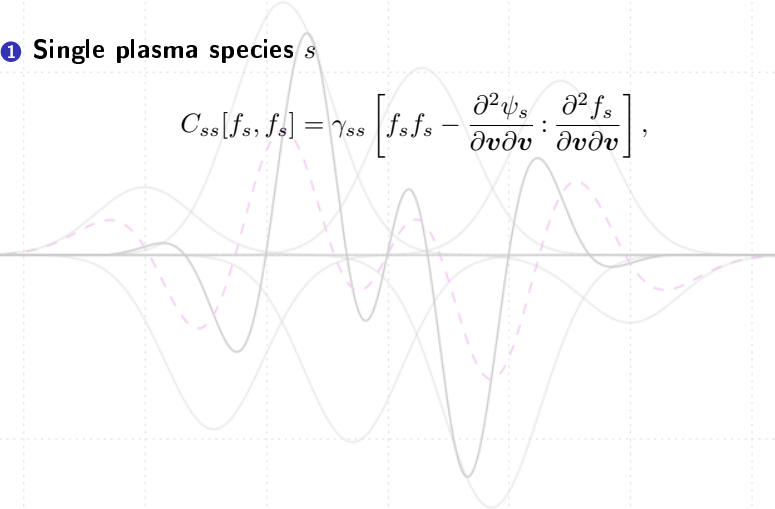
Does it work?

Outline

- ① What are Radial Basis Functions (RBFs)?
- ② Discretization of the Collision operator
- ③ Non-linear relaxation problem**
- ④ What about the Vlasov-Maxwell part?
- ⑤ Summary and where to proceed

Non-linear relaxation problem

① Single plasma species s

$$C_{ss}[f_s, f_s] = \gamma_{ss} \left[f_s f_s - \frac{\partial^2 \psi_s}{\partial v \partial v} : \frac{\partial^2 f_s}{\partial v \partial v} \right],$$


Non-linear relaxation problem

1 Single plasma species s

$$C_{ss}[f_s, f_s] = \gamma_{ss} \left[f_s f_s - \frac{\partial^2 \psi_s}{\partial v \partial v} : \frac{\partial^2 f_s}{\partial v \partial v} \right],$$

2 Normalize time

$$\tau = \gamma_{ss} t = \left(\frac{e_s^2}{m_s \epsilon_0} \right)^2 \ln \Lambda t$$

Non-linear relaxation problem

1 Single plasma species s

$$C_{ss}[f_s, f_s] = \gamma_{ss} \left[f_s f_s - \frac{\partial^2 \psi_s}{\partial v \partial v} : \frac{\partial^2 f_s}{\partial v \partial v} \right],$$

2 Normalize time

$$\tau = \gamma_{ss} t = \left(\frac{e_s^2}{m_s \epsilon_0} \right)^2 \ln \Lambda t$$

3 Solve the initial value problem

$$\frac{\partial f_s}{\partial \tau} = f_s f_s - \frac{\partial^2 \psi_s}{\partial v \partial v} : \frac{\partial^2 f_s}{\partial v \partial v},$$

Non-linear relaxation problem

1 Single plasma species s

$$C_{ss}[f_s, f_s] = \gamma_{ss} \left[f_s f_s - \frac{\partial^2 \psi_s}{\partial v \partial v} : \frac{\partial^2 f_s}{\partial v \partial v} \right],$$

2 Normalize time

$$\tau = \gamma_{ss} t = \left(\frac{e_s^2}{m_s \epsilon_0} \right)^2 \ln \Lambda t$$

3 Solve the initial value problem **by applying the RBF expansion**

$$\frac{\partial f_s}{\partial \tau} = f_s f_s - \frac{\partial^2 \psi_s}{\partial v \partial v} : \frac{\partial^2 f_s}{\partial v \partial v},$$

Initial state

- 1 Choose a non-linear double peaked distribution function

$$f_s(\mathbf{v}, \tau_0) = \sum_{i=1}^2 \exp[-\beta_i (\mathbf{v} - \mathbf{v}^i)^2]$$

with $\mathbf{v}^{1,2} = (\pm 3, 0, 0)$ and
 $\beta_{1,2} = 1/5$

Initial state

- 1 Choose a non-linear double peaked distribution function

$$f_s(\mathbf{v}, \tau_0) = \sum_{i=1}^2 \exp[-\beta_i (\mathbf{v} - \mathbf{v}^i)^2]$$

with $\mathbf{v}^{1,2} = (\pm 3, 0, 0)$ and
 $\beta_{1,2} = 1/5$

- 2 Project the initial state to the RBF basis (with center collocation)

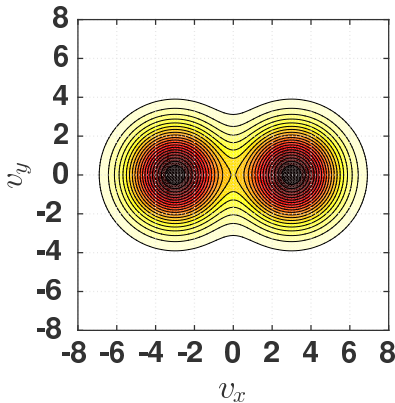
Initial state

- 1 Choose a non-linear double peaked distribution function

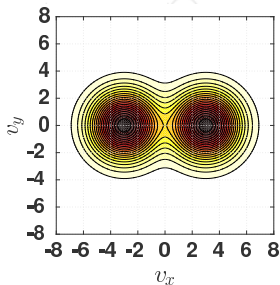
$$f_s(\mathbf{v}, \tau_0) = \sum_{i=1}^2 \exp[-\beta_i (\mathbf{v} - \mathbf{v}^i)^2]$$

with $\mathbf{v}^{1,2} = (\pm 3, 0, 0)$ and
 $\beta_{1,2} = 1/5$

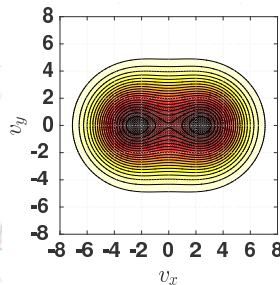
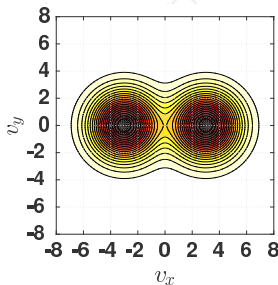
- 2 Project the initial state to the RBF basis (with center collocation)



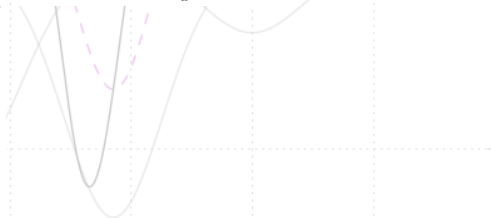
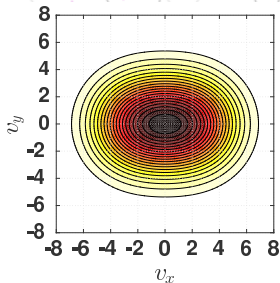
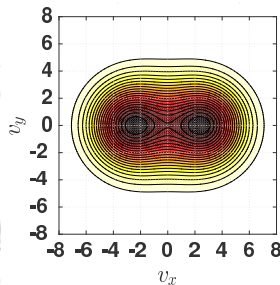
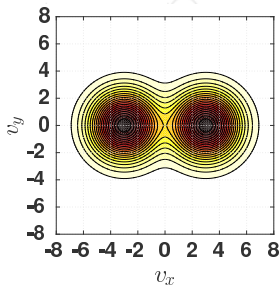
Time slices $\tau = 0, 3, 6, 40$



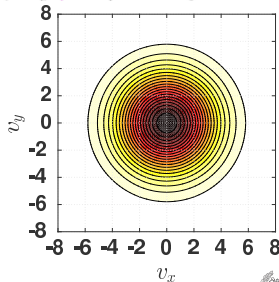
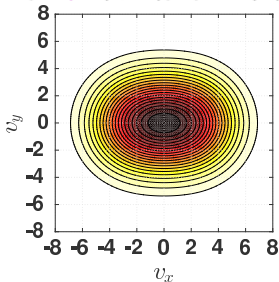
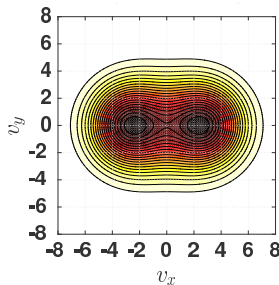
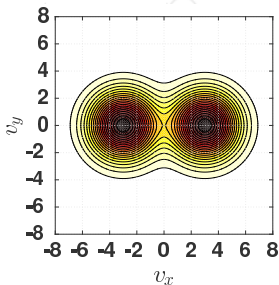
Time slices $\tau = 0, 3, 6, 40$



Time slices $\tau = 0, 3, 6, 40$

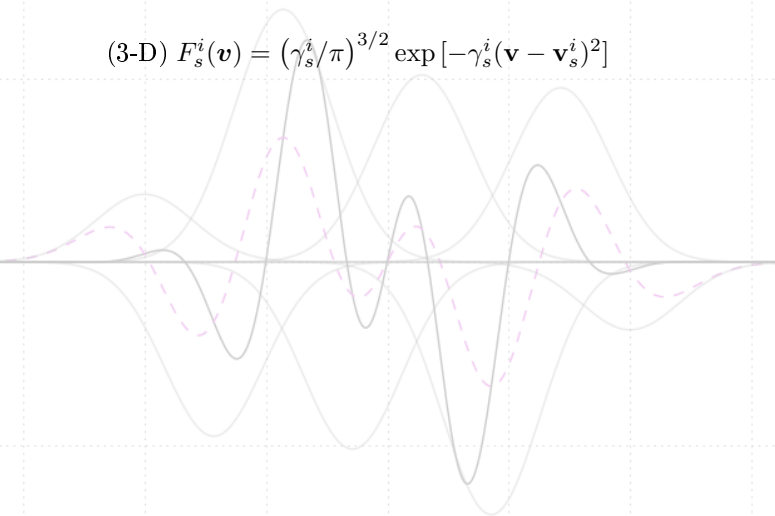


Time slices $\tau = 0, 3, 6, 40$



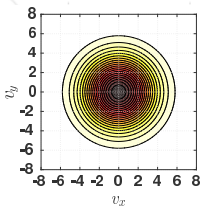
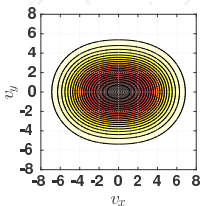
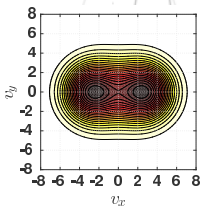
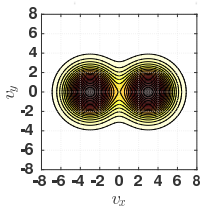
3-D vs 2-D RBF

$$(3\text{-D}) F_s^i(\mathbf{v}) = \left(\gamma_s^i/\pi\right)^{3/2} \exp[-\gamma_s^i(\mathbf{v} - \mathbf{v}_s^i)^2]$$



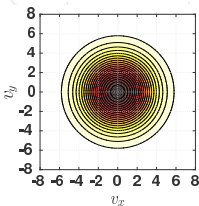
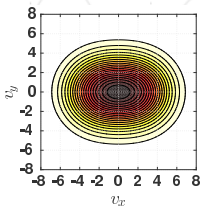
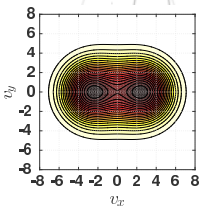
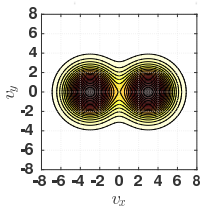
3-D vs 2-D RBF

$$(3-D) F_s^i(\mathbf{v}) = \left(\gamma_s^i/\pi\right)^{3/2} \exp[-\gamma_s^i(\mathbf{v} - \mathbf{v}_s^i)^2]$$



3-D vs 2-D RBF

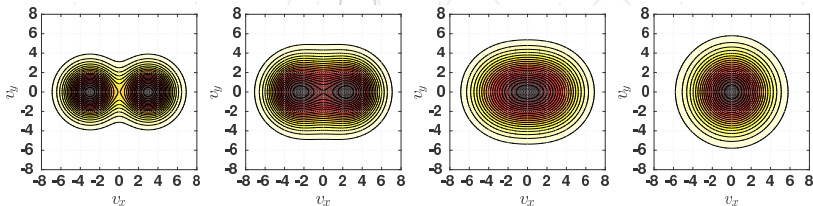
$$(3-D) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} \exp[-\gamma_s^i(\mathbf{v} - \mathbf{v}_s^i)^2]$$



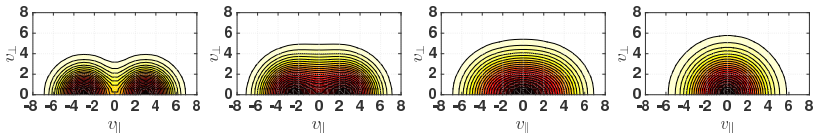
$$(2-D) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} I_0(2\gamma_s^i v_{s,\perp}^i v_\perp) e^{-\gamma_s^i(v_\parallel - v_{s,\parallel}^i)^2 - \gamma_s^i(v_\perp^2 + (v_{s,\perp}^i)^2)}$$

3-D vs 2-D RBF

$$(3-D) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} \exp[-\gamma_s^i(\mathbf{v} - \mathbf{v}_s^i)^2]$$

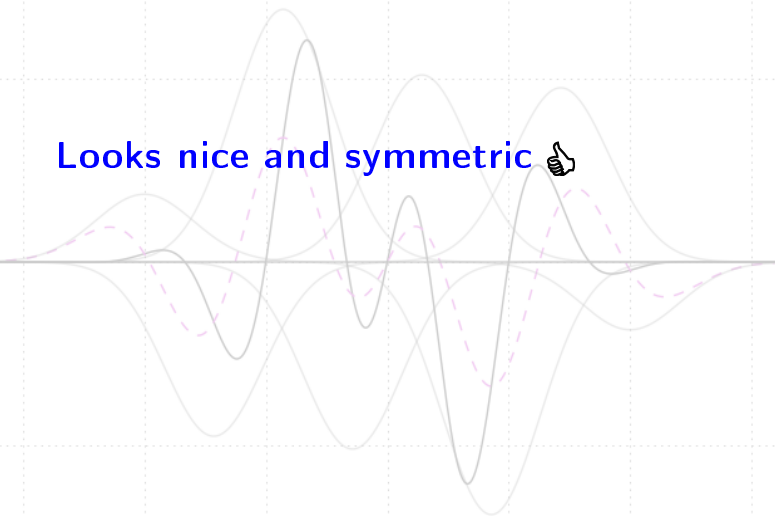


$$(2-D) F_s^i(\mathbf{v}) = (\gamma_s^i/\pi)^{3/2} I_0(2\gamma_s^i v_{s,\perp} v_{\perp}) e^{-\gamma_s^i(v_{\parallel} - v_{s,\parallel})^2 - \gamma_s^i(v_{\perp}^2 + (v_{s,\perp}^2))}$$



The RBF method

Looks nice and symmetric 👍



The RBF method

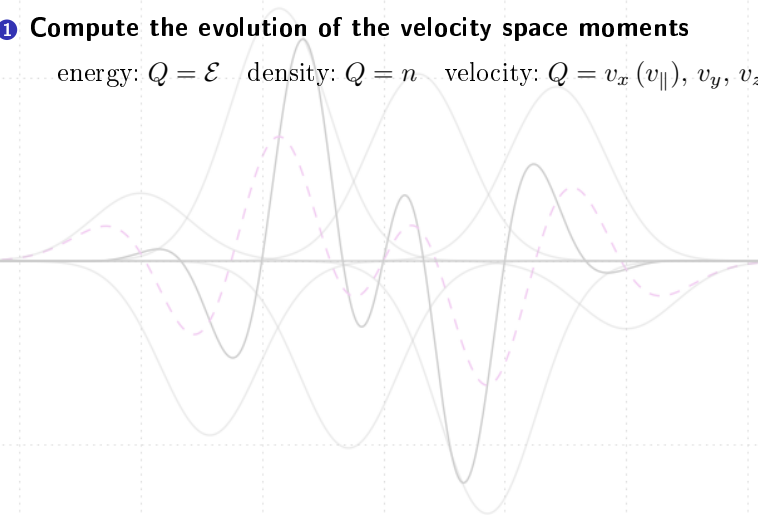
Looks nice and symmetric 👍

What about exact numbers?

Conservation properties

1 Compute the evolution of the velocity space moments

energy: $Q = \mathcal{E}$ density: $Q = n$ velocity: $Q = v_x (v_{\parallel}), v_y, v_z$

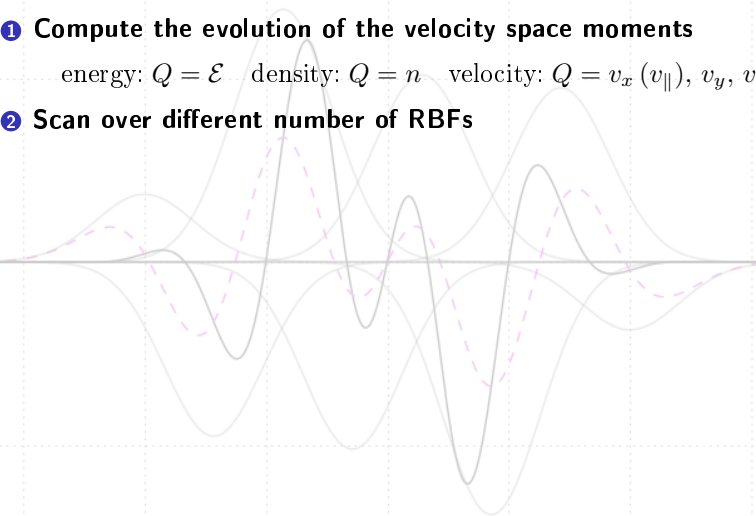


Conservation properties

1 Compute the evolution of the velocity space moments

energy: $Q = \mathcal{E}$ density: $Q = n$ velocity: $Q = v_x (v_{\parallel}), v_y, v_z$

2 Scan over different number of RBFs

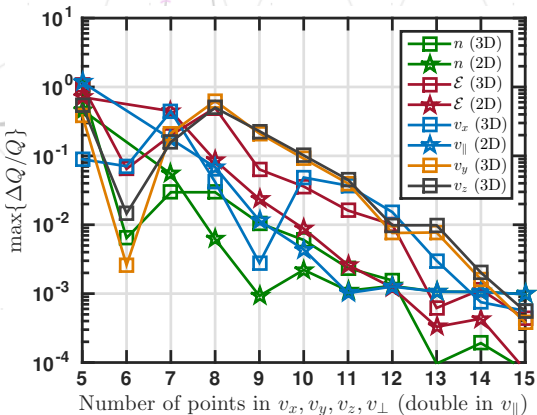


Conservation properties

1 Compute the evolution of the velocity space moments

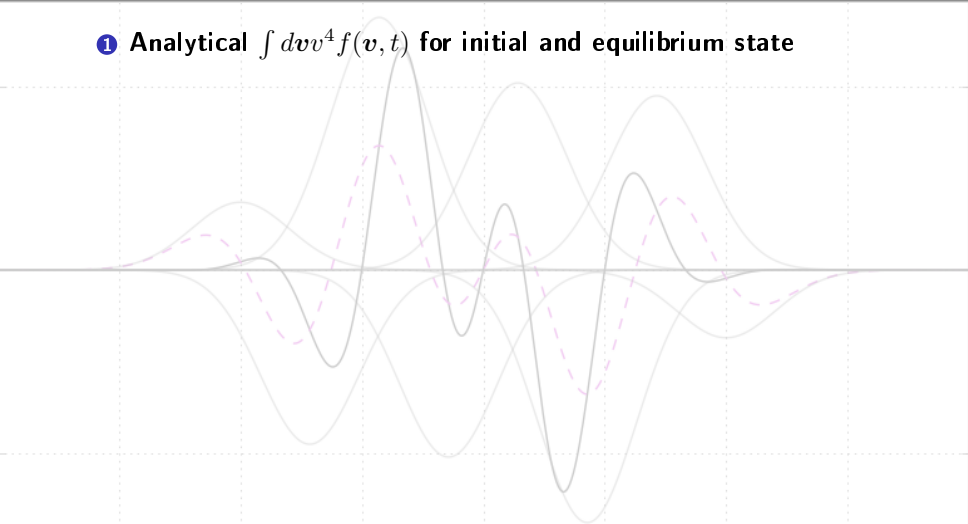
energy: $Q = \mathcal{E}$ density: $Q = n$ velocity: $Q = v_x (v_{\parallel}), v_y, v_z$

2 Scan over different number of RBFs



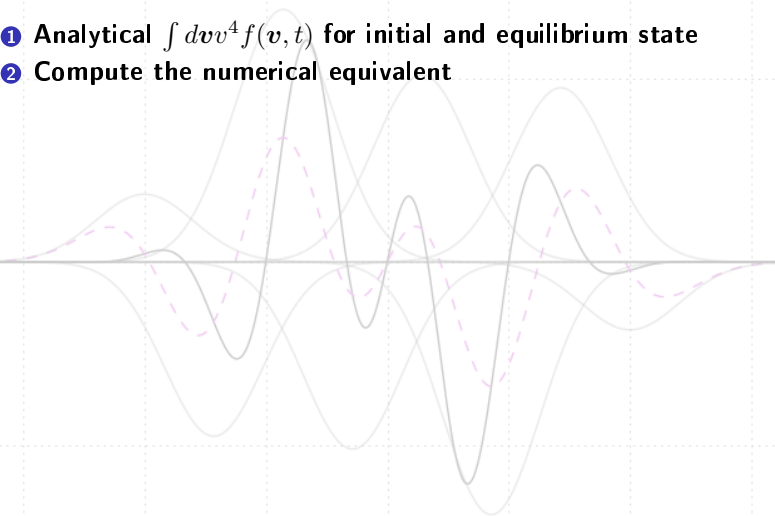
Non-conserved moments:

- 1 Analytical $\int dv v^4 f(v, t)$ for initial and equilibrium state



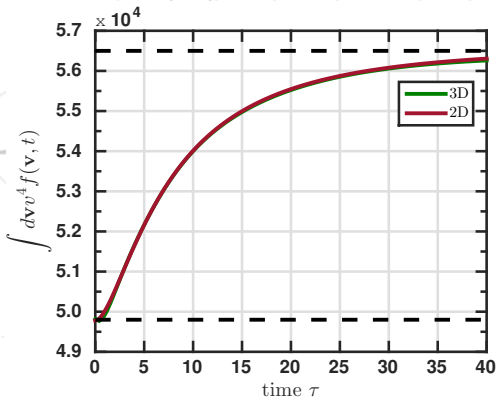
Non-conserved moments:

- 1 Analytical $\int dv v^4 f(v, t)$ for initial and equilibrium state
- 2 Compute the numerical equivalent



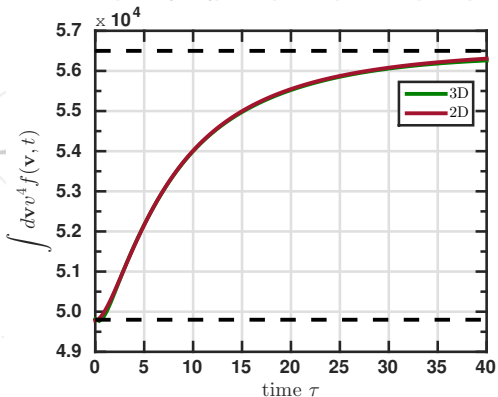
Non-conserved moments:

- 1 Analytical $\int dv v^4 f(v, t)$ for initial and equilibrium state
- 2 Compute the numerical equivalent



Non-conserved moments:

- 1 Analytical $\int dv v^4 f(v, t)$ for initial and equilibrium state
- 2 Compute the numerical equivalent



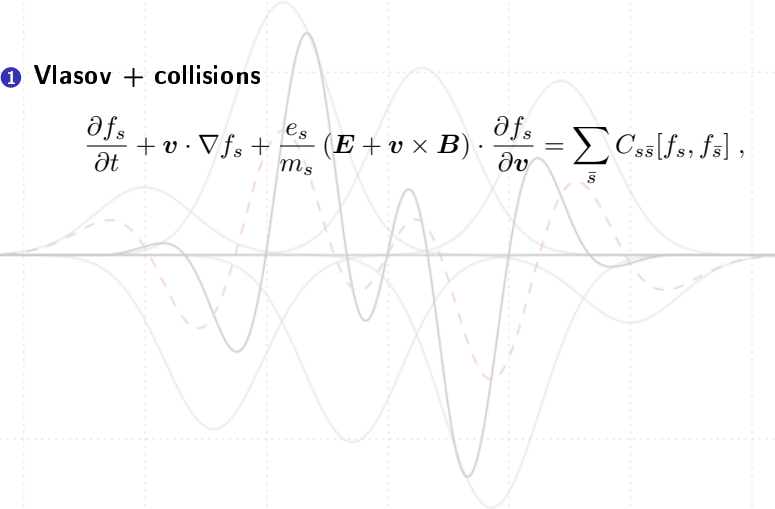
Not bad 👍

Outline

- ① What are Radial Basis Functions (RBFs)?
- ② Discretization of the Collision operator
- ③ Non-linear relaxation problem
- ④ **What about the Vlasov-Maxwell part?**
- ⑤ Summary and where to proceed

MHD, gyrokinetics, you name it!

1 Vlasov + collisions

$$\frac{\partial f_s}{\partial t} + \mathbf{v} \cdot \nabla f_s + \frac{e_s}{m_s} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f_s}{\partial \mathbf{v}} = \sum_{\bar{s}} C_{s\bar{s}}[f_s, f_{\bar{s}}],$$


MHD, gyrokinetics, you name it!

1 Vlasov + collisions

$$\frac{\partial f_s}{\partial t} + \mathbf{v} \cdot \nabla f_s + \frac{e_s}{m_s} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f_s}{\partial \mathbf{v}} = \sum_{\bar{s}} C_{s\bar{s}}[f_s, f_{\bar{s}}],$$

2 Electric and Magnetic fields

$$\nabla \cdot \mathbf{E} = (1/\epsilon_0) \sum_s \int d\mathbf{v} f_s(\mathbf{v}, t)$$

$$\nabla \times \mathbf{B} = \mu_0 \sum_s e_s \int d\mathbf{v} \mathbf{v} f_s(\mathbf{v}, t) + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$$

MHD, gyrokinetics, you name it!

1 Vlasov + collisions

$$\frac{\partial f_s}{\partial t} + \mathbf{v} \cdot \nabla f_s + \frac{e_s}{m_s} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f_s}{\partial \mathbf{v}} = \sum_{\bar{s}} C_{s\bar{s}}[f_s, f_{\bar{s}}],$$

2 Electric and Magnetic fields

$$\nabla \cdot \mathbf{E} = (1/\epsilon_0) \sum_s \int d\mathbf{v} f_s(\mathbf{v}, t)$$

$$\nabla \times \mathbf{B} = \mu_0 \sum_s e_s \int d\mathbf{v} \mathbf{v} f_s(\mathbf{v}, t) + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$$

**RBF already tamed the collision operator.
What about the rest?**

RBF discretization of the Vlasov-Maxwell system

1 RBF Vlasov-Fokker-Planck

$$\sum_j \mathcal{M}_s^{ij} \mathcal{L}_s^{ij} w_s^j = \sum_{k, \ell, \bar{s}} w_s^k w_s^\ell C_{s\bar{s}}^{ik\ell}, \quad \forall i \in 1, 2, 3, \dots, \quad (1)$$

where \mathcal{M}_s^{ij} is the RBF matrix as previously and the operator \mathcal{L}_s^{ij} is

$$\mathcal{L}_s^{ij} \doteq \frac{\partial}{\partial t} + \mathbf{v}_s^i \cdot \nabla + 2\gamma_s^j \frac{e_s}{m_s} \left[(\mathbf{v}_s^j - \mathbf{v}_s^i) \cdot \mathbf{E} + (\mathbf{v}_s^j \times \mathbf{v}_s^i) \cdot \mathbf{B} \right]$$

RBF discretization of the Vlasov-Maxwell system

1 RBF Vlasov-Fokker-Planck

$$\sum_j \mathcal{M}_s^{ij} \mathcal{L}_s^{ij} w_s^j = \sum_{k, \ell, \bar{s}} w_s^k w_{\bar{s}}^{\ell} C_{s\bar{s}}^{ik\ell}, \quad \forall i \in 1, 2, 3, \dots, \quad (1)$$

where \mathcal{M}_s^{ij} is the RBF matrix as previously and the operator \mathcal{L}_s^{ij} is

$$\mathcal{L}_s^{ij} \doteq \frac{\partial}{\partial t} + \mathbf{v}_s^i \cdot \nabla + 2\gamma_s^j \frac{e_s}{m_s} \left[(\mathbf{v}_s^j - \mathbf{v}_s^i) \cdot \mathbf{E} + (\mathbf{v}_s^j \times \mathbf{v}_s^i) \cdot \mathbf{B} \right]$$

2 Maxwell's equations

$$\nabla \cdot \mathbf{E} = (1/\epsilon_0) \sum_{s,i} e_s w_s^i,$$

$$\nabla \times \mathbf{B} = \mu_0 \sum_{s,i} e_s w_s^i \mathbf{v}_s^i + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$$

RBF discretization of the Vlasov-Maxwell system

1 RBF Vlasov-Fokker-Planck

$$\sum_j \mathcal{M}_s^{ij} \mathcal{L}_s^{ij} w_s^j = \sum_{k,\ell,\bar{s}} w_s^k w_{\bar{s}}^{\ell} C_{s\bar{s}}^{ik\ell}, \quad \forall i \in 1, 2, 3, \dots, \quad (1)$$

where \mathcal{M}_s^{ij} is the RBF matrix as previously and the operator \mathcal{L}_s^{ij} is

$$\mathcal{L}_s^{ij} \doteq \frac{\partial}{\partial t} + \mathbf{v}_s^i \cdot \nabla + 2\gamma_s^j \frac{e_s}{m_s} \left[(\mathbf{v}_s^j - \mathbf{v}_s^i) \cdot \mathbf{E} + (\mathbf{v}_s^j \times \mathbf{v}_s^i) \cdot \mathbf{B} \right]$$

2 Maxwell's equations

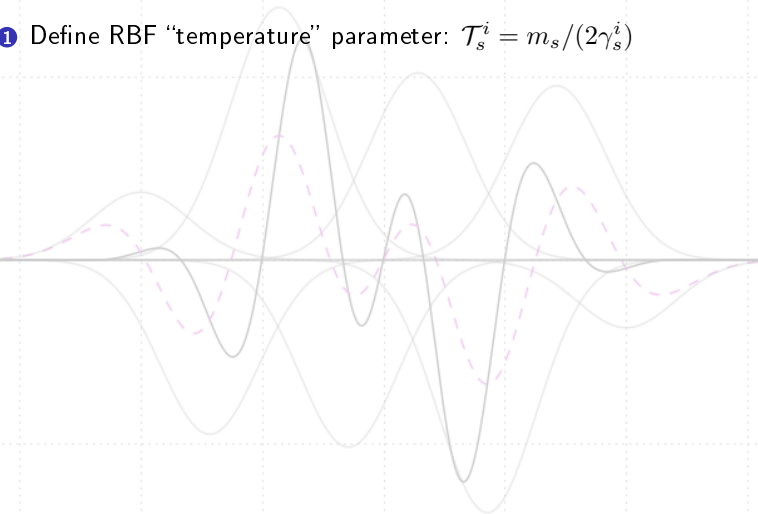
$$\nabla \cdot \mathbf{E} = (1/\epsilon_0) \sum_{s,i} e_s w_s^i,$$

$$\nabla \times \mathbf{B} = \mu_0 \sum_{s,i} e_s w_s^i \mathbf{v}_s^i + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$$

Resemblance to moment equations...

RBF fluid moments

- 1 Define RBF “temperature” parameter: $\mathcal{T}_s^i = m_s / (2\gamma_s^i)$



RBF fluid moments

- 1 Define RBF “temperature” parameter: $\mathcal{T}_s^i = m_s / (2\gamma_s^i)$
- 2 Compute the moments

Density $n_s = \sum_i w_s^i$

RBF fluid moments

- 1 Define RBF “temperature” parameter: $\mathcal{T}_s^i = m_s / (2\gamma_s^i)$
- 2 Compute the moments

Density $n_s = \sum_i w_s^i$

Velocity $n_s \mathbf{V}_s = \sum_i w_s^i \mathbf{v}_s^i$

RBF fluid moments

- 1 Define RBF “temperature” parameter: $\mathcal{T}_s^i = m_s / (2\gamma_s^i)$
- 2 Compute the moments

$$\text{Density } n_s = \sum_i w_s^i$$

$$\text{Velocity } n_s \mathbf{V}_s = \sum_i w_s^i \mathbf{v}_s^i$$

$$\text{Temperature } \frac{3}{2} n_s T_s = \sum_i w_s^i \left[\frac{3}{2} \mathcal{T}_s^i + \frac{1}{2} m_s (\mathbf{v}_s^i - \mathbf{V}_s)^2 \right]$$

$$\text{Momentum flux tensor } \mathbf{\Pi}_s = \sum_i w_s^i m_s \left[\mathcal{T}_s^i \mathbb{I} + \mathbf{v}_s^i \mathbf{v}_s^i \right]$$

$$\text{Energy flux } \mathbf{Q}_s = \sum_i w_s^i \mathbf{v}_s^i \left[\frac{5}{2} \mathcal{T}_s^i + \frac{1}{2} m_s (v_s^i)^2 \right]$$

RBF fluid moments

- 1 Define RBF “temperature” parameter: $\mathcal{T}_s^i = m_s / (2\gamma_s^i)$
- 2 Compute the moments

$$\text{Density } n_s = \sum_i w_s^i$$

$$\text{Velocity } n_s \mathbf{V}_s = \sum_i w_s^i \mathbf{v}_s^i$$

$$\text{Temperature } \frac{3}{2} n_s T_s = \sum_i w_s^i \left[\frac{3}{2} \mathcal{T}_s^i + \frac{1}{2} m_s (\mathbf{v}_s^i - \mathbf{V}_s)^2 \right]$$

$$\text{Momentum flux tensor } \mathbf{\Pi}_s = \sum_i w_s^i m_s \left[\mathcal{T}_s^i \mathbb{I} + \mathbf{v}_s^i \mathbf{v}_s^i \right]$$

$$\text{Energy flux } \mathbf{Q}_s = \sum_i w_s^i \mathbf{v}_s^i \left[\frac{5}{2} \mathcal{T}_s^i + \frac{1}{2} m_s (v_s^i)^2 \right]$$

Easy to compute even for a large number of RBFs

Outline

- ① What are Radial Basis Functions (RBFs)?
- ② Discretization of the Collision operator
- ③ Non-linear relaxation problem
- ④ What about the Vlasov-Maxwell part?
- ⑤ **Summary and where to proceed**

Summary

Introduction to RBFs

- How to interpolate scattered data
- Convergence of the interpolation

New idea to address the collision operator

- Gaussian RBF expansion for the distributions
- Analytic expansion for the non-linear Fokker-Planck operator

Demonstration of a non-linear relaxation problem

- Both 2-D and 3-D RBF methods seem to work
- Conservation properties reasonable
- Consistent with analytical results

Where to proceed

Applied mathematics

- Well-known conditioning problem with the RBF matrix
- Experiment with different mesh configurations
- Galerkin projection instead of the center collocation

Physics

- Heating scenarios
- Loss cones
- Non-linear stuff

Vlasov-Maxwell-Fokker-Planck

- Multispecies scenarios
- Advection terms
- Coupling to field solvers